Neutrino interactions

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VSON8, ICISE, 23 July 2024



Disclaimers

- To prepare for this lecture, I have borrowed from many sources: books, papers, lectures, internet, ... I can not cite them all here.
- Time is limitted, content is unlimitted ! I can not conver all things and go into detail with all contents. If your questions can not be answered in the lecture, discuss later in break or refer to references for more information.
- If you see some complicated formulae and don't even know any thing about them. Don't worry, me too !!!





References

- Carlo Giunti and Chung W. Kim. Fundamentals of Neutrino Physics and Astrophysics. 2007.
- Mark Thomson. Modern particle physics. Cambridge University Press, New York, 2013.
- Rev. Mod. Phys. 84, 1307 (2012) (arXiv:1305.7513v1) From eV to EeV: Neutrino Cross-Sections Across Energy Scales

Contents

- 1. A brief history and properties of neutrinos
- 2. Cross sections
- 3. Neutrino interactions 3.1: Neutrino - electron scattering 3.2: Neutrino - nucleon scattering 3.3: Cross section measurement
 - 3.4: Summary
- 4. Introduction to NEUT an event generator 4.1: Neutrino event generator

 - 4.2: Introduction to NEUT
 - 4.3: Practice with NEUT
 - 4.4: Exercises

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- 2. Cross sections
- 3. Neutrino interactions 3.1: Neutrino - electron scattering 3.2: Neutrino - nucleon scattering 3.3: Cross section measurement 3.4: Summary
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1. A brief history and properties of neutrinos



$$X_Z^A \to Y_{Z-1}^A + e^-$$

$$K_e = (m_X - m_Y - m_e)c^2$$

- Beta spectrum was expected to be discrete as alpha and gamma spectra
- But it was not

$$X_Z^A \to Y_{Z-1}^A + e^- + \overline{\nu}_e$$
$$K_e = (m_X - m_Y - m_e)c^2 - k$$

• Beta energy now can have any value between 0 and maximum



Tips for doing physics and math







1+1=3 Calculations don't add up?

1+1+1=3 Add a hypothetical "dark number" to account for inaccuracy.

Now you're doing math like a physicist!

1. A brief history and properties of neutrinos



"I have done a terrible thing I have postulated a particle that cannot be detected"



- Neutrinos:
- Mass ~ 0: almost no gravitational interaction
- Neutral: no electromagnetic interaction
- No strong interaction
- **Only participate in weak interaction**
- Neutrinos are detectable but they interacts extremely weak with matter



Extremely small cross section

 $\sigma_{\nu} \approx 10^{-62} m^2 - 10^{-35} m^2$



1. A brief history and properties of neutrinos



• An example to see how weak neutrino-matter interaction is



1 light year $\approx 9.5 \times 10^{15} m$

- Need a block of lead with 1-light-year length to stop neutrinos!!!!
- It is a real challenge to study neutrino interactions (to measure cross sections)



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- of particle b:
- $v = v_a + v_b$ ($\overrightarrow{v_a}$, $\overrightarrow{v_b}$: velocities of a and b)
- $N_a = n_a V$: total number of incident particles in volume V
- $N_b = n_b V$: total number of target particles in volume V
- $\phi_a = \frac{N_a}{A\Delta t} = n_a v$: flux of particle a $\left[\frac{1}{[time] \cdot [area^2]}\right]$
- $N = \phi_a N_b \sigma$: #interactions = flux x number of target particles x cross section

• Consider a beam of particles a with flux ϕ_a , pass through a volume $V = Av\Delta t$

 $\phi_a \cdot N_b$



Cross section is a measure of the probability of an interaction to be occured

$$[area^2]$$

- Consider a scattering: $a + b \rightarrow 1 + 2$
- Cross section is related to the transition rate: $\sigma = \frac{f_i}{\phi_a N_b} = \frac{n_a f_i}{n_a n_b (v_a + v_b) V}$
- $v_a + v_b$
- Because the factors of V in the expression for the flux will ultimately be and phase space (~ 1/V), the volume V will not appear in the final result

 $v_a + v_b$





• Normalising the wavefunctions to one particle per unit volume $n_a = n_b = 1/V$

cancelled by the corresponding factors from the wavefunction normalisation

• In terms of matrix element M_{fi}

$$\sigma = \frac{(2\pi)^4}{4\sqrt{(p_a \cdot p_b)^2 - m_a^2 m_b^2}} \int |M_{fi}|^2 \sigma(E_i - E_f) \sigma^3(\vec{p_i} - \vec{p_f}) \frac{d^3 \vec{p_1}}{(2\pi)^3 2E_1} \frac{d^3 \vec{p_2}}{(2\pi)^3 2E_2}$$

• $M_{fi} = \sqrt{2E_a \cdot 2E_b \cdot 2E_1 \cdot 2E_2} T_{fi}$, T_{fi} : transition matrix element

•
$$E_i = E_a + E_b; \ E_f = E_1 + E_2; \ \overrightarrow{p_i} =$$

- Lorentz-invariant flux factor $F = 4(v_a + v_b)E_aE_b = 4\sqrt{(p_a \cdot p_b)^2}$
- In center-of-mass (COM) frame:

$$\sigma_{COM} = \frac{p_f}{64\pi^2 s p_i} \int |N|$$

•
$$\overrightarrow{p_i} = \overrightarrow{p_a} = -\overrightarrow{p_b}$$
 and $\overrightarrow{p_f} = \overrightarrow{p_1} = -\overrightarrow{p_2}$

- $s = (E_a + E_b)^2$ (in COM frame)
- $d\Omega = d(\cos \theta) d\phi$ (θ : polar or zenith angle; ϕ : azimuthal angle)

 $\overrightarrow{p_a} + \overrightarrow{p_b}; \ \overrightarrow{p_f} = \overrightarrow{p_1} + \overrightarrow{p_2}$

$$-m_a^2 m_b^2$$

 $M_{fi}|^2 d\Omega$

 \vec{b}_{2} (in COM frame)

- Differential cross section for solid angle distribution $d\sigma$ $d\Omega$ • In COM: $\frac{d\sigma_{COM}}{d\Omega} = \frac{p_f}{64\pi^2 sp_i} \int |M_{fi}|^2$ Differential cross section: $\frac{d\sigma}{dE}, \frac{d\sigma}{dp}, \frac{d\sigma}{dQ^2}, \frac{d\sigma}{d(\cos\theta)}$
- Double differential cross section: $\frac{d^2\sigma}{dEd\Omega}$, $\frac{d^2\sigma}{d(\cos\theta)d(Q^2)}$



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3. Neutrino interactions



0.1 - 20 GeV

- Basically there are 2 types of interactions:
 - bosons
 - boson



$$i_{CC}^{\mu} = \frac{g_W}{\sqrt{2}} \bar{\psi} \gamma^{\mu} \frac{1}{2} (1 - \gamma^5) q$$

$$j_{NC}^{\mu} = g_Z \bar{\psi} \gamma^{\mu} \frac{1}{2} (g_V - g_A \gamma)$$

• Neutrino - electron elastic scattering:

$$\nu_{\alpha} + e^{-} \rightarrow \nu_{\alpha} + e^{-}$$

 $\overline{\nu}_{\alpha} + e^{-} \rightarrow \overline{\nu}_{\alpha} + e^{-}$
• Neutrino - electron quasi-elastic scatter
 $\nu_{\mu} + e^{-} \rightarrow \nu_{e} + \mu^{-}$ ($E_{\nu} > 10.92 \ GeV$
 $\nu_{\tau} + e^{-} \rightarrow \nu_{e} + \tau^{-}$ ($E_{\nu} > 3089 \ GeV$)
• Consider process: $\nu + A \rightarrow \sum_{X} X$
• A at rest, neglecting neutrino mass, in center of
 $E_{\nu}^{th} = \frac{\left(\sum_{X} m_{X}\right)^{2}}{2m_{A}} - \frac{m_{A}}{2}$

ering:

Did you calculate this value as Sanjib's request?

of mass frame, $s = 2E_{\nu}m_A + m_A^2$ must greater than $\left(\sum_X m_X\right)^{-1}$



$$m_e = 0.511 \ MeV$$

 $m_\mu = 105.66 \ MeV$
 $m_\tau = 1776.9 \ MeV$



o Neutrino - electron elastic scattering





• Neutrino - electron elastic scattering

The numerical values are in units of $10^{-46} \,\mathrm{cm}^2$.

Process	
$\nu_e + e^-$	$\left(G_{\mathrm{F}}^{2} s/4 \pi\right) \Big[\left(1+2 \right.$
$\bar{\nu}_e + e^-$	$\left(G_{\mathrm{F}}^2 s/4\pi\right) \left[\frac{1}{3}\left(1+2\right)\right]$
$ u_{\mu, au} + e^-$	$\left(G_{ m F}^2s/4\pi ight)\left[\left(1-2 ight.$
$ar{ u}_{\mu, au}+e^-$	$\left(G_{\mathrm{F}}^2 s/4 \pi\right) \left[\frac{1}{3} \left(1-2 s\right)\right]$
• $G_F = 1$	$.17 \times 10^{-5} GeV^2$

- $\sin^2 \theta_W = 0.2324 \pm 0.0058 \pm 0.0059$ (1994)

TABLE 5.1. Total neutrino-electron elastic scattering cross-sections for $\sqrt{s} \gg m_e$.

Total cross-section

$$2 \sin^2 \vartheta_{\rm W} \Big)^2 + \frac{4}{3} \sin^4 \vartheta_{\rm W} \Big] \simeq 93 \, s/{\rm MeV^2}$$
$$2 \sin^2 \vartheta_{\rm W} \Big)^2 + 4 \sin^4 \vartheta_{\rm W} \Big] \simeq 39 \, s/{\rm MeV^2}$$
$$2 \sin^2 \vartheta_{\rm W} \Big)^2 + \frac{4}{3} \sin^4 \vartheta_{\rm W} \Big] \simeq 15 \, s/{\rm MeV^2}$$
$$2 \sin^2 \vartheta_{\rm W} \Big)^2 + 4 \sin^4 \vartheta_{\rm W} \Big] \simeq 13 \, s/{\rm MeV^2}$$

: Fermi constant

• $s = (E_{\nu} + E_e)_i^2 = m_e^2 + 2m_e E_{\nu}$: Mandelstam variable

- **o** Neutrino electron elastic scattering
- **o** Differential cross section

$$\frac{d\sigma}{dQ^2} = \frac{G_F^2}{\pi} \left[g_1^2 + g_2^2 \left(1 - \frac{Q^2}{2p_\nu p_e} \right)^2 - g_1 g_2 m_e^2 \frac{Q^2}{2(p_\nu p_e)^2} \right]$$

•
$$g_1(\nu_e) = g_2(\overline{\nu}_e) = \frac{1}{2} + si$$

•
$$g_2(\nu_e) = g_1(\overline{\nu}_e) = \sin^2 \theta_W$$

•
$$g_1(\nu_{\mu,\tau}) = g_2(\overline{\nu}_{\mu,\tau}) = -\frac{1}{2}$$

•
$$g_2(\nu_{\mu,\tau}) = g_1(\overline{\nu}_{\mu,\tau}) = \sin^2$$

 $in^2 \theta_W \approx 0.73$

 $_{V} \approx 0.23$

 $\frac{1}{2} + \sin^2 \theta_W \approx -0.27$

 $e^2 \theta_W \approx 0.23$

• Neutrino - electron quasi-elastic scattering

$$\nu_{\mu} + e^- \rightarrow \nu_e + \mu^- \quad E_{\nu_{\mu}} >$$

o Total cross section

$$\sigma = \frac{G_F^2 s}{\pi} \left(1 - \frac{m_\mu^2}{s} \right)$$
$$s = (p_{\nu_\mu} + p_e)^2$$

10.92 *GeV*





- o Neutrino interactions at 0-1 MeV energy:
 - Coherent scattering (COH):
 - **Neutrino capture on radioative** nuclei:
- o Neutrino interactions at 1-100 MeV energy:
 - **Inverse beta decay**
- **Neutrino interactions at 0.1-20 GeV energy:**
 - **Quasi-elastic scattering (QE):**
 - **Resonant meson production** (**RES**):
 - **Coherent pion production** (COH):
 - **Deep inelastic scattering (DIS):**
- **Neutrino interactions at high and** ultra-high energy (> 20 GeV):
 - **Deep inelastic scattering (DIS):**





O Neutrino interactions at 0-1MeV energy

 $\nu_l + A_N^Z \to \nu_l + A_N^{*Z}$ (QR << 1)

• Coherent scattering (COH):



- $\beta_{\rho}, \beta_{\nu}$: electron, neutrino velocities
- $E_e, p_e, \cos\theta$: electron energy, momentum, scattering angle
- λ^2 : axial-to-vector coupling ratio, $f_V^2(0)$: coupling strength
- $|V_{ud}|^2$: Cabbibo angle
- $F(Z_f, E_e)$: Fermi function

ergy:	• $G_F = 1.17 \times 10^{-5} \ GeV^2$
	 Q: momentum transfer
	 R: nuclear radius
	 T: recoil kinetic energy
	• M_A : target mass
	• E_{ν} : neutrino energy
ıclei:	• $F(Q^2)$: nucleon form factor
	• Q_W : weak current term

Not experimentally measured yet

• Neutrino interactions at 1-100MeV energy:

Inverse beta decay:

 $\overline{\nu}_{\rho} + p \rightarrow e^+ + n$

of neutrino



Scattering process used in Reines & Cowan experiment to confirm the existence

Delayed coincident detection of γ from ^{109}Cd with pair of $\gamma's$ from $e^+ - e^-$ annihilation.

0.5 MeV promt signals of gamma rays

$$e^+ + e^- \rightarrow \gamma + \gamma$$

Delayed gamma signal after $5\mu s$

 $n + {}^{108}Cd \rightarrow {}^{109}Cd^* \rightarrow {}^{109}Cd + \gamma$







• Neutrino interactions at 1-100MeV energy:

• Inverse beta decay:

 $\overline{\nu}_e + p \rightarrow e^+ + n$

$$\frac{d\sigma}{d\cos\theta} = \frac{G_F^2 |V_{ud}|^2 E_e p_e}{2\pi} \left(f_V^2(0)(1 + \frac{1}{2}) - \frac{1}{2} \frac{1}$$

- f_V : nuclear vector (Fermi transition), f_A : axial-vector (Gamov-Teller transition)
- Important for detecting reactor neutrino experiments (1.806 < $E_{\nu} \leq 10 MeV$) Important for understanding supernova explosion mechanism
- $(10 < E_{\nu} \le 20 MeV)$





3.2 Neutrino - nucleon interactions • Neutrino interactions at 0.1-20GeV energy:

Charge current: W^+ exchange

- Quasi-elastic scattering **(target change but not break up):**
 - CCQE: $\nu_l + n \rightarrow l^- + p$
- **Resonant meson production** (target is exited):
 - CCRES:

$$\nu_l + N \rightarrow l^- + \Delta(1232) \rightarrow l^- + N' + m$$

- **Coherent pion production** (target unchange):
 - CCCOH: $\nu_l + A \rightarrow l^- + A + \pi^+$
- **Deep inelastic scattering** (nucleon broken):

• CCDIS:
$$\nu_l + N \rightarrow l^- + X$$

Where $l^- = \{e^-, \mu^-, \tau^-\}; N = \{p, n\}; m = \{\pi, \eta, K\}; A :$ nucleus; X : hadrons

Neutral current: Z^0 exchange

• Elastic scattering
(target change but not break up):
• NC:
$$\nu_l + N \rightarrow \nu_l + N'$$

• Resonant meson production
(target is exited):
• NCRES:
 $\nu_l + N \rightarrow \nu_l + \Delta(1232) \rightarrow \nu_l + N' + m$
• Coherent pion production
(target unchange):
• NCCOH: $\nu_l + A \rightarrow \nu_l + A + \pi^0$
• Deep inelastic scattering
(nucleon broken):
• NCDIS: $\nu_l + N \rightarrow \nu_l + X$



o Anti-neutrino interactions at 0.1-20GeV energy:

Charge current: W^- exchange

Quasi-elastic scattering (target change but not break up):

• CCQE:
$$\overline{\nu}_l + p \rightarrow l^+ + n$$

- **Resonant meson production** (target is exited):
 - CCRES:

$$\overline{\nu}_l + N \rightarrow l^+ + \Delta(1232) \rightarrow l^+ + N' + m$$

- **Coherent pion production** (target unchange):
 - CCCOH: $\overline{\nu}_l + A \rightarrow l^+ + A + \pi^-$
- **Deep inelastic scattering (nucleon broken):**

• CCDIS:
$$\overline{\nu}_l + N \rightarrow l^+ + X$$

Where $l^- = \{e^-, \mu^-, \tau^-\}; N = \{p, n\}; m = \{\pi, \eta, K\}; A :$ nucleus; X : hadrons

Neutral current: Z^0 exchange

• Elastic scattering
(target change but not break up):
• NC:
$$\overline{\nu}_l + N \rightarrow \overline{\nu}_l + N'$$

• Resonant meson production
(target is exited):
• NCRES:
 $\overline{\nu}_l + N \rightarrow \overline{\nu}_l + \Delta(1232) \rightarrow \overline{\nu}_l + N' + m$
• Coherent pion production
(target unchange):
• NCCOH: $\overline{\nu}_l + A \rightarrow \overline{\nu}_l + A + \pi^0$
• Deep inelastic scattering
(nucleon broken):
• NCDIS: $\overline{\nu}_l + N \rightarrow \overline{\nu}_l + X$







- Neutrino interactions at 0.1-20GeV energy:
- Quasi-elastic scattering (QE): scattering variables

• CCQE:
$$\nu_l + n \rightarrow l^- + p$$

- Mandelstam variables in lab frame (4-r $s = (p_{\nu} + p_{n})^{2} = (p_{l} + p_{p})^{2} = (p_{\nu} + p_{p})^{2}$ $t = (p_{\nu} - p_l)^2 = (p_n - p_p)^2 = (p_{\nu} - p_l)^2$ $u = (p_{\nu} - p_{p})^{2} = (p_{l} - p_{n})^{2} = (p_{\nu} - p_{p})^{2}$ $s + t + u = m_{\nu}^2 + m_n^2 + m_l^2 + m_n^2$
- In detail: $s = m_N^2 + 2m_N E_{\nu}$ $t = q^2 = -Q^2 = m_l^2 - 2E_{\nu}(E_l - p_l \cos \theta)$ $u = m_N^2 - 2m_N E_\nu + 2E_\nu (E_l - p_l \cos \theta)$

momentum):

$$p_n(p_l + p_p)$$

 $p_l(p_p - p_n)$
 $p_p(p_l - p_n)$

• Inelasticity:

$$y = \frac{p_l \cdot q}{p_l \cdot p_{\nu}}$$

 $\nu(E_{\nu}, \overline{p}_{\nu})$

• Hadronic invariant mass:
$$W = \sqrt{E_X^2 - p_X^2}$$

Π

4-momentum transferred: $Q^2 = -q^2 = 2E_{\nu}(E_l - p_l \cos \theta) - m_l^2$





- Neutrino interactions at 0.1-20GeV energy: **Quasi-elastic scattering (QE): (Llewellyn-Smith, 1972)**

• CCQE:
$$\nu_l + n \rightarrow l^- + p$$

Neutrino reconstructed energy (HW: derive it)

$$E_{\nu} = \frac{m_p^2 - m_l^2 - (m_n - E_b)^2}{2[(m_n - E_b) - E_l]}$$

 p_l, E_l, θ : out-going lepton momentum, energy, scattered angle $E_b = 27 MeV$: neutron binding energy in Oxygen m_p, m_n, m_l : proton, neutron, lepton masses

 $E_{\nu} = \frac{2m_n E_l - m_l^2}{2(m_n - E_l + p_l \cos \theta)}$ • Remove E_b (free neutron) and consider $m_p \approx m_n$:

 $+ 2(m_n - E_b)E_l$

 $\nu(E_{\nu}, \overline{p}_{\nu})$

- $+ p_l \cos \theta$



 $l(E_l, p)$



- Neutrino interactions at 0.1-20GeV energy:
- Quasi-elastic scattering (QE): (Llewellyn-Smith, 1972)

$$\frac{d\sigma}{dQ^2} = \frac{G_F^2 m_N^2 |V_{ud}|^2}{8\pi (P_\nu \cdot p_{N_i})_\nu^2} \left[A \pm \frac{(s-u)}{m_N^2} B + \frac{(s-u)^2}{m_N^4} C \right]$$

$$A = \frac{(m_l^2 + Q^2)}{m_N^2} \left[(1 + \eta) F_A^2 - (1 - \eta) F_1^2 + \eta (1 - \eta) F_2^2 + 4\eta F_1 F_2 - \frac{m_l^2}{4m_N^2} \left((F_1 + F_2)^2 + (F_A + 2F_P)^2 - \left(\frac{Q^2}{m_N^2} + 4\right) F_P^2 \right) \right]$$

$$B = \frac{Q^2}{m_N^2} F_A(F_1 + F_2) \qquad C = \frac{1}{4} (F_A^2 + F_1^2 + \eta F_2^2)$$

rgy: n-Smith. 19

- (-) +: (anti-) neutrino scattering
- G_F : Fermi constant
- Q^2 : squared 4-momentum transfered ($Q^2 = -q^2 > 0$)
- m_N : nucleon mass
- m_l : lepton mass
- E_{ν} : neutrino energy

•
$$(s-u) = 4m_N E_\nu - Q^2 - Q$$

- F_1, F_2, F_A, F_P : Dirac, Pauli, axial-vector, and preudoscalar form factors
- $\eta = Q^2 / 4m_N^2$



• Neutrino interactions at 0.1-20GeV energy: • Quasi-elastic scattering (QE): (Llewellyn-Smith, 1972) • Axial-vector form factor $F_A(Q^2) = \frac{g_A}{\left(1 + Q^2/M_A^2\right)^2}$ $g_A = F_A(0) = 1.2694 \pm 0.0028$ • Axial mass $M_A = 1.026 \pm 0.021 \ GeV$ $F_1^p(0) = 1, \qquad \qquad F_1^n(0) = 0,$ • Dirac form factor $F_1(Q^2) = F_1^p(Q^2) + F_1^n(Q^2)$ $F_2^p(0) = rac{\mu_p}{\mu_{
m N}} - 1, \qquad F_2^n(0) = rac{\mu_n}{\mu_{
m N}},$ • Pauli form factor $F_2(Q^2) = F_2^p(Q^2) + F_2^n(Q^2)$ • Nuclear magneton $\mu_N = e\hbar/2m_p$

- For $\nu_e, \overline{\nu}_e, \nu_\mu, \overline{\nu}_\mu$ scattering, the term related to A can be neglected. In this case, cross sections do not depend on the preudoscalar form factor F_{P}

- Proton magnetic moment μ_p
- Neutron magnetic moment μ_n



- Neutrino interactions at 0.1-20GeV energy:
- Quasi-elastic scattering (QE): measurement



- IHEP-ITEP, SJNP 30, 527 (1979)

- Neutrino interactions at 0.1-20GeV energy:
- form factors

$$\frac{d\sigma}{dQ^2} = \frac{G_F^2 m_N^2}{8\pi (P_\nu \cdot p_{N_i})_\nu^2} \left[A \pm \frac{(s-u)}{m_N^2} B + \frac{(s-u)}{m_N^4} \right]^2$$
$$F_1(Q^2) = \left(\frac{1}{2} - \sin^2 \theta_W\right) \left[\frac{\tau_3(1+\eta(1+\mu_p-\mu_p))}{(1+\eta)(1+Q^2/M_V^2)} \right]^2$$
$$F_2(Q^2) = \left(\frac{1}{2} - \sin^2 \theta_W\right) \frac{\tau_3(\mu_p-\mu_n)}{(1+\eta)\left(1+\frac{Q^2}{M_V^2}\right)^2} - \sin^2 \theta_W$$

- τ₃ = +1 (-1) for proton (neutron) scattering
 F^s_{1,2}(Q²): strange vector form factors
- Strange axial-vector form factor (Δs : stragne quark contribution to nucleon spin)

• NC elastic scattering: Similar to formulae of CCQE but without $|V_{ud}|^2$, and different





- Neutrino interactions at 0.1-20GeV energy:
- Resonant meson production (CC1 π dominates at low energy)

$$\begin{array}{ll}
\nu_{\mu} \, p \to \mu^{-} \, p \, \pi^{+}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, p \, \pi^{-} \\
\nu_{\mu} \, n \to \mu^{-} \, p \, \pi^{0}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, n \, \pi^{0} \\
\nu_{\mu} \, n \to \mu^{-} \, n \, \pi^{+}, & \overline{\nu}_{\mu} \, n \to \mu^{+} \, n \, \pi^{-} \\
\end{array}$$

$$\begin{aligned} \frac{d\sigma(\nu N \to lN\pi)}{dk^2 dW d\Omega_{\pi}} &= \frac{G_F^2}{2} \frac{1}{(2\pi)^4} \frac{|\mathbf{q}|}{4} \frac{-k^2}{(k^L)^2} \sum_{\lambda_2,\lambda_1} \left\{ \frac{|\mathbf{q}|}{k^L dW d\Omega_{\pi}} + \left[C_{L_-} (\tilde{F}_{\lambda_2\lambda_1}^{e_L}(\theta,\phi) - \tilde{G}_{\lambda_2\lambda_1}^{e_L}(\theta,\phi)) + C_{R_-} (\tilde{F}_{\lambda_2\lambda_1}^{e_R}(\theta,\phi) - \tilde{G}_{\lambda_2\lambda_1}^{e_R}(\theta,\phi)) + C_- (\tilde{F}_{\lambda_2\lambda_1}^{e_-}(\theta,\phi) - \tilde{G}_{\lambda_2\lambda_1}^{e_-}(\theta,\phi)) + C_+ (\tilde{F}_{\lambda_2\lambda_1}^{e_-}(\theta,\phi) - \tilde{G}_{\lambda_2\lambda_1}^{e_-}(\theta,\phi)) + C_+ (\tilde{F}_{\lambda_2\lambda_1}^{e_+}(\theta,\phi) - \tilde{G}_{\lambda_2\lambda_1}^{e_+}(\theta,\phi)) + C_+ (\tilde{F}_{\lambda_2\lambda_1}^{e_+}(\theta,\phi)) + C_+ (\tilde{F$$

• For anti-neutrino, swap C_{L_+} and C_{R_+}

_	$ u_{\mu} p ightarrow u_{\mu} p \pi^0,$	$\overline{\nu}_{\mu} p o \overline{\nu}_{\mu} p \pi^0$
)	$ u_{\mu} p \rightarrow \nu_{\mu} n \pi^+, $	$\overline{ u}_{\mu} n ightarrow \overline{ u}_{\mu} n \pi^0$
_	$ u_{\mu} n ightarrow u_{\mu} n \pi^0,$	$\overline{ u}_{\mu} n ightarrow \overline{ u}_{\mu} n \pi^0$
	$ u_{\mu} n ightarrow u_{\mu} p \pi^{-},$	$\overline{\nu}_{\mu} n o \overline{\nu}_{\mu} p \pi^-$



- Neutrino interactions at 0.1-20GeV energy:
- Resonant meson production (CC1 π dominates at low energy)

$$\begin{array}{ll}
\nu_{\mu} \, p \to \mu^{-} \, p \, \pi^{+}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, p \, \pi \\
\nu_{\mu} \, n \to \mu^{-} \, p \, \pi^{0}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, n \, \pi \\
\nu_{\mu} \, n \to \mu^{-} \, n \, \pi^{+}, & \overline{\nu}_{\mu} \, n \to \mu^{+} \, n \, \pi
\end{array}$$

$$\begin{array}{ll}
\nu_{\mu} \, p \to \nu_{\mu} \, p \, \pi^{0}, & \overline{\nu}_{\mu} \, p \to \overline{\nu}_{\mu} \, p \, \pi \\
\nu_{\mu} \, p \to \nu_{\mu} \, n \, \pi^{+}, & \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, n \, \pi \\
\nu_{\mu} \, n \to \nu_{\mu} \, n \, \pi^{0}, & \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, n \, \pi \\
\nu_{\mu} \, n \to \nu_{\mu} \, p \, \pi^{-}, & \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, p \, \pi
\end{array}$$


- Neutrino interactions at 0.1-20GeV energy:
- Resonant meson production (CC1 π dominates at low energy)

$$\begin{array}{ll}
\nu_{\mu} \, p \to \mu^{-} \, p \, \pi^{+}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, p \, \pi^{-} \\
\nu_{\mu} \, n \to \mu^{-} \, p \, \pi^{0}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, n \, \pi^{0} \\
\nu_{\mu} \, n \to \mu^{-} \, n \, \pi^{+}, & \overline{\nu}_{\mu} \, n \to \mu^{+} \, n \, \pi^{-} \\
\end{array}$$

$$\begin{split} \nu_{\mu} \, p &\to \nu_{\mu} \, p \, \pi^{0}, \qquad \overline{\nu}_{\mu} \, p \to \overline{\nu}_{\mu} \, p \, \pi^{0} \\ \nu_{\mu} \, p &\to \nu_{\mu} \, n \, \pi^{+}, \qquad \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, n \, \pi^{0} \\ \nu_{\mu} \, n \to \nu_{\mu} \, n \, \pi^{0}, \qquad \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, n \, \pi^{0} \\ \nu_{\mu} \, n \to \nu_{\mu} \, p \, \pi^{-}, \qquad \overline{\nu}_{\mu} \, n \to \overline{\nu}_{\mu} \, p \, \pi^{-} \end{split}$$

0

 Main background 	s of $CC0\pi$ at T2K
can mimic μ/e sig	ynals
$\pi^{+} \rightarrow \mu^{+} + \nu_{\mu}$ $\pi^{-} \rightarrow \mu^{-} + \overline{\nu}_{\mu}$	<i>BR</i> = 0.999877
$\pi^0 o 2\gamma$	BR = 0.988

 π can be absorbed into the nucleus: the event is indistinguisable with a QE event (similar observed final particles).







- Neutrino interactions at 0.1-20GeV energy:
- Resonant meson production (CC1 π dominates at low energy)



- Neutrino interactions at 0.1-20GeV energy:
- Resonant meson production (CC1 π dominates at low energy)



t Target	NC/CC Ratio	Value	Reference
H_2	$\sigma(u_\mup ightarrow u_\mup\pi^0)/\sigma(u_\mu p ightarrow\mu^-p\pi^+)$	0.51 ± 0.25	(Barish et al., 197
H_2	$\sigma(u_\mup o u_\mup\pi^0)/\sigma(u_\mu p o \mu^-p\pi^+)$	$0.09\pm0.05^*$	(Derrick et al., 19
H_2	$\sigma(u_\mup o u_\mun\pi^+)/\sigma(u_\mu p o \mu^-p\pi^+)$	0.17 ± 0.08	(Barish et al., 197
H_2	$\sigma(u_\mup o u_\mun\pi^+)/\sigma(u_\mu p o \mu^-p\pi^+)$	0.12 ± 0.04	(Derrick et al., 19
D_2	$\sigma(u_\mun o u_\mup\pi^-)/\sigma(u_\mu n o \mu^-n\pi^+)$	0.38 ± 0.11	(Fogli and Nardul
$C_3H_8 \ CF_3Br$	$\sigma(u_{\mu} N ightarrow u_{\mu} N \pi^{0})/2 \sigma(u_{\mu} n ightarrow \mu^{-} p \pi^{0})$	0.45 ± 0.08	(Krenz et al., 197
Al	$\sigma(u_\muN o u_\muN\pi^0)/2\sigma(u_\mu n o \mu^-p\pi^0)$	0.40 ± 0.06	(Fogli and Nardul
Al	$\sigma(u_\muN o u_\muN\pi^0)/2\sigma(u_\mu n o \mu^-p\pi^0)$	0.17 ± 0.04	(Lee et al., 1977)
Al	$\sigma(u_\muN o u_\muN\pi^0)/2\sigma(u_\mu n o \mu^-p\pi^0)$	$0.25 \pm 0.09^{**}$	(Nienaber, 1988)
D_2	$\sigma(u_\mun o u_\mup\pi^-)/\sigma(u_\mu p o \mu^-p\pi^+)$	0.11 ± 0.022	(Derrick et al., 19

 The cross sections of the other channels can be deduced from the above three ones by experimental measurements







• Neutrino interactions at 0.1-20GeV energy: Coherent pion production

$$\nu_{\mu} A \to \nu_{\mu} A \pi^{0} \quad \overline{\nu}_{\mu} A \to \overline{\nu}_{\mu} A \pi^{0}$$
$$\nu_{\mu} A \to \mu^{-} A \pi^{+} \quad \overline{\nu}_{\mu} A \to \mu^{+} A \pi^{0}$$



- $y = E_{\pi}/E_{\nu}$
- f_{π} : pion decay constant
- $Q^2 = -q^2$: 4-momentum transferred to hadronic system

- Negligible transferred energy
- Coherently produce pion
- No nuclear recoil

Adler's PCAC theorem

• |t| : magnitude of the square of the 4-momentum transfer to the nucleus



• Neutrino interactions at 0.1-20GeV energy: Coherent pion production



• CC COH Xsec measurement on C at T2K for ν_{μ} (left) and $\overline{\nu}_{\mu}$ (right)

- Negligible transferred energy
- Coherently produce pion
- No nuclear recoil





- Neutrino interactions at 20 500 GeV energy:
- **Deep inelastic scattering (CC)** $\nu_l + N \to l^- + X$

$$\frac{d^2 \sigma^{\nu,\overline{\nu}}}{dxdy} = \frac{G_F^2 m_N E_{\nu}}{\pi (1 + Q^2/m_{W,Z}^2)^2} \left[xy^2 F_1(x, Q^2) \right]$$

$$x = \frac{Q^2}{2p_N \cdot q}$$

$$y = \frac{p_N \cdot q}{p_N \cdot p_\nu}$$

$$xy = \frac{Q^2}{s - m_N^2}$$

$$s = (p_\nu + p_N)^2 = m_N^2 + 2p_\nu \cdot p_N$$

$$Q^2 = -m_l^2 + 2E_\nu(E_l - p_l \cos \theta)$$

 $\overline{\nu}_l + N \to l^+ + X$

(-)+ for (anti-)neutrino scattering

 $(1 - y) F_2(x, Q^2) \pm y (1 - y/2) xF_3(x, Q^2)$





Table 52.1: List of beam properties, targets, and run durations for modern accelerator-based neutrino experiments studying neutrino scattering.

PDG2023		$\langle E_{\mu} \rangle, \langle E_{\overline{\mu}} \rangle$	neutrino	run
Experiment	beam	(GeV)	target(s)	period
ArgoNeuT	$ u,\overline{ u} $	4.3, 3.6	Ar	2009 - 2010
$\mathrm{FASER}\nu$	$ u,\overline{ u}$	> 200	W, emulsion	2021 -
ICARUS (at CNGS)	ν	20	\mathbf{Ar}	2010 - 2012
ICARUS (at FNAL)	ν	0.8	\mathbf{Ar}	2021 -
K2K	ν	1.3	CH, H_2O	2003 - 2004
MicroBooNE	ν	0.8	Ar	2015 - 2020
MINERvA	$ u,\overline{ u}$	3.5 (LE),	He, C, CH,	2009 - 2019
		5.5 (ME)	H_2O , Fe, Pb	
MiniBooNE	$ u,\overline{ u}$	0.8, 0.7	CH_2	2002 - 2019
MINOS	$ u,\overline{ u}$	3.5, 6.1	\mathbf{Fe}	2004 - 2016
NINJA	$ u,\overline{ u}$	0.6, 0.6	Fe, emulsion	2015 -
NOMAD	$ u,\overline{ u}$	23.4, 19.7	\mathbf{C}	1995 - 1998
NOvA	$ u,\overline{ u}$	2.0, 2.0	CH_2	2010 -
SBND	ν	0.8	\mathbf{Ar}	2024 -
SciBooNE	$ u,\overline{ u}$	0.8, 0.7	\mathbf{CH}	2007-2008
T2K	$ u,\overline{ u}$	0.6, 0.6	$\mathrm{CH},\mathrm{H}_{2}\mathrm{O},\mathrm{Fe}$	2010 -

How do experiments measure Xsec?

- Recall that the true number of interactions at the detector: $N = \phi_a N_b \sigma$ or in general: $N(E_\nu) = \phi(E_\nu) \cdot T \cdot \sigma(E_\nu)$ Consider an observable x that describes the interaction

$$N(x_{true}) = \int \frac{d\sigma(E_{\nu}, x_{true})}{dx_{true}} \phi(E_{\nu})$$

- There is no perfect detector, what you measured (reconstrued) is smeared from the true information $N(x_{recon}) \neq N(x_{true})$. You need to account for this difference and make your measurement closer to the true information as much as possible:
- Take into account the detector efficiency *E*
 - Construct a smearing matrix (migration matrix) U_{smearing} to migrate from true quantity to reconstructed quantity $N(x_{recon}) = \int \frac{U_{smearing}}{d\sigma(E_{\nu}, x_{true})} \frac{d\sigma(E_{\nu}, x_{true})}{dx_{true}}$

). $T. dx_{true}$

$$\frac{de}{de}\phi(E_{\nu}) \cdot T \cdot \epsilon \cdot dx_{true}$$



- Solving for differential cross section: $\frac{d\sigma(E_{\nu}, x_{true})}{dx_{true}} = \frac{U_{smearing}^{-1} N(x_{recon})}{\phi(E_{\nu}) \cdot T \cdot \epsilon \cdot dx_{true}} \equiv \frac{U_{unfolding} N(x_{recon})}{\phi(E_{\nu}) \cdot T \cdot \epsilon \cdot dx_{true}}$
 - Where $U_{unfolding} = U_{smearing}^{-1}$
- In general

 $\left(\frac{d\sigma}{dx}\right)_{\alpha} = \frac{\sum_{i} U_{i\alpha} (N_{data,i}^{selected} - N_{data,i}^{bckg})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$

- $U_{i\alpha}$: unfolding matrix (unsmearing matrix)
- N^{selected} : total number of selected events after applying all cuts
- N_{data}^{bckg} : number of backgrounds in selected events
- ϕ : neutrino integrated flux
- T: number of target nucleons
- ϵ : detector efficiency
- Δx : bin width



An example of ν_{μ} CC0 π Xsec measurement at T2K

 $\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - N_{data,i}^{bckg})$ $\left(\frac{d\sigma}{dx}\right)_{\alpha} = \frac{\sum_{i} U_{i\alpha}(N_{data,i}) - N_{data,i}}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$

• Analysis flow:







An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut





An example of ν_{μ} CC0 π Xsec measurement at T2K

- **o Event selection:**
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut

$$\left(\frac{d\sigma}{dx}\right)_{\alpha} = \frac{\sum_{i} U_{i\alpha} (N_{data,i}^{selected} - \phi \cdot T \cdot \epsilon \cdot \Delta x)}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$$

 $\nu_{\mu} + A \rightarrow \mu^{-} + X + 0\pi^{\pm}$



Figure 1: Signal topology for $CC0\pi^{\pm}$ sample



An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
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 - Michel electron tagging
 - Contained volume cut

$$\left(\frac{d\sigma}{dx}\right)_{\alpha} = \frac{\sum_{i} U_{i\alpha} (N_{data,i}^{selected} - I_{\alpha})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$$

 Whole fundamental cuts related to 2D and 3D reconstruction including track matching between detectors



An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut

 $\frac{\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - N_{data,i}^{bckg})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$ $d\sigma$



Figure 9: X vertex distribution for H₂O target

What is the fiducial volume?



An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut

Figure 13: the number of track distribution broken down by interaction topologies for the charged current inclusive sample on water target





 $CC0\pi$ typically has 1 or 2 tracks, events with large number of tracks are more likely the backgrounds





An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut

 $= \frac{\sum_{i} U_{i\alpha} (N_{data,i}^{selected} - N_{data,i}^{bckg})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$ $d\sigma$





An example of ν_{μ} CC0 π Xsec measurement at T2K

• Event selection:

•	Define signal sample ($CC0\pi$)	ack	300
•	Define selections (cuts)	of tra	250
	 Pre-selection 	#	
	 Fiducial volume cut 		200
	 Number-of-track cut 		150
	 Muon confidence level cut Particle charge cut 		100
			100
	 Track/cluster ratio cut 		50
	 Michel electron tagging 		
	 Contained volume cut 		





An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut





 Remove backgrounds from interactions from experimental hall or other subdectors

An example of ν_{μ} CC0 π Xsec measurement at T2K

- Event selection:
- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut



 $= \frac{\sum_{i} U_{i\alpha} (N_{data,i}^{selected} - N_{data,i}^{bckg})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$



 Remove main backgrounds from Wall **MRDs**





7 CM

An example of ν_{μ} CC0 π Xsec measurement at T2K

• Event selection:

- Define signal sample ($CC0\pi$)
- Define selections (cuts)
 - Pre-selection
 - Fiducial volume cut
 - Number-of-track cut
 - Muon confidence level cut
 - Particle charge cut
 - Track/cluster ratio cut
 - Michel electron tagging
 - Contained volume cut



• An event which is **not actually a true signal** event can pass your selection even after appying all cuts, it is called a background event



(a) all selection applied



An example of ν_{μ} CC0 π Xsec measurement at T2K

o Unfolding:

- Our reconstruction is not perfect, an event can be reconstructed into a wrong bin
- $U_{i\alpha}$ presents the probability that an event observed in bin *i* actually happened in bin α
- We can use our Monte Carlo to construct a smearing matrix (migration matrix) which indicates the fraction of events generated in bin α be mis-reconstructed into bin i
- If the detector has good resolution, the matrix should be close to diagonal
- Bin width is also important. If the bins are too small compared to our resolution, the matrix may not be diagonal

 $= \frac{\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - N_{data,i}^{bckg})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$ $\frac{d\sigma}{dx}$





An example of ν_{μ} CC0 π Xsec measurement at T2K

Bins • Unfolding: 10 • Unfolding matrix True (unsmearing matrix) can be obtained by inverting the smearing matrix (migration matrix

$$U_{unfolding} = U_{smearing}^{-1}$$

 $\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - N_{data,i}^{bckg})$ $d\sigma$ $\phi \cdot T \cdot \epsilon \cdot \Delta x$







An example of ν_{μ} CC0 π Xsec measurement at T2K

o Flux:

Integrated flux over all neutrino energies

$$\phi = \frac{\int dE_{\nu}\phi(E_{\nu})}{dE_{\nu}}$$



 $\frac{\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - 1)}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$ $\left(\frac{d\sigma}{dx}\right)$





An example of ν_{μ} CC0 π Xsec measurement at T2K

- **o T: number of nucleons in target** • $T = \frac{M}{m} \cdot N_A * (n_p + n_n)$
- M: target mass in grams
- m: atomic mass
- n_p, n_n : number of protons and neutrons per atom
- $N_A = 6.02214 \times 10^{23} mol^{-1}$: Avogadro's number

$$\left(\frac{d\sigma}{dx}\right)_{\alpha} = \frac{\sum_{i} U_{i\alpha}(N_{data,i}^{selected} - N_{data,i})}{\phi \cdot T \cdot \epsilon \cdot \Delta x}$$

- **o** How many nucleons are there in Super-K detector?
- $M = 50kt = 5 \times 10^{10} g$
- $m_{H_2O} = 18$
- $(n_p + n_n)_{H_2O} = (10 + 8) = 18$ $T = \frac{5 \times 10^{10}}{18} * 6.0214 \times 10^{23} * (10 + 8) = 3 \times 10^{34}$







An example of ν_{μ} CC0 π Xsec measurement at T2K

• $\epsilon = \frac{\text{number of signal events after selection}}{\text{number of signal events in MC}}$



(a) with respect to μ angle

Figure 37: Selection efficiency in one-dimensional phase space for H_2O target

 $\sum_{i} U_{i\alpha}(N_{data,i}^{selected} \cdot$ $d\sigma$

(b) with respect to μ momentum



An example of ν_{μ} CC0 π Xsec measurement at T2K

Bin width Δx

- Bin width is defined based on the number of expected events in each bin
- More bins (smaller bin width) you have, better you can distinguish the features of the distribution
- However, more bins you have, less events in each bin you get (worse statistics)







Contents

- 1. A brief history and properties of neutrinos
- 2. Cross sections
- 3. Neutrino interactions
 - 3.1: Neutrino electron scattering 3.2: Neutrino - nucleon scattering
 - 3.3: Cross section measurement

3.4: Summary

4. Introduction to NEUT - an event generator 4.1: Neutrino event generator 4.2: Introduction to NEUT 4.3: Practice with NEUT 4.4: Exercises

3.4: Summary

- Three active neutrinos (ν_e, ν_μ, ν_τ) an exchange W^{\pm}, Z^0 via CC or NC
- Active neutrinos only participate in weak interaction with extremely small Xsec
- A neutrino can elastic, quasi-elastic, or inelastic scattering with electron, nucleus, nucleon, or quarks depending on its energy
- Theoretical calculation of Xsec is difficult when dealing with nuclear effects
- Experimental measurement of Xsec is difficult due to tiny Xsec and technical barriers
- Interactions in medium energy region (~GeV) is important because they cover mostly accessible and abundant neutrino sources (solar, reactor, atmospheric, and accelerator)

- Three active neutrinos (ν_e, ν_μ, ν_τ) and their anti-particles interact with matter by

Contents

- 1. A brief history and properties of neutrinos
- 2. Cross sections
- 3. Neutrino interactions
 - 3.1: Neutrino electron scattering
 - 3.2: Neutrino nucleon scattering
 - 3.3: Cross section measurement
 - 3.4: Summary
- 4. Introduction to NEUT an event generator
 - 4.1: Neutrino event generator
 - 4.2: Introduction to NEUT
 - 4.3: Practice with NEUT
 - 4.4: Exercises



4.1: Neutrino event generators

- Neutrino event generators are MC softwares used for simulating neutrino interactions
- They are bridges connecting between theory and experiment
- They play an important role in oscillation analysis:
 - Predicting neutrino flux
 - Simulate interaction channels
 - Simulate signals and bacgrounds
 - Calculate efficiencies and systematics uncertainties



4.1: Neutrino event generators

• How do they work?

- Neutrino type
- Target

...

Nu energy or flux

Inputs

- Interaction mode
- Other parameters



From H. Van



4.1: Neutrino event generators

• GENIE:

- Developed by international colaboration for neutrino interactions from MeV to PeV ($10^{15} eV$)
- ROOT based neutrino generator written in C++.
- Well maintainance and open source
- Recently implemented models of quasielastic and 2p2h interactions
- Used in many neutrino oscillation experiments such as T2K, NOvA, ...

4.1: Neutrino event generators

- GiBUU (Giessen-Boltzmann-Uehling-Uhlenbeck Project):
 - neutrino and hadron with nuclei)
 - complete final state of an event
 - Not only used as neutrino generator but also in nuclear reactions
 - Written in Fortran and open source

 Provides a unified theory and transport framework in the MeV and GeV energy regimes for elementary reactions on nuclei (electron, photon,

Provides a full dynamical description of the reaction and delivers the

4.1: Neutrino event generators

• NuWro:

- Has been developed at the University of Wroclaw since 2006
- Written in C++ \bullet
- target from 100MeV to TeV

 Simulate neutrino interactions taking into account beam profile and composition, detailed detector geometry as well as FSI in the nuclear

4.2: Introduction to NEUT

• NEUT:

- Originally developed for Kamioka experiment to predict neutrino-induced backgrounds for nucleon decay search
- Has been officially used for Super-K and T2K (currently NEUT 5.4.0)
- Mainly written in Fortran, not open source yet
- Primarilly simulate neutrino-nucleus/nucleon interactions in wide range of energy from 100 MeV to few TeV ($10^{12} eV$)
- Incorporates nuclear effects
- Recently impemented state-of-the-art CCQE and multi-nucleon model, single pion production models, and electron - nucleus scattering

4.2: Introduction to NEUT

- NEUT interaction models (used at T2K analysis):
 - 1p1h (CCQE)
 - Produce single-ring electron-like and muon-like events
 - Contribute 70% of the $1R_{\mu}$ selection at far detector
 - Use Benhar Spectral Function (Nucl. Phys. A 579, 493–517 (1994))
 - Nucleon vector form factors: BBBA05 description (Nucl. Phys. B Proc. <u>Suppl. 159, 127–132 (2006)</u>
 - Nucleon axial mass $M_{\Lambda}^{QE} = 1.03 \pm 0.06 \ GeV$
4.2: Introduction to NEUT

- NEUT interaction models (used at T2K analysis):
 - CC 2p2h (2-particles 2-holes), NC 2p2h not included yet
 - Neutrinos interactions with 2 nucleons producing 2 holes
 - Produce single-ring electron-like and muon-like events
 - Contribute 12% of the $1R_{\mu}$ selection at far detector
 - Use Nieves model (Phys. Lett. B 707, 72–75 (2012))



4.2: Introduction to NEUT

- NEUT interaction models (used at T2K analysis):
 - Single pion production
 - Coherent 1π (Berger–Sehgal model)
 - invariant mass $1.3 \leq W \leq 2 \ GeV$
 - Contribute 13% of the $1R_{\mu}$ selection at far detector

• Resonant and non-resonant 1π (Rein–Sehgal (RS) model) with hadronic

4.2: Introduction to NEUT

- NEUT interaction models (used at T2K analysis):
 - Deep inelastic scattering

 - Use <u>GRV98</u> Parton Distribution Functions (PDFs) • Bodek–Yang (BY) correction is used for $Q^2 \le 1.5 \ GeV^2$
 - Processes begin where hadronic invariant mass W > 1.3 GeV
 - <u>Pythia 5.72 for</u> W > 2 GeV
 - <u>A custom model interpolating between the $\Delta(1232)$ and DIS interactions</u> is employed for $W < 2 \ GeV$

4.3: Practice with NEUT

- Working space is IOP cluster
- Source code: /home/vson/neut_5.4.0

```
[[vson@tcp ~]$ pwd
/home/vson
[[vson@tcp ~]$ ls
neut_5.4.0
neut_5.4.0/src/neutsmpl/Cards/neut_5.4.0_nd5_C_ccqe.card
```

• CERNLIB and ROOT are needed for ROOT installtion (already available on IOP cluster)

• Car file: /home/vson/neut_5.4.0/src/neutsmpl/Cards/neut_5.4.0_nd5_C_ccqe.card

[[vson@tcp ~]\$ ls neut_5.4.0/src/neutsmpl/Cards/neut_5.4.0_nd5_C_ccqe.card



4.3: Practice with NEUT

- Working space is IOP cluster
 - o NEUT input: Card file
 - Used to specify models, parameters, interaction modes, …
 - Card file: /home/vson/neut_5.4.0/src/neutsmpl/Cards/neut_5.4.0_nd5_C_ccqe.card
 - When setting your card file, check carefully there is no letter C at the beginning of the line (C: comment out that line)

```
C-----
C
C EVCT-NEVT
C Number of events to generate
C
C EVCT-NEVT 1000000
EVCT-NEVT 1000
C-----
```



4.3: Practice with NEUT

- Working space is IOP cluster
- 1. Log in IOP cluster (2024) \$ ssh -XY vson@202.151.162.171 pass: son...
- **2. Create a new folder with your name**
- \$ cp -r -p neut_5.4.0/ tvngoc/
- 4. Go to your working directory \$ cd tvngoc/neut_5.4.0/ \$ ls

include lib setup_env_neut540.sh src

\$ mkdir your_name (in the following steps, replace tvngoc by your_name) 3. Copy source folder to your working directory you have just created

4.3: Practice with NEUT

- Working space is IOP cluster
- **5. Check where you are** \$ pwd

/home/vson/tvngoc/neut_5.4.0 6. Before working with neut, we need to setup environment. Use vim or gedit or other text editors to open and modify the file setup_env_neut540.sh Change the default directory to your directory, don't forget to replace tvngoc by your_name

\$ vi setup_env_neut540.sh

7. Press "i" to modify:

#export NEUT_ROOT=/home/vson/neut_5.4.0 export NEUT_ROOT=/home/vson/tvngoc/neut_5.4.0 (same as output when you pwd) Press "esc", then ":wq", press "enter" to save and close the file

4.3: Practice with NEUT

- Working space is IOP cluster
- **8. Source the environment file** \$ source setup_env_neut540.sh 9. Go to working directory \$ cd /src/neutsmpl **10. Clean and make to create binary files** \$ /bin/csh Cleanneutsmpl.csh (or ./Cleanneutsmpl) \$ /bin/csh Makeneutsmpl.csh (or ./Makeneutsmpl) 11. Press "Is" then "enter" to see if there are "neutroot2" binary file created **12. Open card file to see what inside** \$ vi Cards/neut_5.4.0_nd5_C_ccqe.card Press "i" to enter modified mode, press "esc" and then ":wq" to save and exit. The letter "C" at the beginning of the line means "comment out" that line.

4.3: Practice with NEUT

- Working space is IOP cluster
- the exercises.
- \$ vi make_histos_standalone_neut540_ccqe_simple.cc

13. Run neut with card file you have just edited. Remember to change the output filename from output_numu_offaxis_tvngoc.root to your_name.root \$./neutroot2 Cards/neut_5.4.0_nd5_C_ccqe.card output_numu_offaxis_tvngoc.root

Press "Is" to if there is a file "output_numu_offaxis_tvngoc.root" in the directory

14. Root macro file for analyzing the output. Edit this file if you want to do

4.3: Practice with NEUT

 Working space is IOP cluster 15. Run the file to create basic histograms. Remember to change tvngoc to

your_name

\$ root -b -q

ngoc.root", "output_numu_offaxis_tvngoc_basichisto.root")

16. Open created root file to check the histograms \$ root -l output_numu_offaxis_tvngoc_basichisto.root root [0] Attaching file output_numu_offaxis_tvngoc_basichisto.root as _file0... root [1] .ls root [2] NEUT_ppro ->Draw();

- 'make_histos_standalone_neut540_ccqe_simple.cc("output_numu_offaxis_tv

4.3: Practice with NEUT

• Summary

1. Define input parameters in Card file

- **Off-axis muon neutrino mode**
 - Flux: nd5_tuned13av1.1_13anom_run1-7c_fine.root Histogram: enu_nd5_tuned13a_numu
- Off-axis muon anti-neutrino mode
 - Flux: nd5_tuned13av1.1_13anom_run5c-7b_antinumode_fine.root Histogram: enu_nd5_tuned13a_numub
- On-axis muon neutrino mode
 - Flux: nd34_tuned_11bv3.1_250ka.root
 - Histogram: ing3_tune_numu
- On-axis muon anti-neutrino mode
 - Flux: run5c_tune_INGRID_13a_1_1.root
 - Histogram: ing3_tune_numub

Change neutrino flavor code accordingly

4. Introduction to NEUT - an event generator Working space is IOP cluster Name it yourself **2. Run NEUT to generate events** ./neutroot2 directory_to_cardfolder/cardfile.card output_neut.root **3. Analyze the output** Same root -b -q 'make_histos_standalone_neut540_ccqe_simple.cc("output_neut.root","

4.3: Practice with NEUT

- Summary
- name_your_root_file_basichisto.root")'
 - The macro is at neut_5.4.0/src/neutsmpl/

4.4: Exercises

- Good information
 - Particle code (ID)

if(nvect->PartInfo(i)->fPID == 13)
{muon_index=i; cout<<"This is a muon
with the index: "<<muon_index<<endl;}</pre>

Particle	Code	Particle	Code
Proton	2212	Xi-zero	3322
Antiproton	-2212	Antixi-zero	-3322
Electron	11	Negative Xi	3312
Positron	-11	Positive Xi	-3312
Electron Neutrino	12	Omega-minus	3334
Electron Antineutrino	-12	Antiomega	-3334
Photon	22	Positive Tau	-15
Neutron	2112	Negative Tau	15
Antineutron	-2112	Tau neutrino	16
Positive Muon	-13	Tau antineutrino	-16
Negative Muon	13	D-plus	411
Kaon-zero long	130	D-minus	-411
Positive Pion	211	D-zero	421
Negative Pion	-211	AntiD-zero	-421
Positive Kaon	321	D_s-plus	431
Negative Kaon	-321	D_s-minus	-431
Lambda	3122	Lambda_c-plus	4122
Antilambda	-3122	Xi_c-plus	4232
Kaon zero short	310	Xi_c-zero	4112
Negative Sigma	3112	Xi'_c-plus	4322
Positive Sigma	3222	Xi'_c-zero	4312
Sigma-zero	3212	Omega_c-zero	4332
Pion-zero	111	Antilambda_c-minus	-4122
Kaon-zero	311	AntiXi_c-minus	-4232
Antikaon-zero	-311	AntiXi_c-zero	-4132
Muon neutrino	14	AntiXi'_c-minus	-4322
Muon antineutrino	-14	AntiXi'_c-zero	-4312
Antisigma-minus	-3222	AntiOmega_c-zero	-4332
Antisigma-zero	-3212		
Antisigma-plus	-3112		

4.4: Exercises

Good information

- Kinematic variables
- Mandelstam variables in lab frame (4-momentum): $s = (p_{\nu} + p_{n})^{2} = (p_{l} + p_{p})^{2} = (p_{\nu} + p_{n})(p_{l} + p_{p})$ $t = (p_{\nu} - p_{l})^{2} = (p_{n} - p_{p})^{2} = (p_{\nu} - p_{l})(p_{p} - p_{n})$ $u = (p_{\nu} - p_{p})^{2} = (p_{l} - p_{n})^{2} = (p_{\nu} - p_{p})(p_{l} - p_{n})$ $s + t + u = m_{\nu}^{2} + m_{n}^{2} + m_{l}^{2} + m_{p}^{2}$
- In detail: $s = m_N^2 + 2m_N E_{\nu}$ $t = q^2 = -Q^2 = m_l^2 - 2E_{\nu}(E_l - p_l \cos \theta)$ $u = m_N^2 - 2m_N E_{\nu} + 2E_{\nu}(E_l - p_l \cos \theta)$



4.4: Exercises

Good information

Neutrino reconstructed energy

$$E_{\nu} = \frac{m_p^2 - m_l^2 - (m_n - E_b)^2 + 2[(m_n - E_b) - E_l + 2]}{2[(m_n - E_b) - E_l + 2]}$$

 p_l, E_l, θ : out-going lepton momentum, energy, scattered angle $E_{h} = 27 MeV$: neutron binding energy in Oxygen m_p, m_n, m_l : proton, neutron, lepton masses

• Remove E_b and consider $m_p \approx m_n$:



$$E_{\nu} = \frac{2m_n E_l - m_l^2}{2(m_n - E_l + p_l \cos \theta)}$$
(Erec2)

4.4: Exercises

• The following histograms have already created KEY: TH1DNEUT_mode;1 NEUT mode KEY: TH1DNEUT_enu;1 Neutrino energy KEY: TH1DNEUT_pmu;1 muon momentum KEY: TH1DNEUT_anglemu;1 muon angle KEY: TH1DNEUT_cosanglemu;1 cos muon angle KEY: TH2DNEUT_mu_pvscosangle;1 muon momentum vs. cos(#theta) KEY: TH1DNEUT_enurec;1 (equation (**Enurec2**) in slide 79) KEY: TH1Dflux_numu;1 #nu_{#mu} flux KEY: TH1Devtrt_numu;1 #nu_{#mu} event rate

4.4: Exercises

o Run NEUT with different inputs from card file

- Off-axis muon neutrino mode

 - Flux: nd5_tuned13av1.1_13anom_run1-7c_fine.root • Histogram: enu_nd5_tuned13a_numu
- Off-axis muon anti-neutrino mode

 - Flux: nd5_tuned13av1.1_13anom_run5c-7b_antinumode_fine.root • Histogram: enu_nd5_tuned13a_numub
- On-axis muon neutrino mode
 - Flux: nd34_tuned_11bv3.1_250ka.root
 - Histogram: ing3_tune_numu
- On-axis muon anti-neutrino mode
 - Flux: run5c_tune_INGRID_13a_1_1.root
 - Histogram: ing3_tune_numub
- Remember to change the output filename accordingly

4.4: Exercises

• Produce the following histograms

- Proton momentum
- Proton angle
- Cosine of proton angle

What is proton ID?

- 4 momentum transfer (equation (Q2rec) in slide <u>Kinematic variables</u>)
- Neutrino reconstructed energy (equation (Erec1) in slide <u>Nu energy</u>)
- Neutrino true energy (get from neut output), plot in the same canvas to compare with **Erec1** and **Erec2**

- if(nvect->PartInfo(i)->fPID == 2212) {proton_index=i; cout<<"This is a proton with the index: "<<pre>roton index<<endl;}</pre>



4.4: Exercises

- Hints 0
- **Define histogram** TH1D *NEUT_Q2 = new TH1D("NEUT_Q2", "Q2", 100,
- **Define what to fill in your histogram** (4 momentum calculation) double Q2 = 2^{*}e_nu*(e_mu-p_mu*cos_mu) - xmmu*xmmu;
- Fill histogram **NEUT_Q2**->Fill(Q2)
- Write histogram into root file

NEUT Q2->Write();







Abbreviation

- ANL: <u>Argonne National Laboratory</u>
- BNL: Brookhaven National Laboratory
- BEBC: Big European Bubble Chamber
- FNAL: Fermi National Accelerator Laboratory
- Adler's PCAC theorem: partially conserved axial vector current

Natural unit



	Table 2.1 Relationship betw	veen S.I. and natural units.	
Quantity	[kg, m, s]	[ħ, c, GeV]	$\hbar = c = 1$
Energy	$kg m^2 s^{-2}$	GeV	GeV
Momentum	$kg m s^{-1}$	GeV/c	GeV
Mass	kg	GeV/c^2	GeV
Time	S	$(\text{GeV}/\hbar)^{-1}$	GeV^{-1}
Length	m	$(\text{GeV}/\hbar c)^{-1}$	GeV^{-1}
Area	m^2	$(\text{GeV}/\hbar c)^{-2}$	GeV ⁻²

length must be multiplied by $\hbar c$, ...

$\hbar = c = 1.$

becomes
$$E^2 = p^2 + m^2$$
.

• To convert from natural unit to SI unit, the mass must be multiplied by c^2 , the

$\hbar c = 0.197 \, \text{GeV} \, \text{fm},$

where one femtometre (fm) = 10^{-15} m.

Michel electron

Experiments soon revealed the manner of muon's decay, into an electron and a pair of neutrinos, and the shape of the electron spectrum was first calculated by Louis Michel in 1950. This is where "Michel electron" comes from. Michel introduced a single parameter ρ to describe the shape. However,

after the non-conservation of parity in weak interactions proposal of Lee and Yang three more parameters (η , ξ , and δ) had to be added, and now all four are referred to

collectively as "Michel parameters".





4.2: Introduction to NEUT

- **NEUT limitations:**
 - Nuclear models are inconsistent between models or steps in the factorisation:

 - Benhar et al. SF can be used for CCQE but no other modes LFG used for FSI nuclear description
 - Benhar et al. SF Pauli blocking uses simple, RFG-like approach
 - Nuclear effects in single pion production are largely ignored
 - Nuclear transparency has no effect on inclusive cross-section

From L. Pickering

4.2: Introduction to NEUT

• Future plan:

Development has begun on NEUT6 - Targeted at HK and final T2K analyses: lacksquare

- Significant reorganization of code-base Ο
- Improved, modern build system Ο
- Removed dependence on an external CERNLIB2005 Ο
- New TOML-based configuration file Ο
- Modern C/Fortran interop Ο
- Automatic C/Fortran interface generation for model integration Ο
- \bullet 2023
 - Will also release the final NEUT5 series release as open source Ο
- \bullet release

Aim is to release NEUT6 as open source under the GPL before the end of

Hope to produce comprehensive data-model comparisons alongside NEUT6

From L. Pickering

Fiducial volume

- largely excluded
- physics analysis
- Reduce the backgrounds

Super-Kamiokande Fiducial Volume

Fiducial mass: 22 kt

• An inner volume of particle detector media in which background events are

The well understood region of the detector where the events are accepted for



Solid angle



	Solid Angle
φ	$d\Omega = \frac{dA_j}{r^2}$
- dA _j lθ d.	$dA_j = AB \ge BC$
	$dA_j = (r \sin\theta d\phi) x (r d\theta)$
t ₁ dφ	$\hat{\mathbf{f}}_1 d\Omega = \frac{r^2 \sin\theta d\theta d\phi}{r^2}$
	$d\Omega = sin\theta d\theta d\phi$

3.2 Neutrino - nucleon interactions

- Neutrino interactions at 0.1-20GeV energy:
- Multi-pion production





3.2 Neutrino - nucleon interactions

- Neutrino interactions at 500 GeV 1 EeV ($10^{18} eV$) energy:
- Ultra-high energy neutrinos

3.2 Neutrino - nucleon interactions

• Neutrino interactions at 0.1-20GeV energy: Kaon production

CC: $u_{\mu} \, n o \mu^- \, K^+ \, \Lambda^0$ $u_{\mu} \, p \to \mu^- \, K^+ \, p$ $u_{\mu} \, n
ightarrow \mu^{-} \, K^{0} \, p$ $\nu_{\mu} n \rightarrow \mu^{-} K^{+} n$ $\nu_{\mu} p \rightarrow \mu^{-} K^{+} \Sigma^{+}$ $u_{\mu} n \to \mu^- K^+ \Sigma^0$ $u_{\mu} \, n o \mu^{-} \, K^{0} \, \Sigma^{+}$

NC: $\nu_{\mu} \, p \to \nu_{\mu} \, K^+ \Lambda^0$ $u_{\mu} n \rightarrow \nu_{\mu} K^0 \Lambda^0$ $\nu_{\mu} p \rightarrow \nu_{\mu} K^+ \Sigma^0$ $u_{\mu} \, p \to \nu_{\mu} \, K^0 \, \Sigma^+$ $u_{\mu} n \to \nu_{\mu} K^0 \Sigma^0$ $\nu_{\mu} n \rightarrow \nu_{\mu} K^+ \Sigma^ \nu_{\mu} n \rightarrow \nu_{\mu} K^{-} \Sigma^{+}$

Potential background of proton decay: $p \rightarrow K^+ \nu$



Fermi's golden rule • Transition rate: $\Gamma_{fi} = 2\pi |T_{fi}|^2 \rho(E_i) = 2\pi \int |T_{fi}|^2 \delta(E_i - E) dn$

• T_{fi} : transition matrix element

•
$$\rho(E_i) = \left| \frac{dn}{dE} \right|_{E_i} = \int \frac{dn}{dE} \delta(E_i - E) dE$$
: density o
• $dn = (2\pi)^3 \frac{d^3 \overrightarrow{p_1}}{(2\pi)^3} \frac{d^3 \overrightarrow{p_2}}{(2\pi)^3} \delta^3 \left(\overrightarrow{p_a} + \overrightarrow{p_b} - \overrightarrow{p_1} - \overrightarrow{p_2} \right)$:
 $E \to E + dE$

• Lorentz invariant matrix element: $M_{fi} = \sqrt{2E_a \cdot 2E_b \cdot 2E_1 \cdot 2E_2} T_{fi}$

•
$$\Gamma_{fi} = \frac{(2\pi)^4}{4E_a E_b} \int |M_{fi}|^2 \delta(E_a + E_b - E_1 - E_2) \delta^3(\overline{p})$$

of states

number of accessible states in the energy range

 $\overrightarrow{p_a} + \overrightarrow{p_b} - \overrightarrow{p_1} - \overrightarrow{p_2}) \frac{d^3 \overrightarrow{p_1}}{(2\pi)^3 2E_1} \frac{d^3 \overrightarrow{p_2}}{(2\pi)^3 2E_2}$

3.1 Neutrino - electron interactions

• Neutrino - electron quasi-elastic scattering

$$\nu_{\mu} + e^- \rightarrow \nu_e + \mu^- \qquad E_{\nu_{\mu}} >$$

• Calculate threshold:

greater than $\left(\sum_{\mathbf{v}} m_{\mathbf{X}}\right)^2$ ~ 2

$$E_{\nu}^{th} = \frac{\left(\sum_{X} m_{X}\right)}{2m_{A}} - \frac{m_{A}}{2}$$

 $\nu + A \rightarrow \sum X$

> 10.92 *GeV*

Did you calculate this value as Sanjib's request?

• A at rest, neglecting neutrino mass, in center of mass frame, $s = 2E_{\nu}m_A + m_A^2$ must

$m_{\rho} = 0.511 \ MeV$ $m_{\mu} = 105.66 \, MeV$

3. Neutrino energies and interactions

- **o Neutrino interactions at 0.1-20GeV energy:**
- Quasi-elastic scattering (QE):
 - CCQE: $\nu_l + n \rightarrow l^- + p$
 - NC: $\nu_l + N \rightarrow \nu_l + N'$
- Resonant meson production (RES):
 - CC RES: $\nu_l + N \rightarrow l^- + N' + m$
 - NC RES: $\nu_l + N \rightarrow \nu_l + N' + m$
- Coherent pion production (COH):
 - CC COH: $\nu_l + A \rightarrow l^- + A + \pi^+$
 - NC COH: $\nu_l + A \rightarrow \nu_l + A + \pi^0$
- Deep inelastic scattering (DIS):
 - CC DIS: $\nu_l + N \rightarrow l^- + N' + hadrons$
 - NC DIS: $\nu_l + N \rightarrow \nu_l + N' + \text{hadrons}$

Where $l^- = \{e^-, \mu^-, \tau^-\}; N = \{p, n\};$ $m = \{\pi, \eta, K\}; A$: nucleus



- There are similar channels for antineutrino
- For RES at low E, $CC1\pi \& NC1\pi$ dominates
- We will focus on neutrino interactions at this energy range



Neutrino energies and interactions

