

# Present Status and Future Prospects of a Light eV-Scale Sterile Neutrino



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# Light Sterile Neutrinos: Huge Interest in the Community

## Light Sterile Neutrinos: A White Paper

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Sterile neutrinos: singlets of  $SU(2) \times U(1)$  gauge group, provide economical extension of the SM

Extensive study of sterile neutrinos at various energy scales

**GUT:** see-saw models of neutrino mass, leptogenesis

**TeV:** production at LHC and impact on electroweak precision observables

**keV:** (warm) dark matter candidates

**eV:** SBL and LBL oscillation experiments

**sub-eV:**  $\theta_{13}$  - reactors and solar neutrinos

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# Light Sterile Neutrinos: Huge Interest in the Community

## White Paper on Light Sterile Neutrino Searches and Related Phenomenology

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## Snowmass White Paper

Acero et al., arXiv:2203.07323 [hep-ex]

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## eV-scale Sterile Neutrinos

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Giunti, Lasserre, arXiv:1901.08330 [hep-ph]

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## Where Are We With Light Sterile Neutrinos?

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We review the status of searches for sterile neutrinos in the  $\sim 1$  eV range, with an emphasis on the latest results from short baseline oscillation experiments and how they fit within sterile neutrino oscillation models. We present global fit results to a three-active-flavor plus one-sterile-flavor model (3+1), where we find an improvement of  $\Delta\chi^2 = 35$  for 3 additional parameters compared to a model with no sterile neutrino. This is a  $5\sigma$  improvement, indicating that an effect that is like that of a sterile neutrino is highly preferred by the data. However we note that separate fits to the appearance

Diaz et al., arXiv:1906.00045 [hep-ex]

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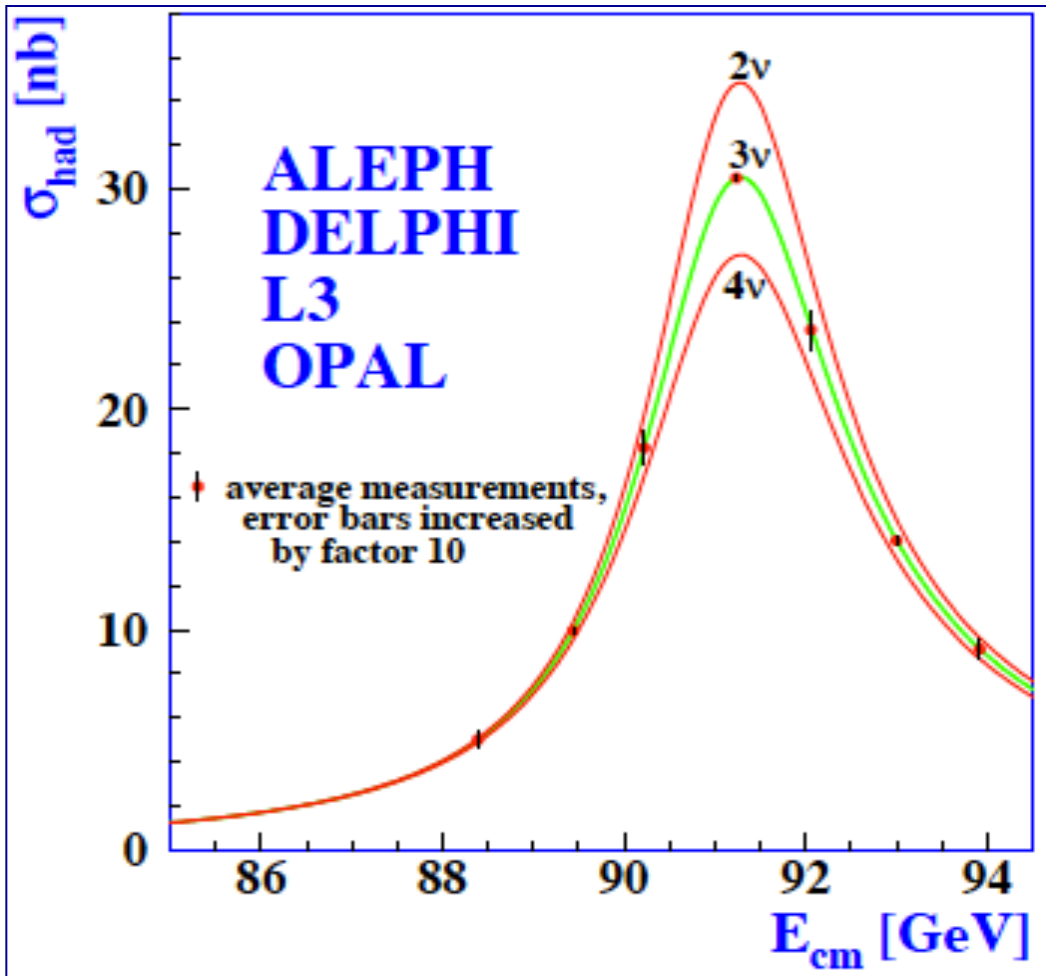
## Status of Light Sterile Neutrino Searches

Sebastian Böser<sup>1</sup>, Christian Buck<sup>2</sup>, Carlo Giunti<sup>3</sup>, Julien Lesgourgues<sup>4</sup>, Livia Ludhova<sup>5,6</sup>, Susanne Mertens<sup>7,8</sup>, Anne Schukraft<sup>9</sup>, Michael Wurm<sup>1,\*</sup>

Böser et al., arXiv:1906.01739 [hep-ex]

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# Why 3 Weak Flavor States?



Precision data on the Z-decay width  
at the  $e^+e^-$  collider at LEP

$$e^+e^- \rightarrow Z \xrightarrow{\text{invisible}} \sum_{a=\text{active}} \nu_a \bar{\nu}_a$$

$$N_{\nu_{\text{active}}} = 2.9840 \pm 0.0082$$

[LEP, Phys. Rept. 427 (2006) 257, hep-ex/0509008]

3 light active flavor neutrinos

$$\nu_e \quad \nu_\mu \quad \nu_\tau$$

<http://pdg.lbl.gov/2016/reviews/rpp2016-rev-light-neutrino-types.pdf>

# Why do we need light sterile neutrinos?

## Long-standing saga of eV-scale anomalies!

- ▶ **1995 LSND Anomaly:  $\sim 3.8\sigma$**  [PRD 64 (2001) 22, 112007]
- ▶ **2008 MiniBooNE Anomaly (combined  $\nu$  & antineutrino):  $\sim 4.8\sigma$**  [PRL 121 (2018) 22, 221801]
- ▶ **2011 Reactor Antineutrino Anomaly:  $\sim 3\sigma$**  [PRD 83 (2011) 073006, PRC 84 (2011) 024617]
- ▶ **2005 Gallium Neutrino Anomaly:  $\sim 2.9\sigma$**  [PRC 83 (2011) 065504, PLB 795 (2019) 542]
- ▶ **NEOS:  $\sim 3\sigma$**  [PRL 118, 121802 (2017)]
- ▶ **DANSS:  $\sim 2.8\sigma$**  [PLB 787 (2018) 56]
- ▶ **Neutrino-4:  $\sim 2.8\sigma$**  [JETP Lett. 109 (2019) 4, 213]

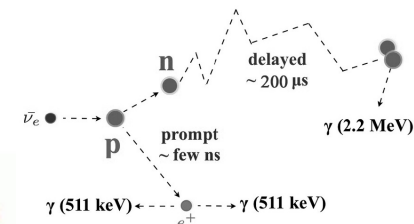
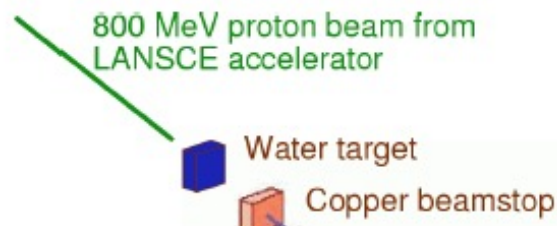
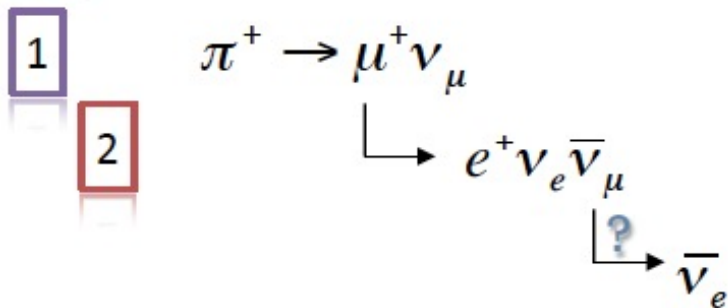
# The Liquid Scintillator Neutrino Detector

## The “LSND Anomaly”

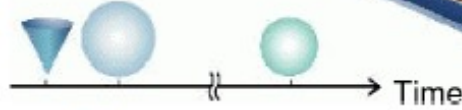
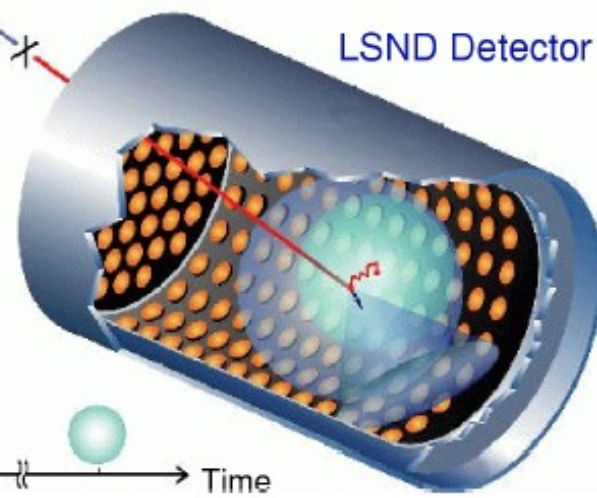
(1993 - 1998)

Los Alamos National Lab (LANL)

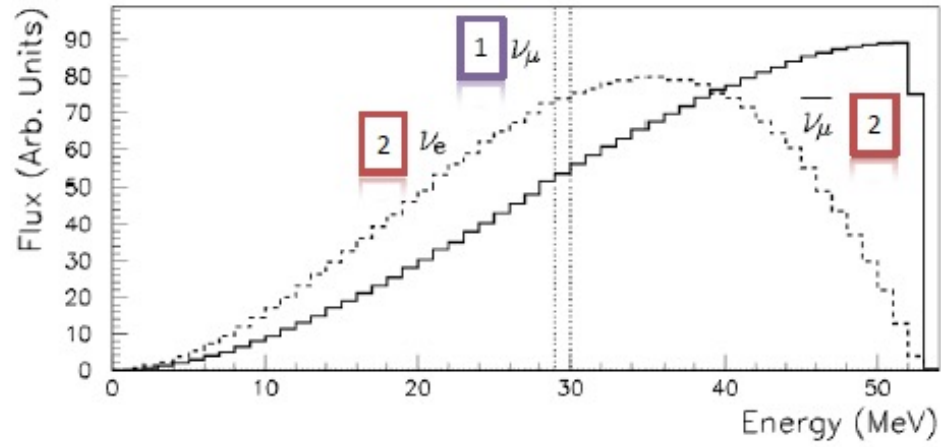
# Liquid Scintillator Neutrino Detector (LSND)



LSND Detector



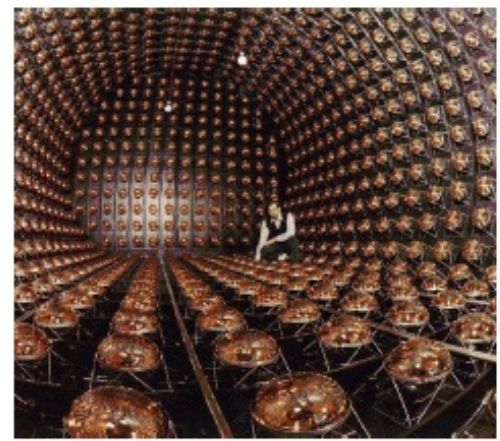
Beam produced by decays at rest is precisely known



Look for electron anti-neutrinos in a beam with well-predicted fluxes and small electron anti-neutrino background

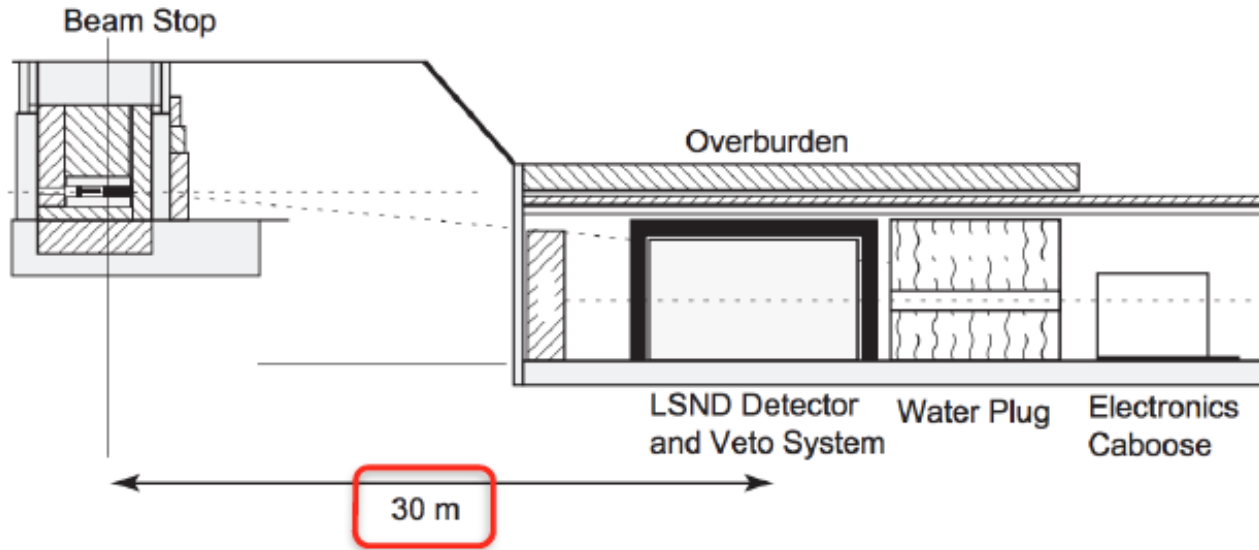
$\bar{\nu}_e$  detection via inverse-beta-decay:  $\bar{\nu}_e + p \rightarrow e^+ + n$  (coincidence signal)

neutron captures to produce a 2.2 MeV gamma

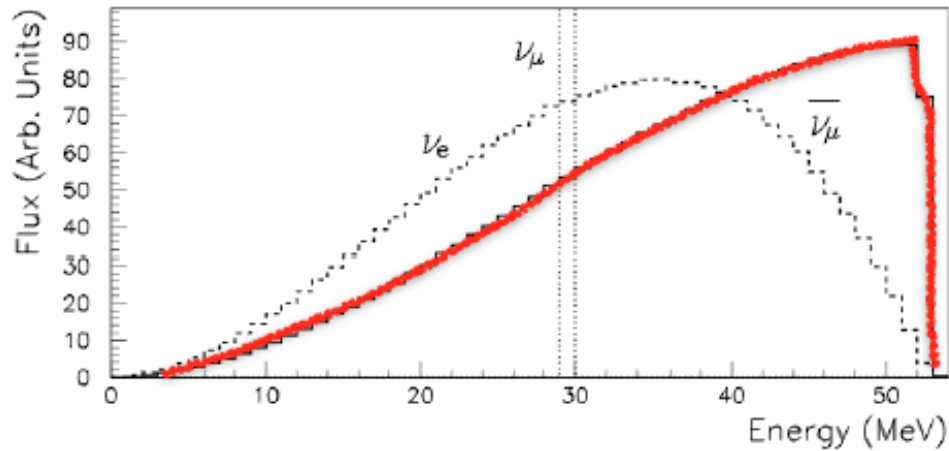




# Liquid Scintillator Neutrino Detector (LSND)



$L$

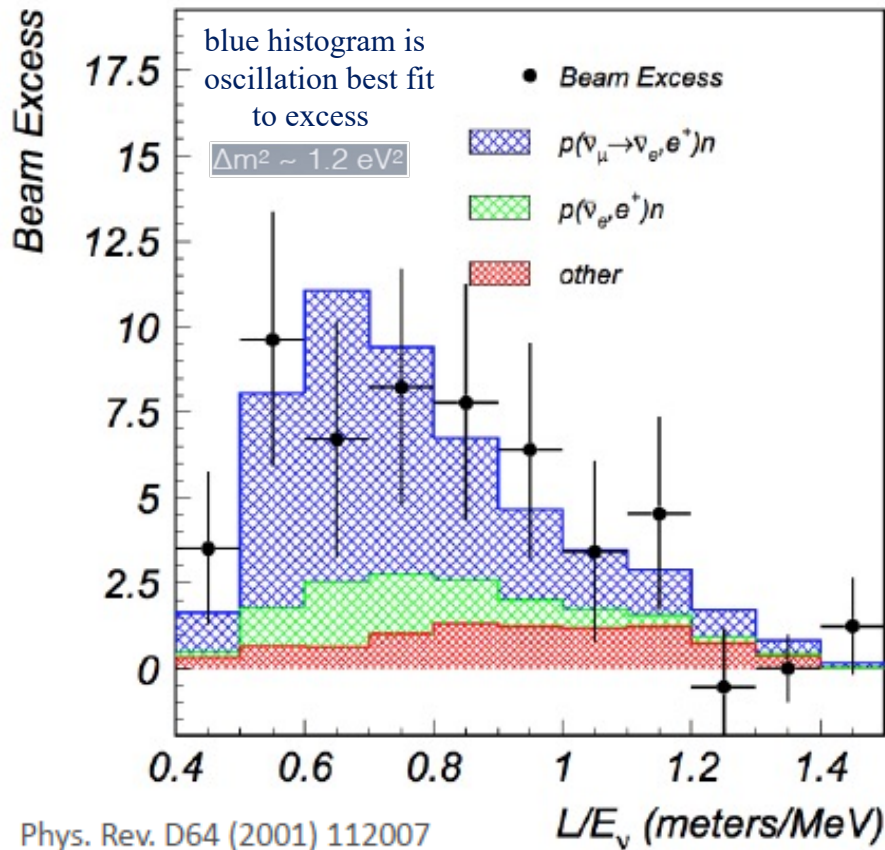


$E$

$$L/E \sim 1 \text{ m/MeV}$$

Courtesy David Schmitz

# Liquid Scintillator Neutrino Detector (LSND)



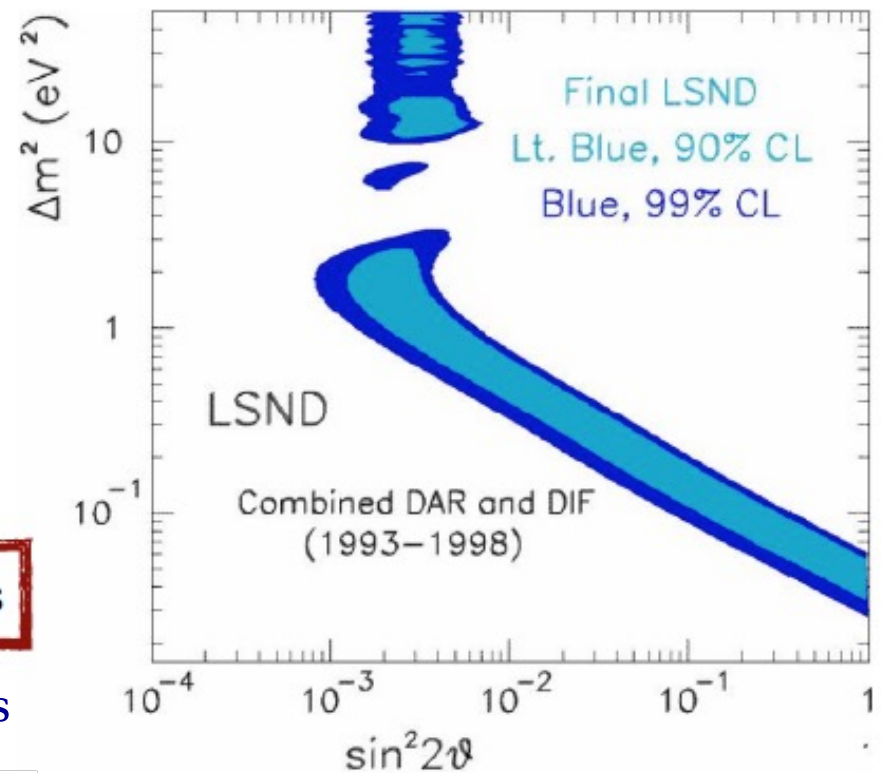
L/E distribution of the most signal-like events

Saw an excess of  $87.9 \pm 22.4 \pm 6.0$  events

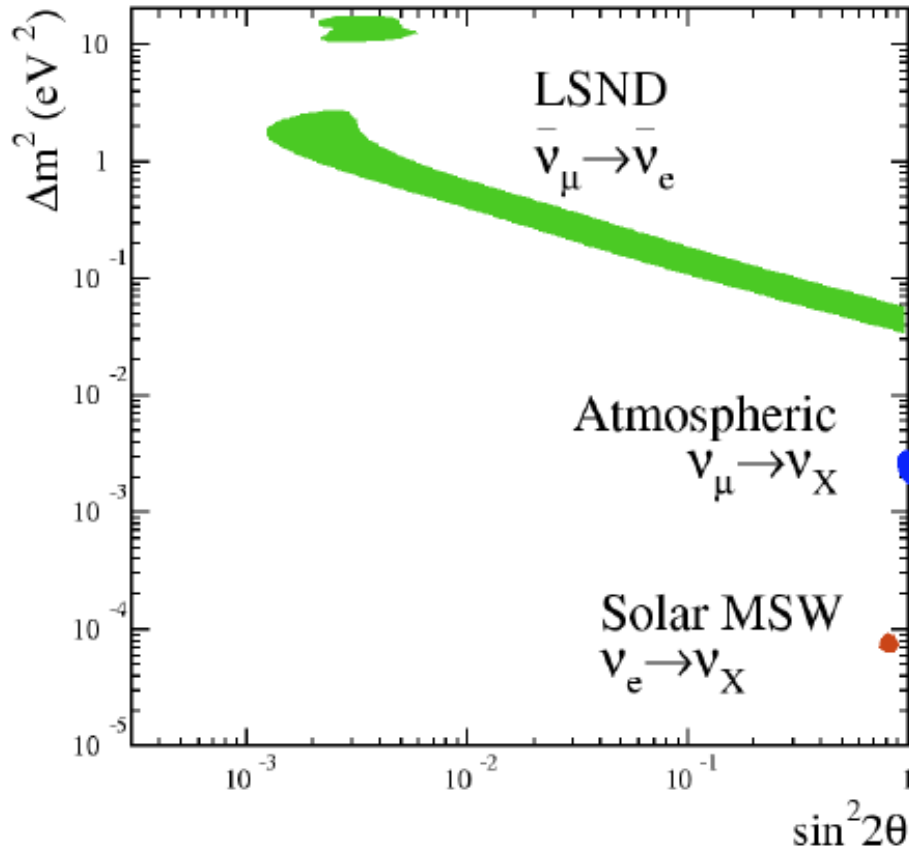
$3.8\sigma$  excess of  $\bar{\nu}_e$  events in a beam of  $\bar{\nu}_\mu$

Observed excess described by best fit oscillation probability:

$$P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) = (0.264 \pm 0.067 \pm 0.045)\%$$



# How to accommodate the LSND result?



$$\Delta m_{LSND}^2$$

$$L/E \sim 1 \text{ km/GeV}$$

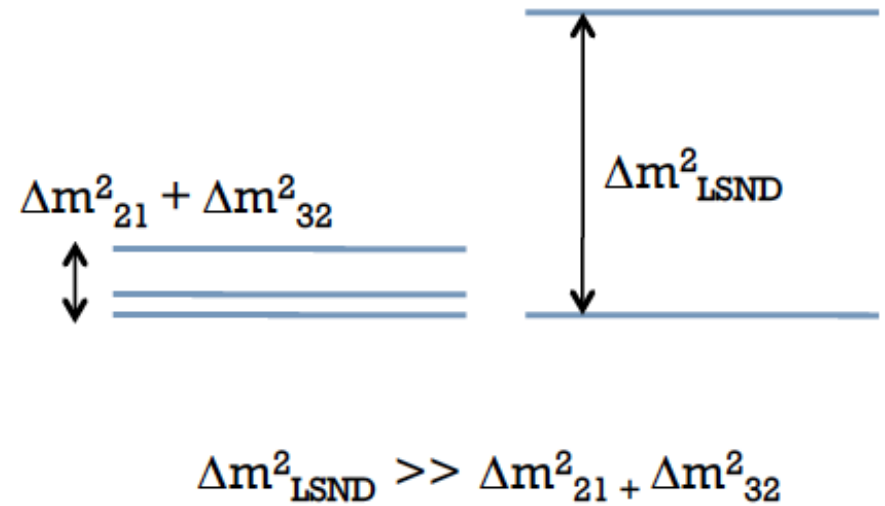
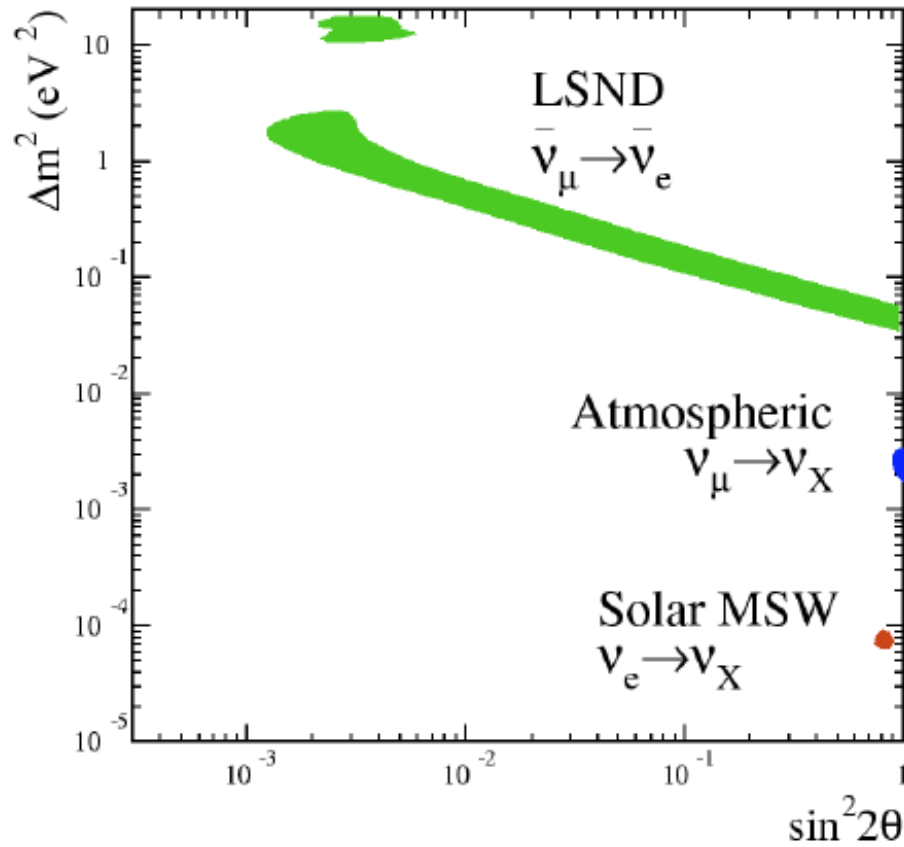
$$|\Delta m_{31}^2| \approx 2.4 \times 10^{-3} eV^2$$

$$L/E = 500 \text{ km/GeV}$$

$$\Delta m_{21}^2 \approx 7.5 \times 10^{-5} eV^2$$

$$L/E = 15,000 \text{ km/GeV}$$

# How to accommodate the LSND result?



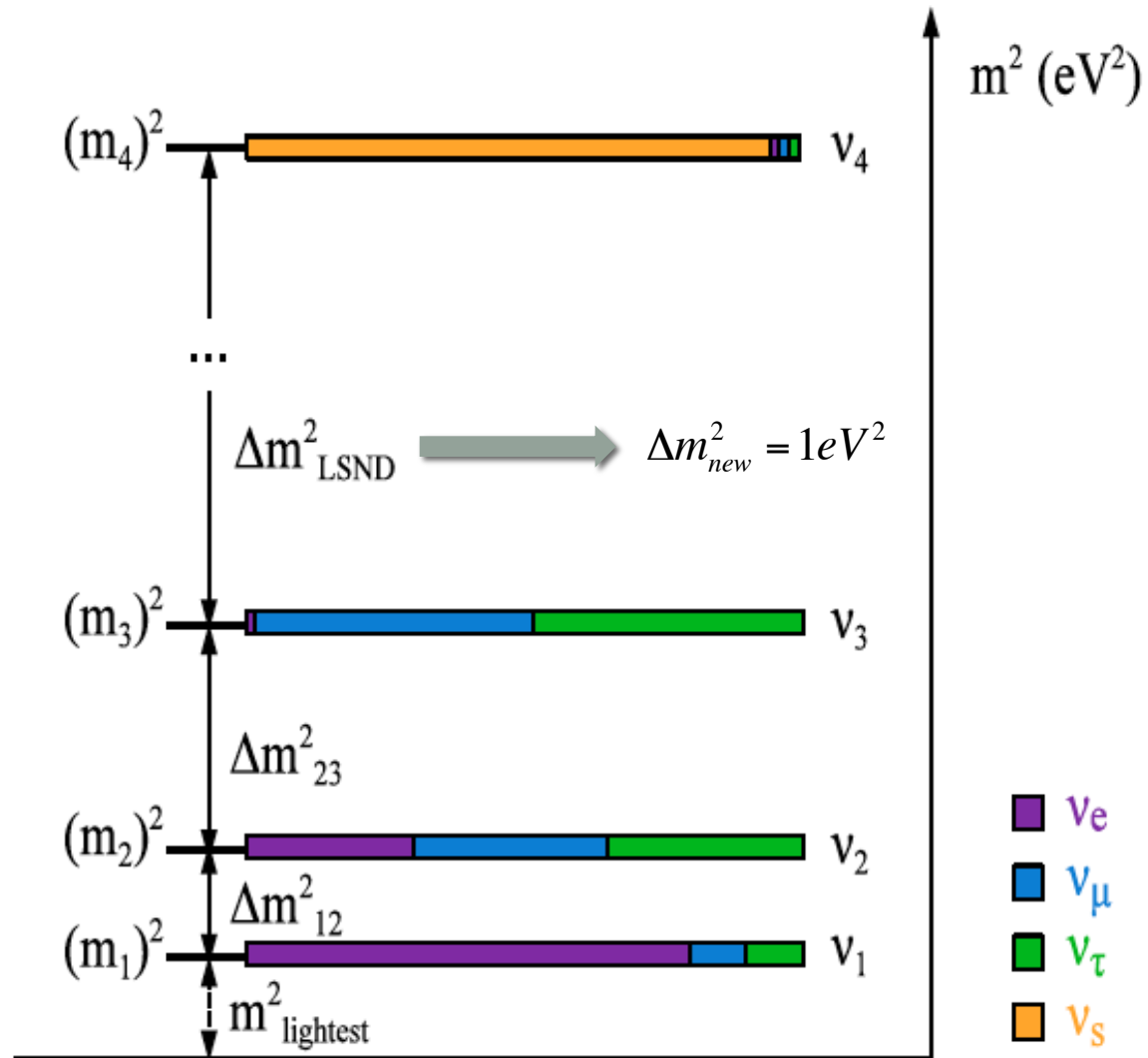
# How to accommodate the LSND result?

## Sterile neutrino

Additional neutrino flavor and mass state that has no weak interactions through the standard W/Z bosons

Mass state accessed only through mixing with standard model neutrinos

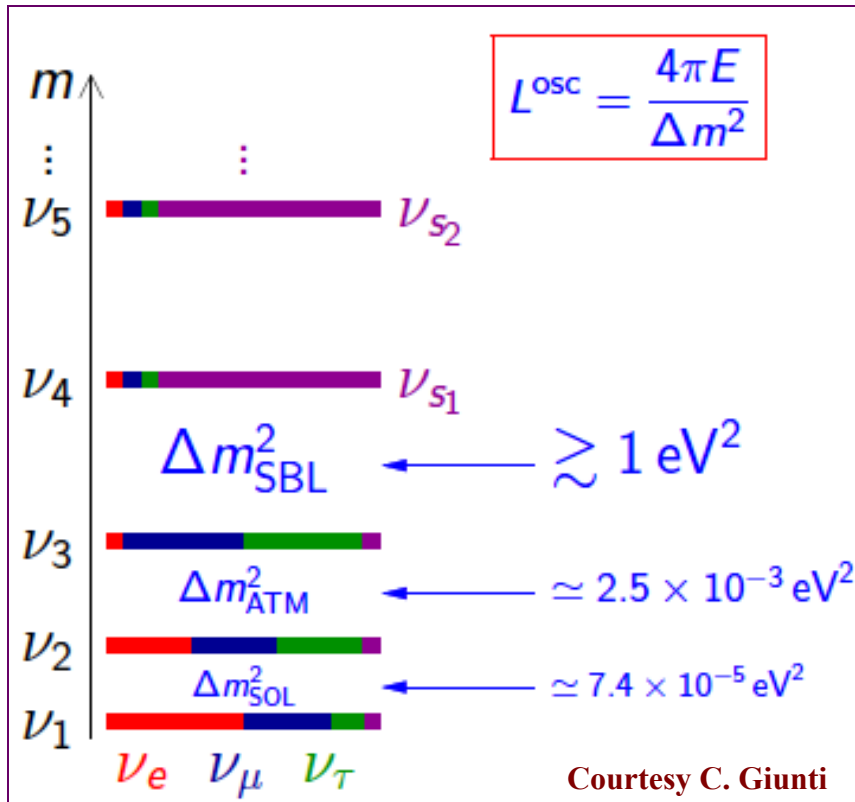
**Sterile neutrinos: singlets of  $SU(3)_C \times SU(2)_L \times U(1)_Y$  gauge group and provide economical extension of the Standard Model**



# LSND data point towards a light eV-Scale Sterile Neutrino

LSND results hint towards high  $\Delta m^2 \approx 0.1 - 10 \text{ eV}^2$  oscillation

Require additional neutrinos with masses at eV-scale



- $\nu_s$ : Sterile states (no weak interactions)
- Can feel gravity
- Can affect oscillations through mixing
- Well postulated in see-saw models
- In all these models, sterile neutrinos couple to active neutrinos via Yukawa couplings

Introduce  $\nu_R$  in the SM:

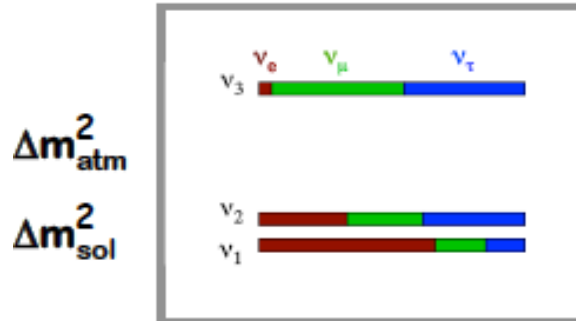
$$\text{Dirac mass } m_D \overline{\nu_R} \nu_L + \text{Majorana mass } m_M \overline{\nu_R^c} \nu_R$$

6 massive Majorana neutrinos :  $(\nu_{eL}, \nu_{\mu L}, \nu_{\tau L}) + (\nu_{eR}, \nu_{\mu R}, \nu_{\tau R})$

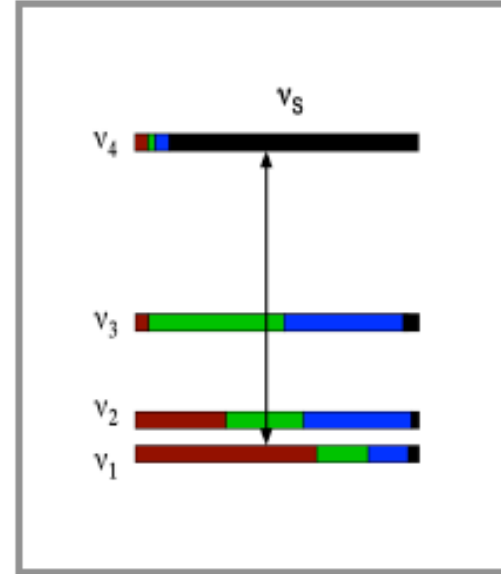
Light left-handed anti- $\nu_R$  = Light left-handed sterile neutrino :  $\nu_R^c \rightarrow \nu_{sL}$

# One Light eV-Scale Sterile Neutrino

## 3ν scheme



## 3+1 scheme



$$\Delta m_{41}^2 \sim 1 \text{ eV}^2$$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

↑  
SBL

Add one sterile  $\nu$  with three active ones at the eV-scale

Small perturbation of 3ν mixing

$$|U_{e4}|^2 \ll 1, |U_{\mu4}|^2 \ll 1, |U_{\tau4}|^2 \ll 1, |U_{s4}|^2 \approx 1$$

# 3+1 Short-Baseline Oscillation

Appearance ( $\alpha \neq \beta$ )

$$P_{\nu_\alpha \rightarrow \nu_\beta}^{\text{SBL}(-)} \simeq \sin^2 2\vartheta_{\alpha\beta} \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{\alpha\beta} = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

Disappearance

$$P_{\nu_\alpha \rightarrow \nu_\alpha}^{\text{SBL}(-)} \simeq 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{\alpha\alpha} = 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

SBL

- ▶ Amplitude of  $\nu_e$  disappearance:

$$\sin^2 2\vartheta_{ee} = 4|U_{e4}|^2 (1 - |U_{e4}|^2) \simeq 4|U_{e4}|^2$$

- ▶ Amplitude of  $\nu_\mu$  disappearance:

$$\sin^2 2\vartheta_{\mu\mu} = 4|U_{\mu 4}|^2 (1 - |U_{\mu 4}|^2) \simeq 4|U_{\mu 4}|^2$$

- ▶ Amplitude of  $\nu_\mu \rightarrow \nu_e$  transitions:

$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2 |U_{\mu 4}|^2 \simeq \frac{1}{4} \sin^2 2\vartheta_{ee} \sin^2 2\vartheta_{\mu\mu}$$

quadratically suppressed for small  $|U_{e4}|^2$  and  $|U_{\mu 4}|^2$



Appearance-Disappearance Tension

See reviews by C. Giunti

- ▶ 6 mixing angles
- ▶ 3 Dirac CP phases
- ▶ 3 Majorana CP phases

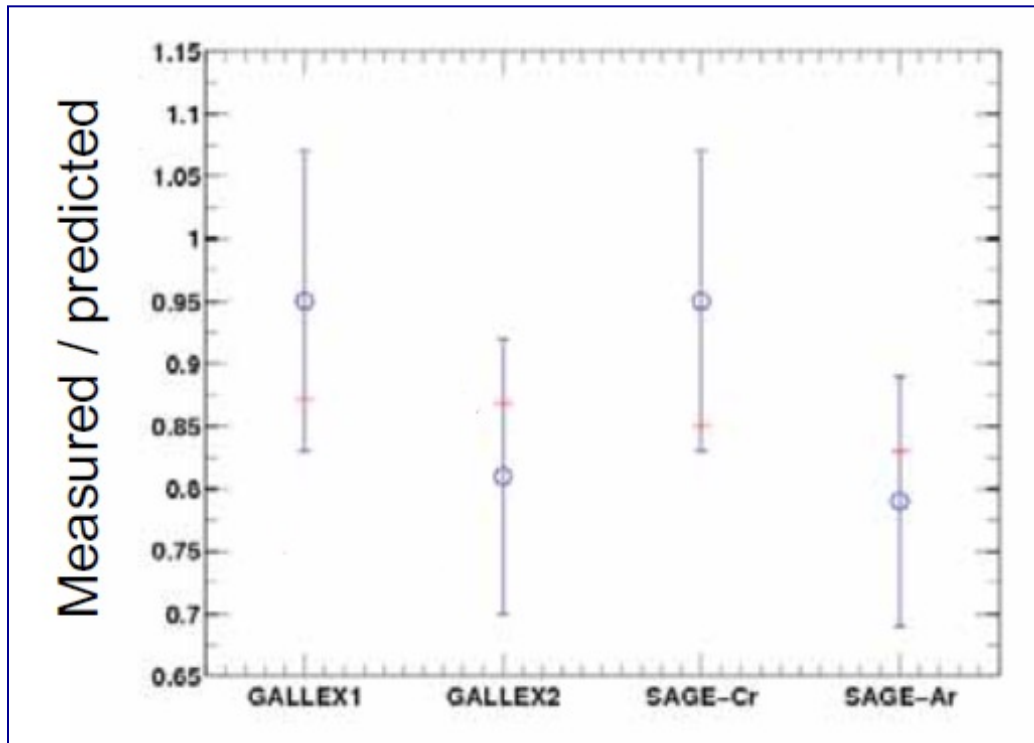


**Short-baseline means :  $L/E \sim 1$  (m/MeV or km/GeV)**

It covers a wide range of experiments

- Radioactive  $\nu_e/\bar{\nu}_e$  Source experiments  
( $L/E \sim 1$  m/1 MeV)
- Reactor  $\bar{\nu}_e$  experiments  
( $L/E \sim 5$  m/5 MeV)
- Accelerator produced  $\nu$  experiments  
( $L/E \sim 1$  km/1 GeV)
- Atmospheric Neutrinos in IceCube  
( $L/E \sim 1000$  km/1 TeV)

# Gallium Neutrino Anomaly

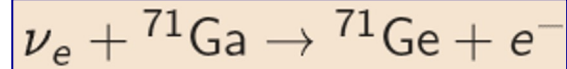


Calibration measurements for the GALLEX & SAGE solar neutrino detectors using intense radioactive  $\nu_e$  fluxes from  $^{51}\text{Cr}$  &  $^{37}\text{Ar}$

$^{51}\text{Cr}$  : 747 KeV (82%)

$^{37}\text{Ar}$  : 811 KeV (90%)

Detection process:



Measurements consistently lower than expectation

Suggests possible  $\nu_e$  disappearance at  $2.7\sigma$  due to active – sterile oscillation

How well do we know the cross sections of the radiochemical detection processes?

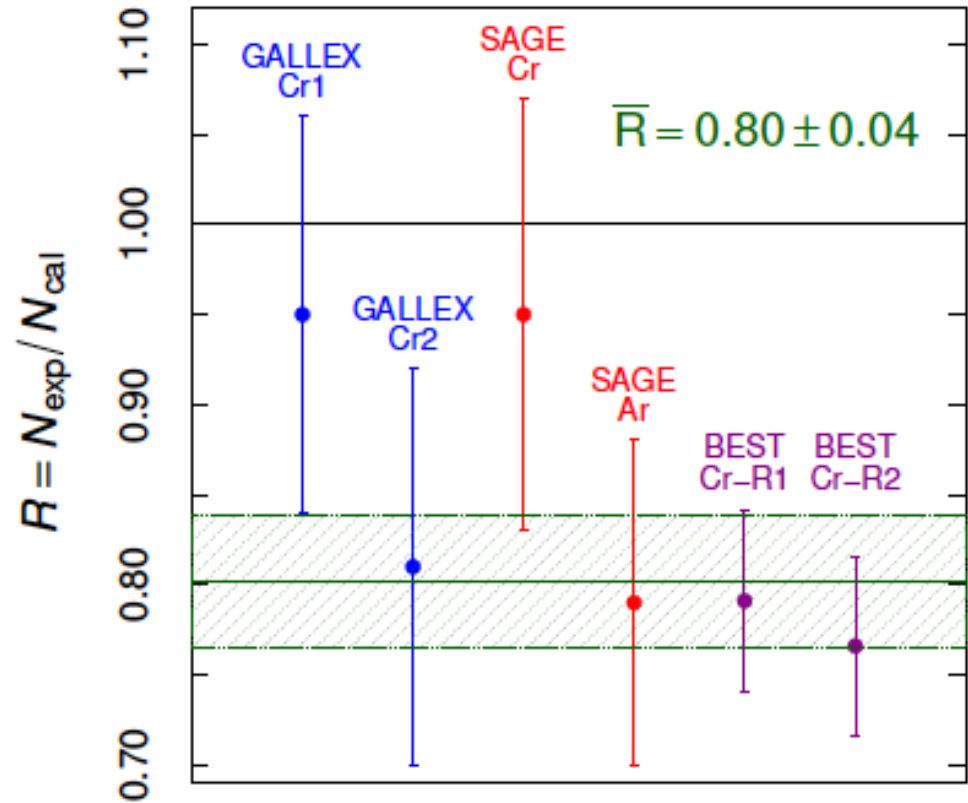
## Baksan Experiment on Sterile Transitions (BEST)

arXiv:2109.11482 and arXiv:2109:14654



Gallium neutrino anomaly has been reinforced after the new measurements from BEST

**More than  $5\sigma$  deficit of  $\nu_e$**

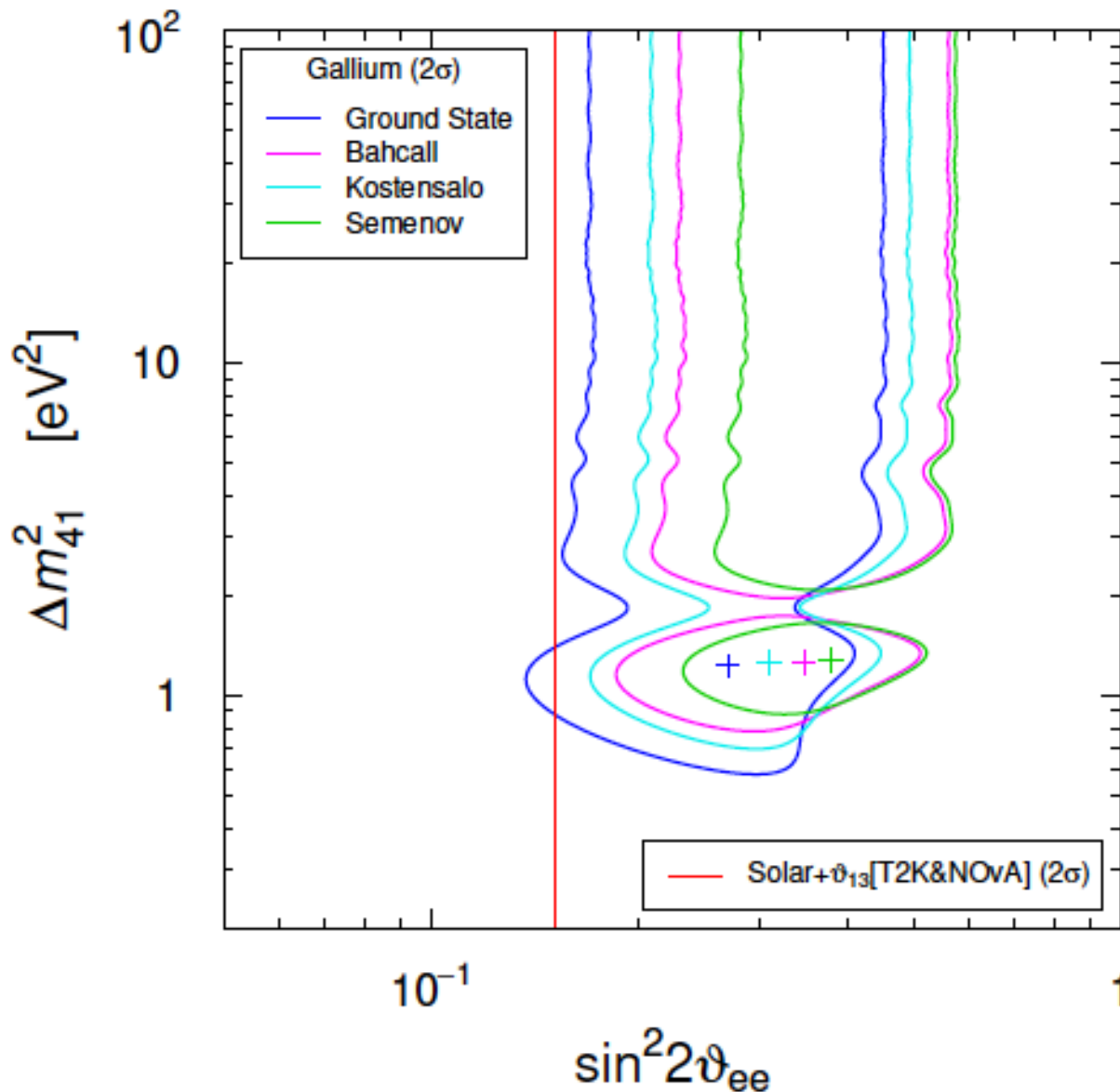


$$\langle L \rangle_{\text{GALLEX}} \simeq 1.9 \text{ m} \quad \langle L \rangle_{\text{SAGE}} \simeq 0.6 \text{ m}$$

$$\langle L \rangle_{\text{BEST}}^{R1} \simeq 0.7 \text{ m} \quad \langle L \rangle_{\text{BEST}}^{R2} \simeq 1.1 \text{ m}$$

$$\Delta m_{\text{SBL}}^2 \gtrsim 1 \text{ eV}^2 \gg \Delta m_{\text{ATM}}^2$$

# Gallium Anomaly vs. Solar Neutrino Bound



Both Gallium and Solar Neutrino Experiments detect  $\nu_e$

Gallium anomaly is in strong tension with the solar neutrino bound

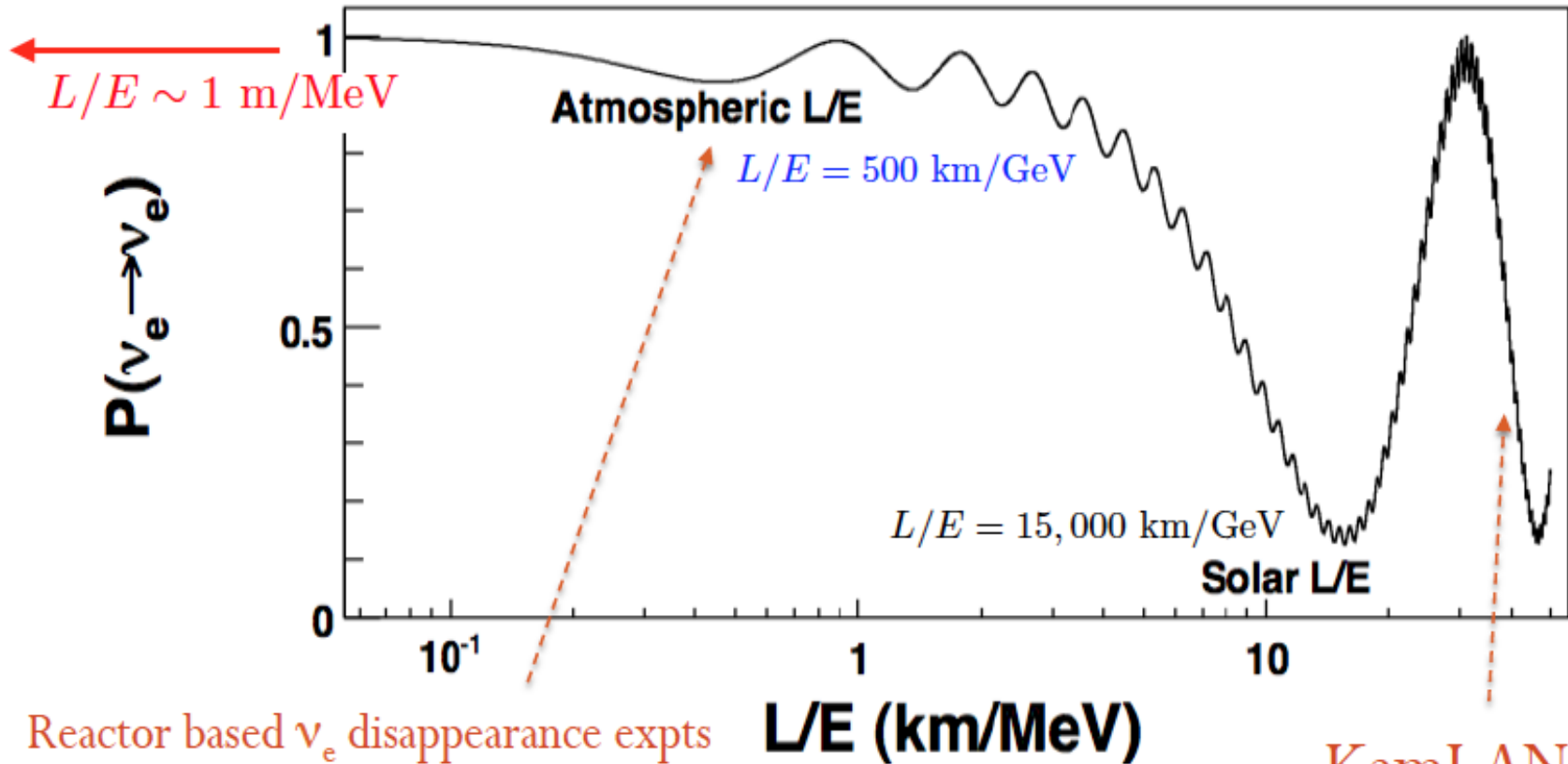
Giunti, Li, Ternes, Tyagi, Xin, arXiv:2209.00916

# **Reactor Antineutrinos at very Short-Baselines**

**The Evolution of**

**“Reactor Antineutrino Anomaly”**

# Oscillations with Reactor Antineutrinos



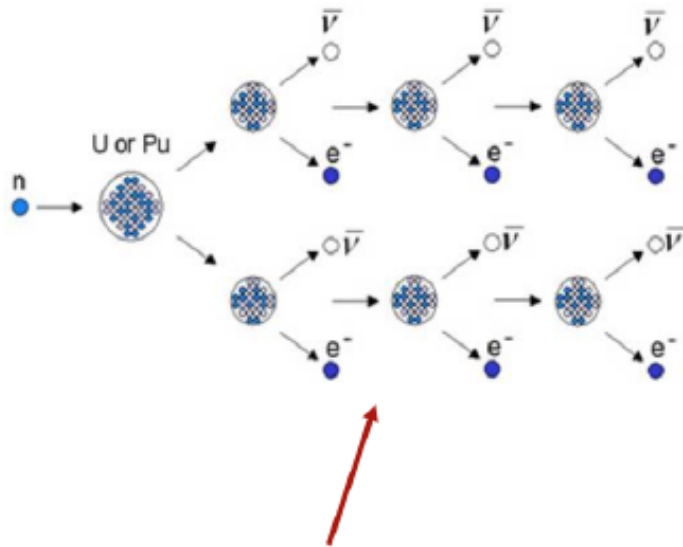
Reactor based  $\bar{\nu}_e$  disappearance expts  
such as Double Chooz and Daya Bay

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{13} \cdot \sin^2(1.27 \cdot \Delta m_{23}^2 \cdot L/E)$$

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{12} \cdot \sin^2(1.27 \cdot \Delta m_{12}^2 \cdot L/E)$$

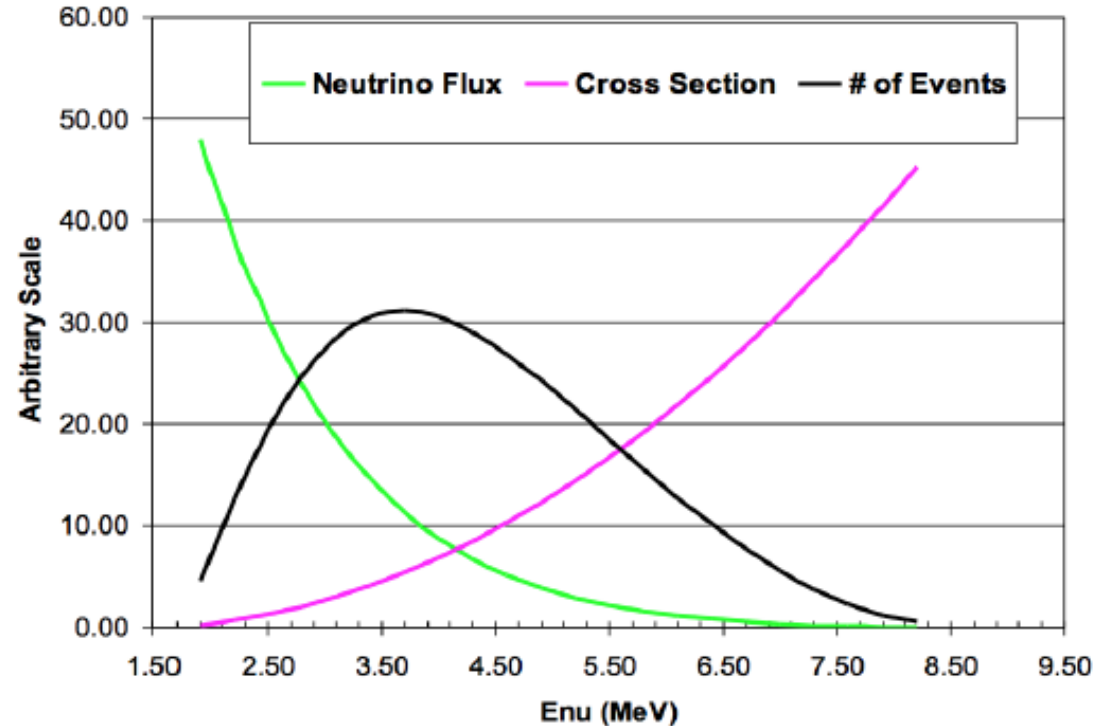
# Reactor Antineutrino Experiments

Reactor  $\bar{\nu}_e$  production



$\beta^-$  decay of neutron rich fission fragments of U and Pu

Detection through inverse  $\beta$  Decay:



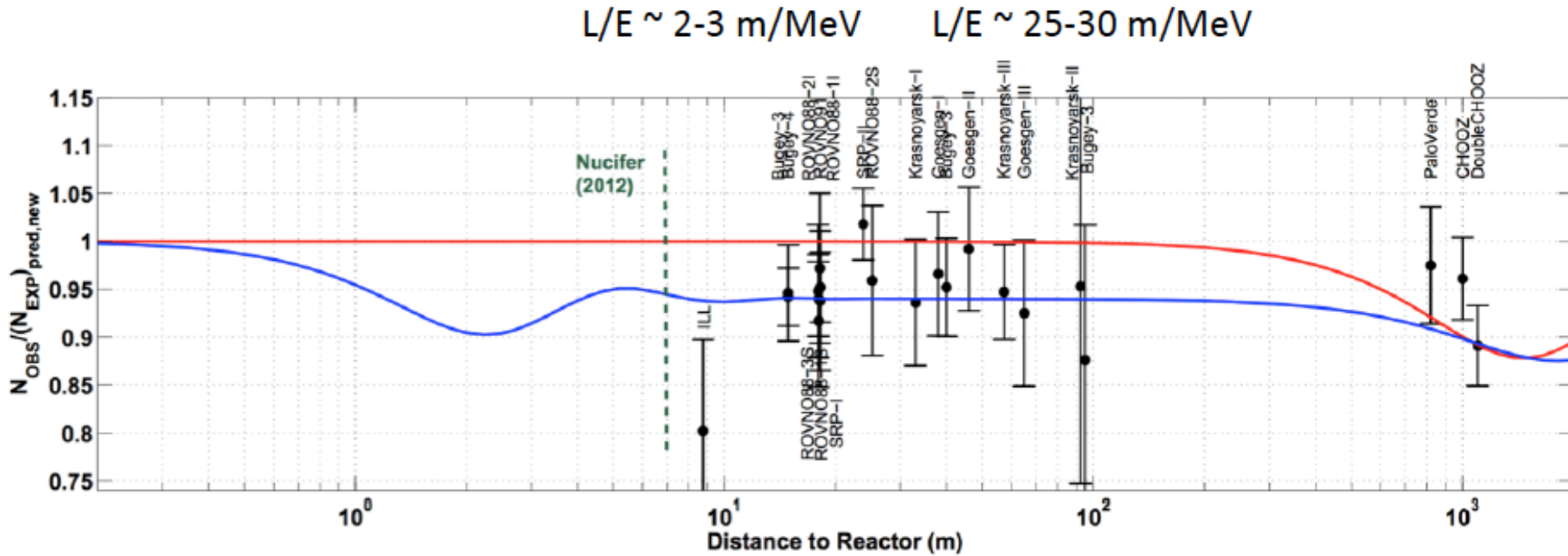
$$E_{\text{prompt}} = E_{\nu} - 0.8 \text{ MeV}$$

# Reactor Antineutrino Anomaly

- In preparation for Double Chooz analysis with a single far detector, Mueller et al. applied an improved procedure to go from measured  $^{235}\text{U}$ ,  $^{239}\text{Pu}$ , and  $^{241}\text{Pu}$   $\beta^-$  spectra (at ILL) to neutrino spectra.
  - Th. A. Mueller et al. “Improved Predictions of Reactor Antineutrino Spectra” Phys. Rev. C83 (2011) 054615; arXiv:1101.2663.
- Result was a net 3% increase in the estimated fluxes relative to previous predictions, reanalysis of past reactor experiments
  - G. Mention et al. “The Reactor Antineutrino Anomaly”, Phys. Rev. D83 (2011) 083006; arXiv:1101.2755.
- P. Huber, using a different method to go from  $\beta^-$  to neutrino spectra, finds a similar shift
  - P. Huber, “On the determination of anti-neutrino spectra from nuclear reactors”; arXiv:1106.0687.

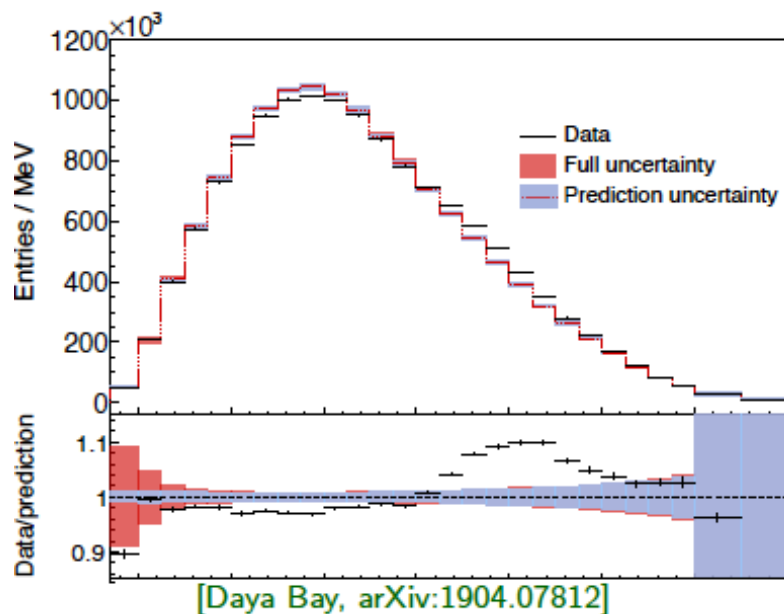
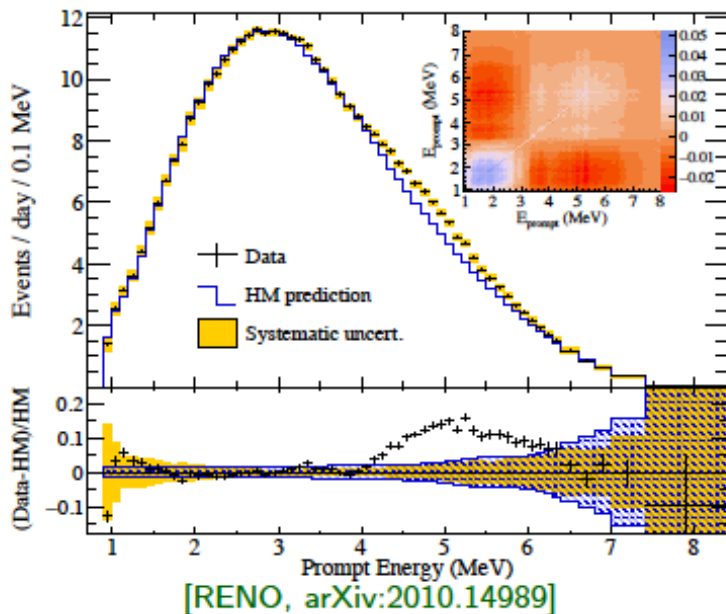


# Reactor Antineutrino Anomaly



- Looked back at ratio of the observed to predicted  $\bar{\nu}_e$  event rate for 19 different reactor neutrino experiments at baselines less than 100 m.
- Mean average ratio including correlations is  $0.927 \pm 0.023$ , indicating a 7.3% deficit at short baseline.
- Curves show fits to data **assuming standard 3 neutrino oscillations** and **assuming oscillations with one additional sterile neutrino**

# Reactor Antineutrino 5 MeV Bump (Shoulder)



- ▶ Discovered in 2014 by RENO, Double Chooz, Daya Bay.
- ▶ **Cannot** be explained by neutrino oscillations (SBL oscillations are averaged in RENO, DC, DB).
- ▶ Most likely it is due to a theoretical miscalculation of the spectrum.
- ▶ A recalculation of the spectrum can have opposite effects on the anomaly:
  - ▶ A 4-6 MeV increase of the predicted flux which explains the bump **increases** the anomaly.
  - ▶ A 1-4 MeV decrease of the predicted flux **decreases** the anomaly.

Talk by C. Giunti in NOW 2022

## Reactor antineutrino anomaly in light of recent flux model refinements

C. Giunti,<sup>1,\*</sup> Y.F. Li,<sup>2,3,†</sup> C.A. Ternes,<sup>1,‡</sup> and Z. Xin<sup>2,3,§</sup>

<sup>1</sup>*Istituto Nazionale di Fisica Nucleare (INFN), Sezione di Torino, Via P. Giuria 1, I-10125 Torino, Italy*

<sup>2</sup>*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing 100049, China*

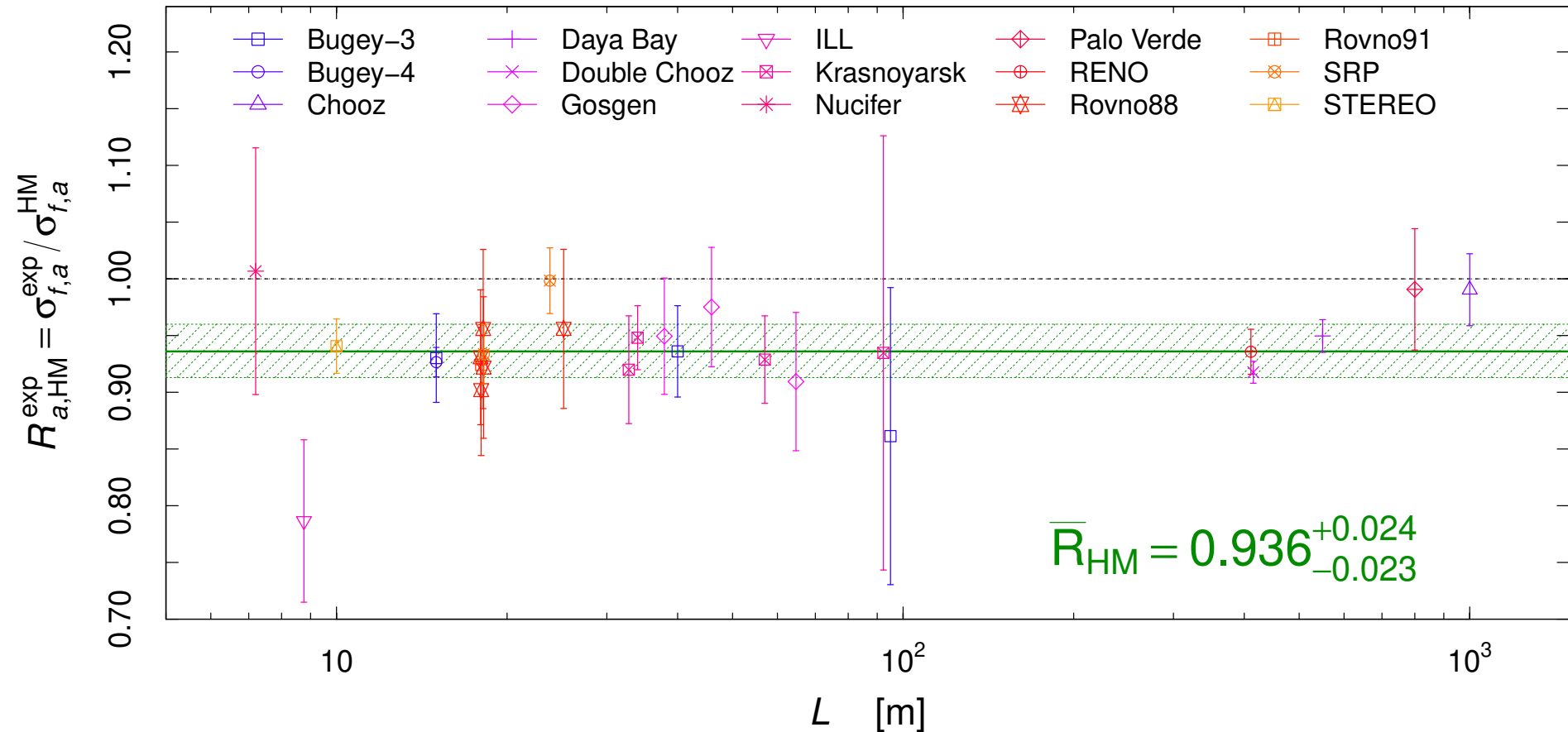
<sup>3</sup>*School of Physical Sciences, University of Chinese Academy of Sciences, Beijing 100049, China*

(Dated: 13 October 2021)

We study the status of the reactor antineutrino anomaly in light of recent reactor flux models obtained with the conversion and summation methods. We present a new improved calculation of the IBD yields of the standard Huber-Mueller (HM) model and those of the new models. We show that the reactor rates and the fuel evolution data are consistent with the predictions of the Kurchatov Institute (KI) conversion model and with those of the Estienne-Fallot (EF) summation model, leading to a plausible robust demise of the reactor antineutrino anomaly. We show that the results of several goodness of fit tests favor the KI and EF models over other models that we considered. We also discuss the implications of the new reactor flux models for short-baseline neutrino oscillations due to active-sterile mixing. We show that reactor data give upper bounds on active-sterile neutrino mixing that are not very different for the reactor flux models under consideration and are in tension with the large mixing required by the Gallium anomaly that has been refreshed by the recent results of the BEST experiment.

**Giunti, Li, Ternes, Xin, PLB 829 (2022) 137054**

# Reactor Antineutrino Anomaly

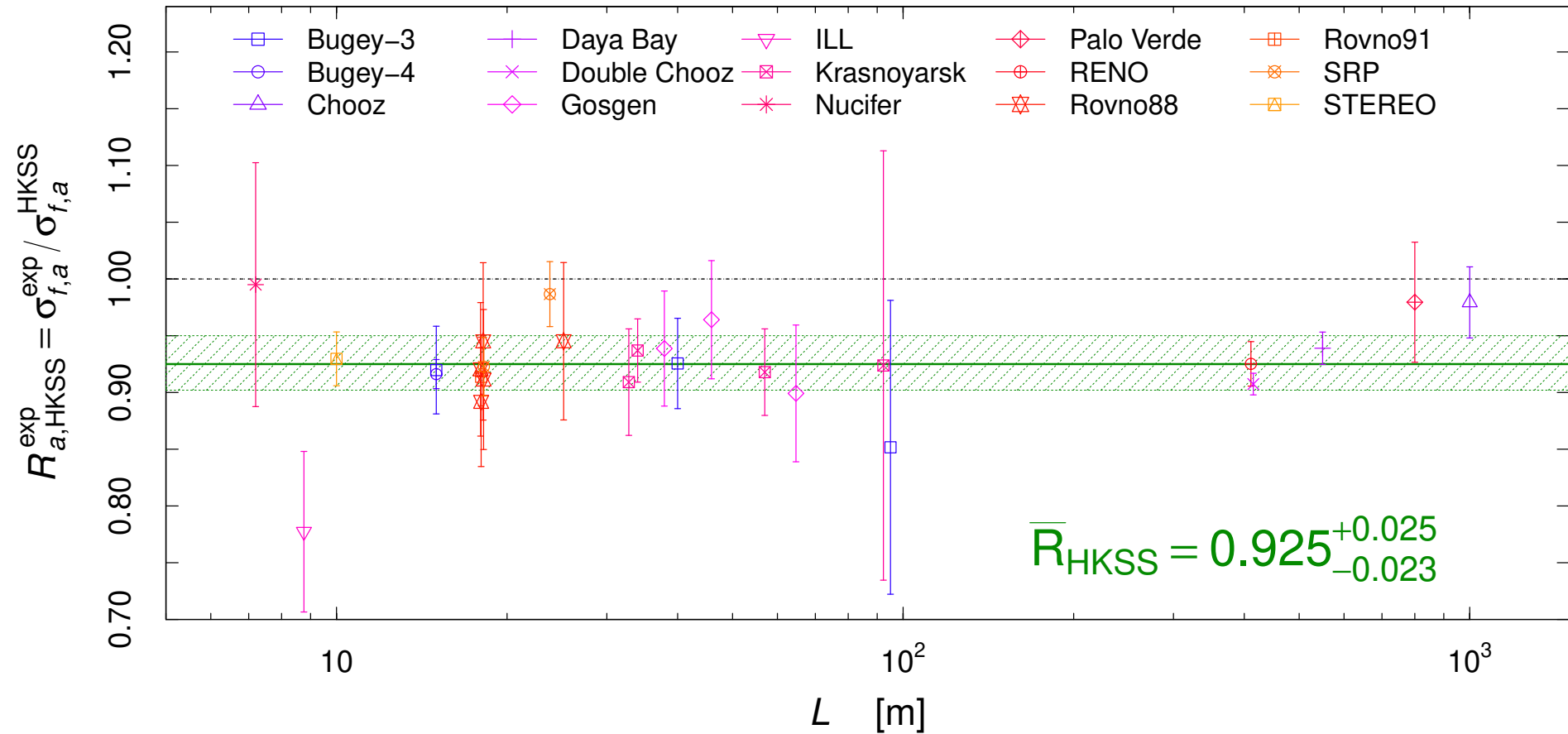


Mueller et al., arXiv:1101.2663; Huber, arXiv:1106.0687; Giunti, Li, Ternes, Xin, arXiv:2110.06820

New improved calculation of the IBD yields of the standard Huber-Mueller (HM) flux model using conversion method (new  $\beta$  spectra from  $^{235}\text{U}$  fission)

**Reactor Antineutrino Anomaly (RAA) is around  $2.5\sigma$**

# Reactor Antineutrino Anomaly

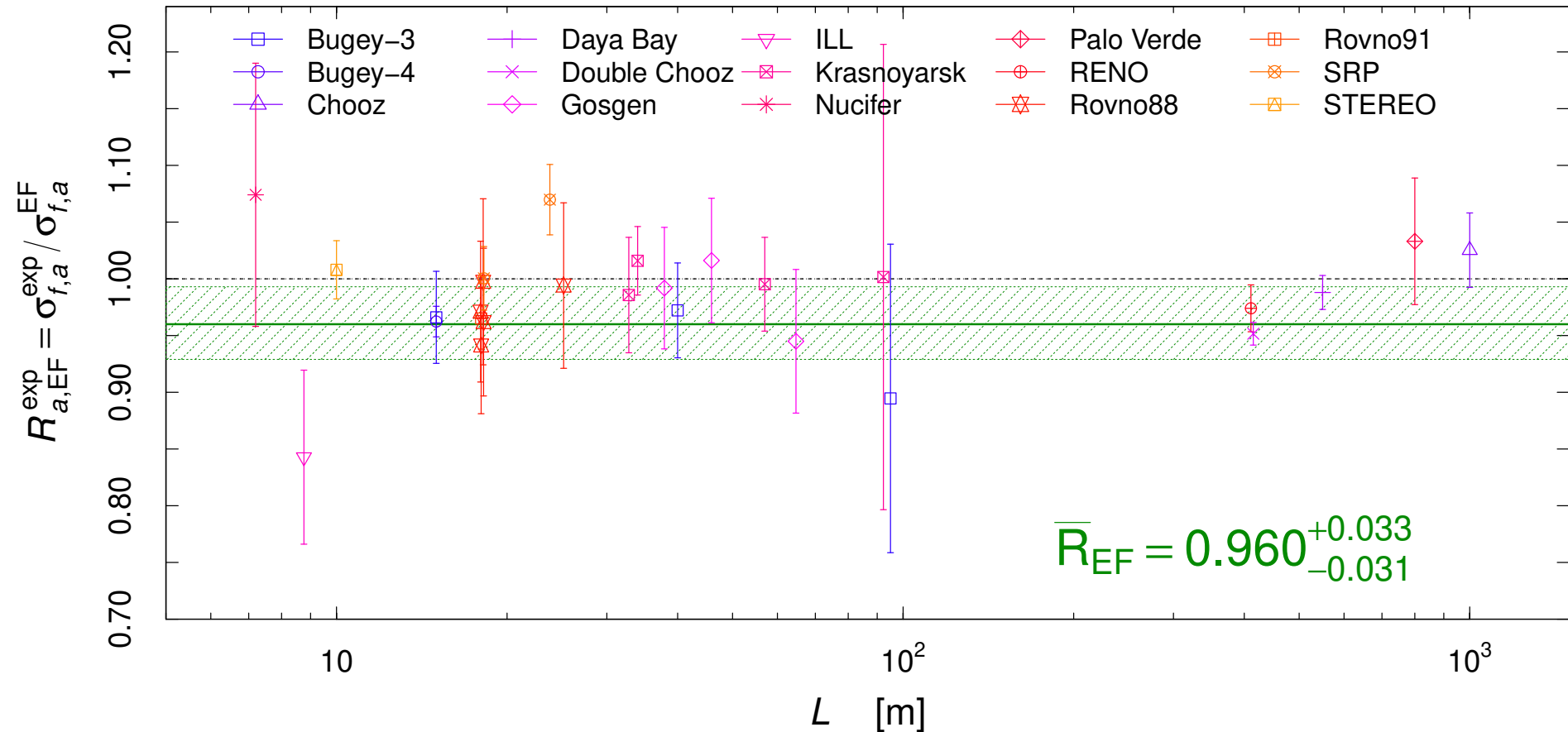


Hayen, Kostensalo, Severijns, Suhonen, arXiv:1908.08302; Giunti, Li, Ternes, Xin, arXiv:2110.06820

New HKSS fluxes in 2019 using conversion method (HM + HKSS uncertainties)

**Reactor Antineutrino Anomaly (RAA) is around  $2.9\sigma$**

# Reactor Antineutrino Anomaly

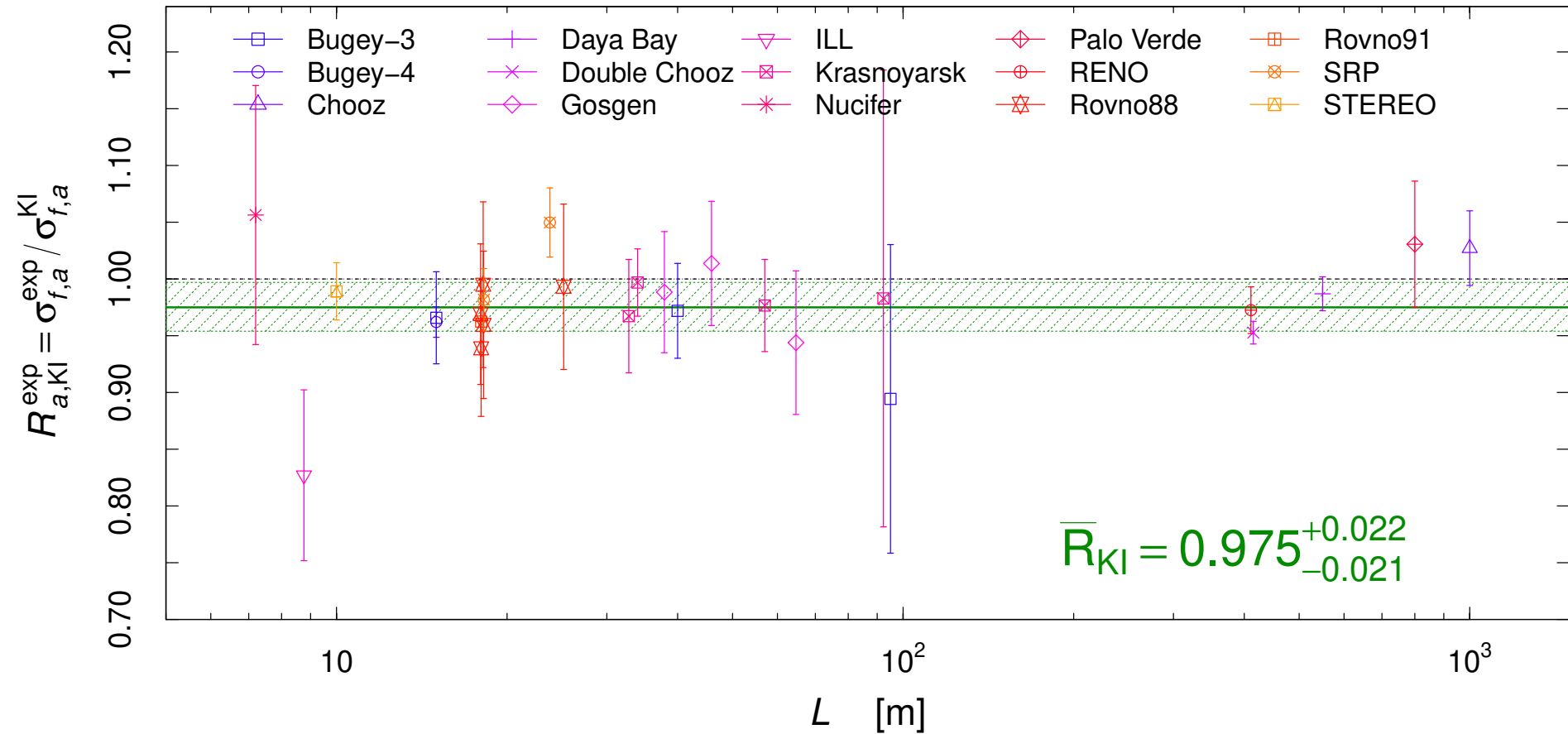


Estienne, Fallot, et al., arXiv:1904.09358; Giunti, Li, Ternes, Xin, arXiv:2110.06820

New EF fluxes in 2019 using summation method  
(5% uncertainties for  $^{235}\text{U}$ ,  $^{239}\text{Pu}$ ,  $^{241}\text{Pu}$  and 10% for  $^{238}\text{U}$ )

**Reactor Antineutrino Anomaly (RAA) is around  $1.2\sigma$**

# Reactor Antineutrino Anomaly

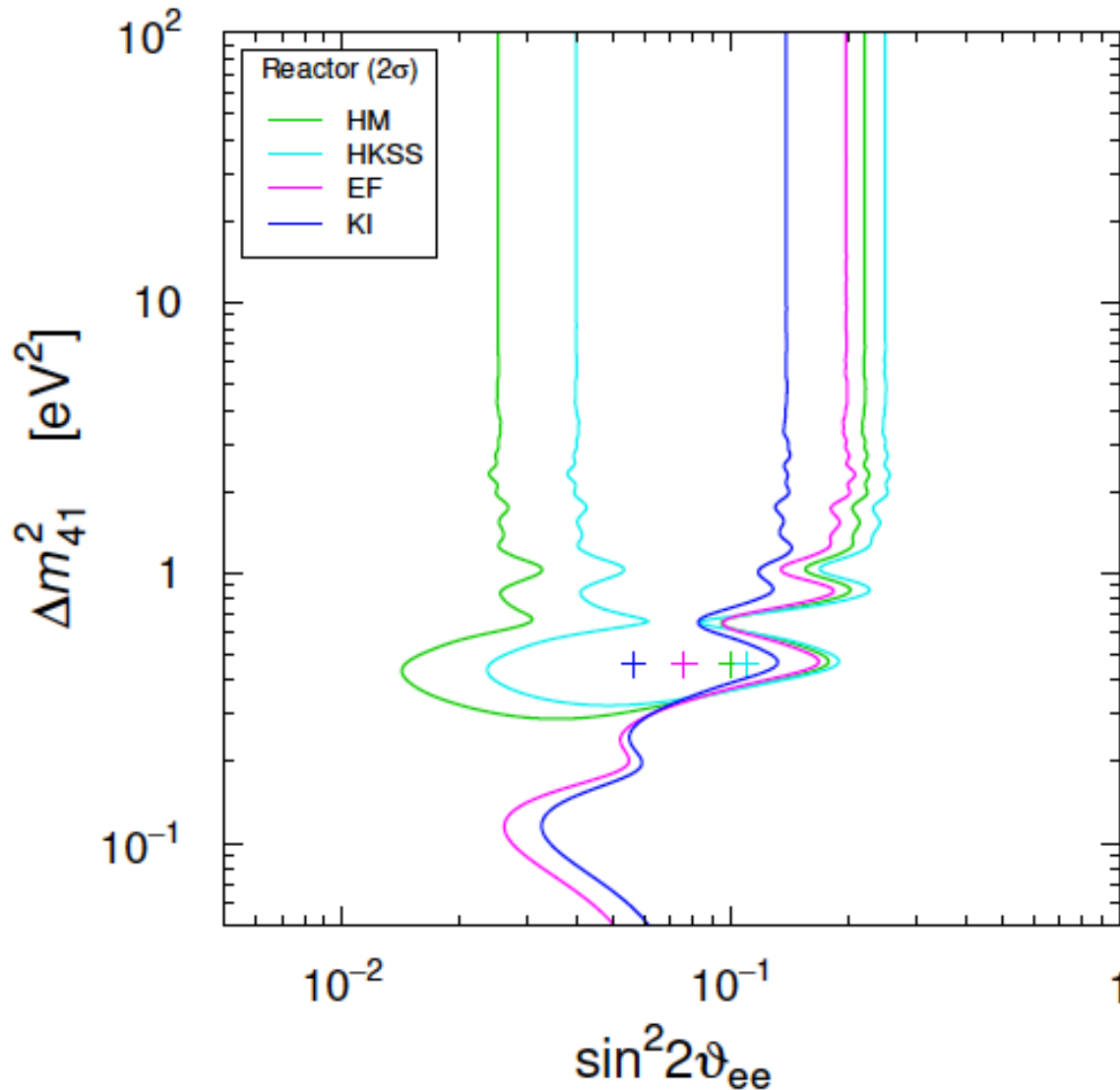


Kurchatov Institute: Kopeikin, Skorokhvatov, Titov, arXiv:2103.01684; Giunti, Li, Ternes, Xin, arXiv:2110.06820

New KI fluxes in 2021 using conversion method (HM + KI uncertainties)  
 Approximate agreement with ab initio (summation method) EF fluxes

**Reactor Antineutrino Anomaly (RAA) is around  $1.1\sigma$**

# Take Home Message



The favored KI and EF models are compatible with no active-sterile oscillations at short baselines

Independently from the reactor antineutrino flux model, we have an upper bound on

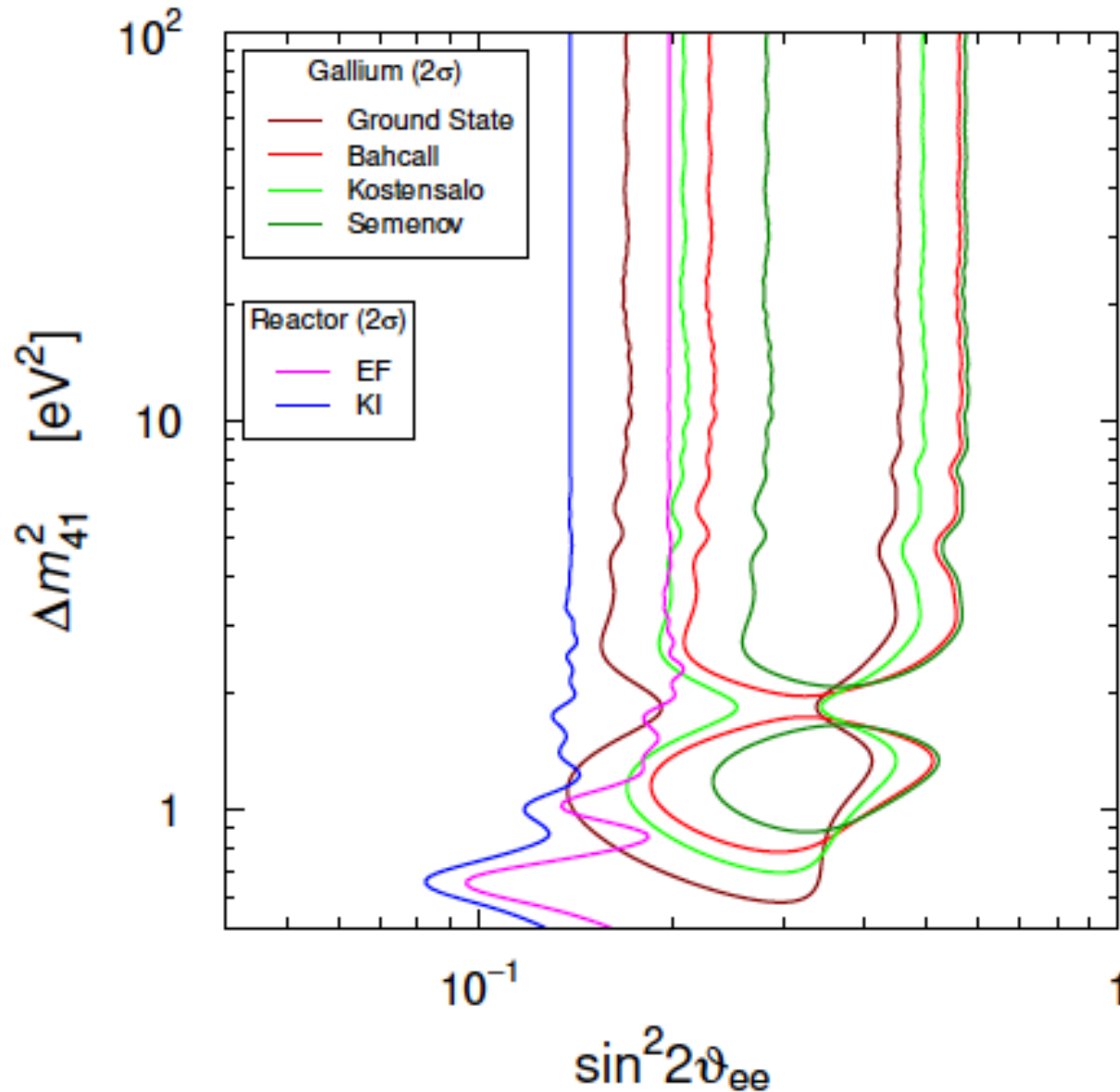
$$\sin^2 2\theta_{ee} \lesssim 0.25 \quad (2\sigma)$$



The reactor rates and the fuel evolution data are consistent with the predictions of the Kurchatov Institute (KI) conversion model and with those of the Estienne-Fallot (EF) summation model, leading to a plausible robust demise of the reactor antineutrino anomaly

Giunti, Li, Ternes, Xin, PLB 829 (2022) 137054

# Reactor Anomaly vs. Gallium Anomaly



**Another Tension!**

**Note:**

Reactor Anomaly deals w/  
electron antineutrinos,  
whereas Gallium Anomaly  
deals w/ electron neutrinos

Giunti, Li, Ternes, Tyagi, Xin, arXiv:2209.00916

# Very Short-Baseline Reactor Experiments

	DANSS	NEOS	NEUTRINO-4	PROSPECT	SoLid	STEREO
Power [MW]	3100	2815	100	85	50-80	58
Core size [cm]	$\varnothing = 320$ $h = 370$	$\varnothing = 310$ $h = 380$	$42 \times 42$ $h = 35$	$\varnothing = 51$ $h = 44$	$\varnothing = 50$ $h = 90$	$\varnothing = 40$ $h = 80$
Overburden [mwe]	50	20	3.5	< 1	10	15
Distance [m]	10.7-12.7 movable	24	6-12 movable	7-9	6-9	9-11
IBD events/day	5000	2000	200	750	~450	400
PSD	No	Yes	No	Yes	Yes	Yes
Readout	3D	1D	2D	3D	3D	2D
S/B	33	23	0.54	1.36	~3	0.9
$\sigma_E/E$ [%] at 1 MeV	34	5	16	4.5	14	8

**DANSS:** Kalinin Nuclear Power Plant (KNPP), Moscow, Russia

**Neutrino-4:** SM-3 Research Reactor at Dmitrovgrad, Russia

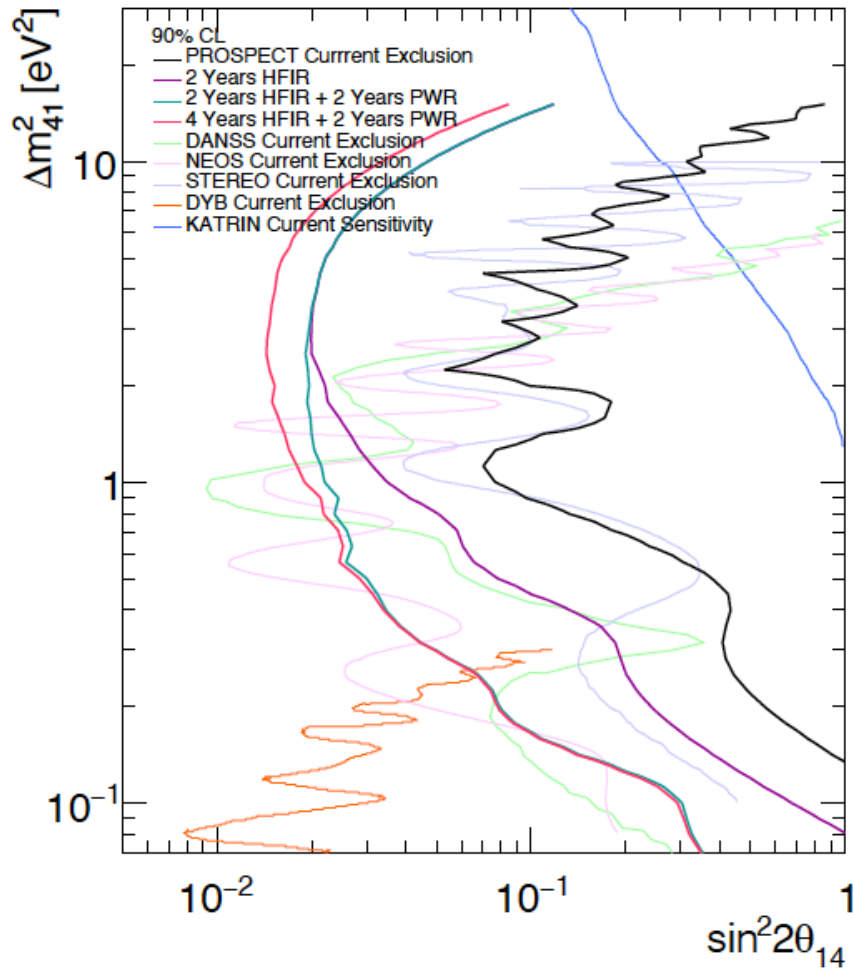
**SoLid:** SCK-CEN BR2 Research Reactor in Belgium

**NEOS:** Hanbit Nuclear Power Complex in Yeong-gwang, Korea

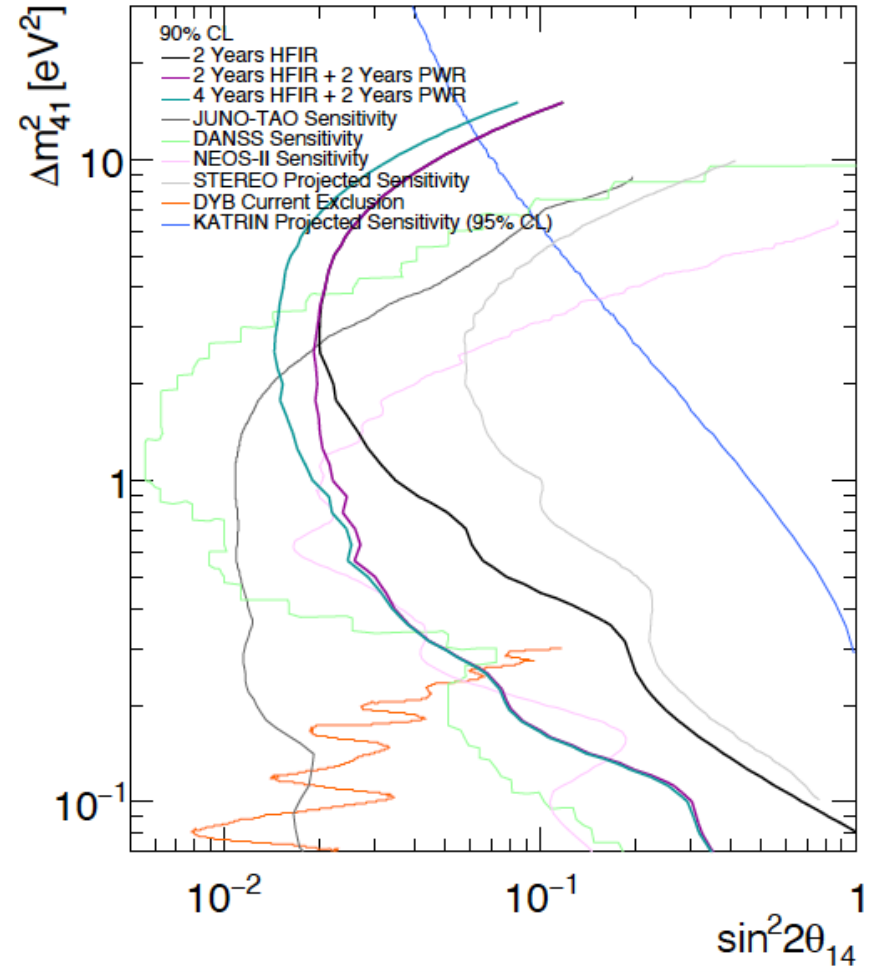
**PROSPECT:** High Flux Isotope Reactor (HFIR) at ORNL, USA

**STEREO:** High Flux Reactor of the Institute Laue-Langevin, France

# Present and Future Sensitivity from Very SBL Reactor Experiments



Present Sensitivity

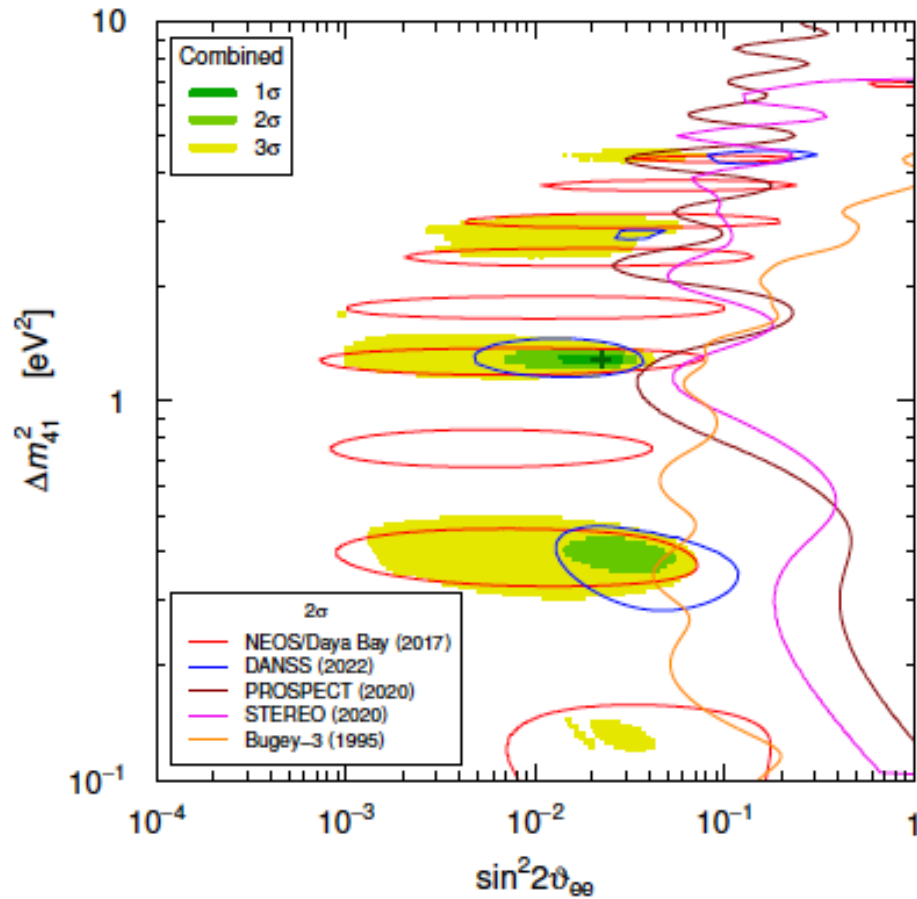


Future Sensitivity

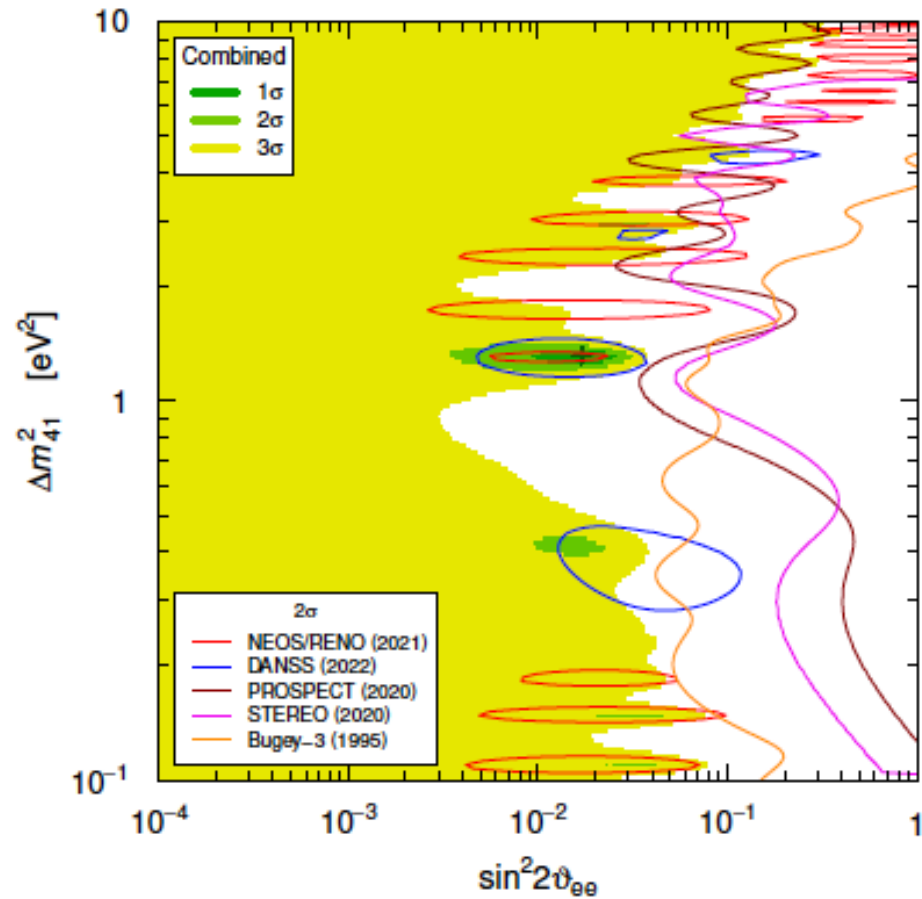
PWR = Pressurized Water Reactor  
HFIR: High Flux Isotope Reactor

Courtesy Bryce Littlejohn

# Global Fit of Existing Data



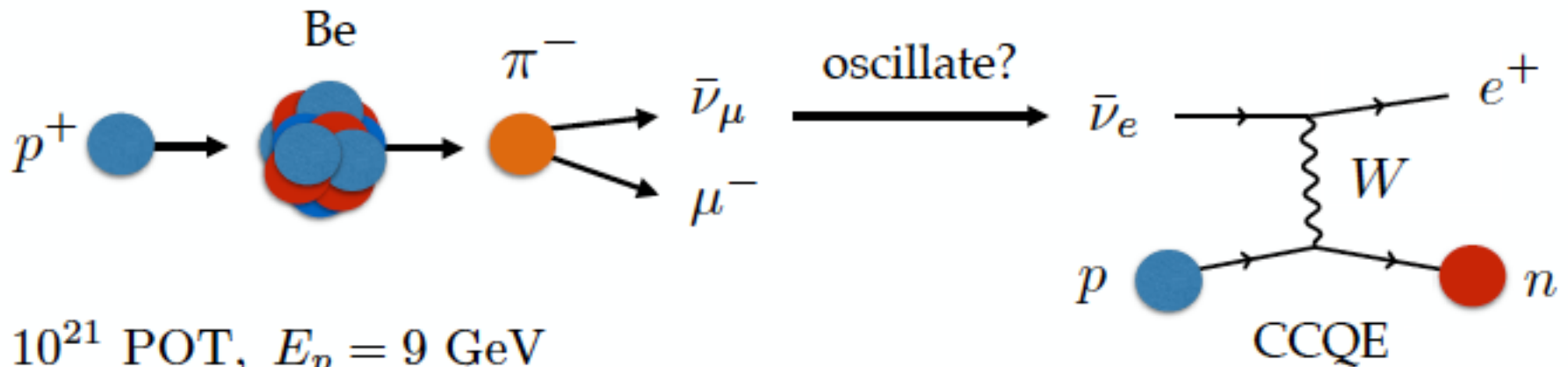
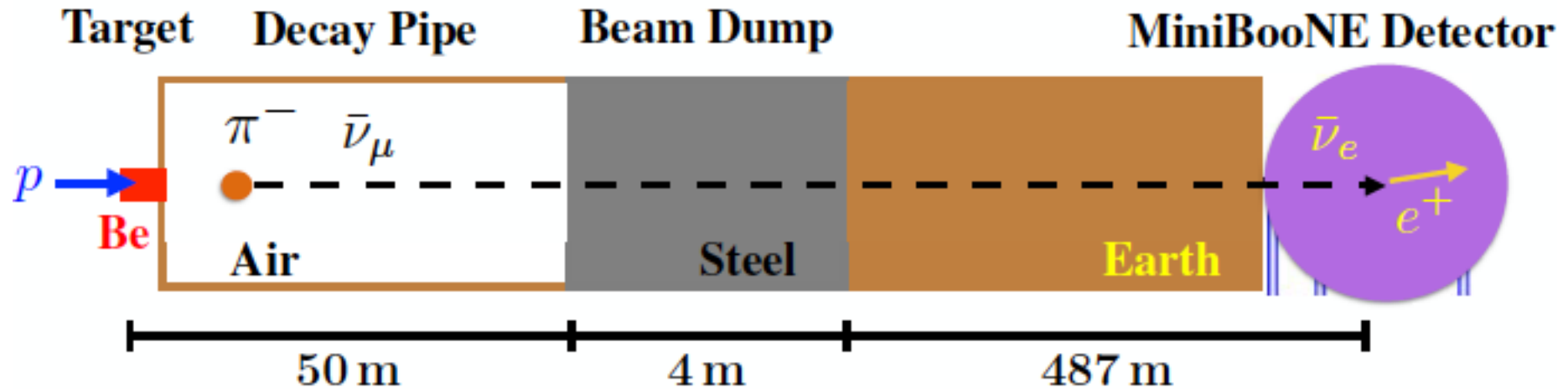
► Fit with NEOS/Daya Bay:  $\Delta\chi^2_{3\nu-4\nu} = 12.6 \Rightarrow 3.1 \sigma$



► Fit with NEOS/RENO:  $\Delta\chi^2_{3\nu-4\nu} = 9.1 \Rightarrow 2.6 \sigma$

Giunti, Li, Ternes, Tyagi, Xin, arXiv:2209.00916

# MiniBooNE Experiment



$10^{21}$  POT,  $E_p = 9$  GeV

Energy and baseline chosen to test LSND

Comparable oscillation probabilities

MiniBooNE Collaboration 1805.12028

Courtesy Gordan Krnjaic

# MiniBooNE Data Analysis

Luminosity	neutrino mode $12.84 \times 10^{20}$ POT	antineutrino mode $11.27 \times 10^{20}$ POT
Reconstructed Neutrino Energy	$200 < E_{\nu}^{QE} < 1250$ MeV	
Excess events BG subtracted	$381.2 \pm 85.2$	$79.3 \pm 28.6$

## Possibly Important Caveat

Mild tension  $\sim 2+$  sigma between neutrino and antineutrino modes

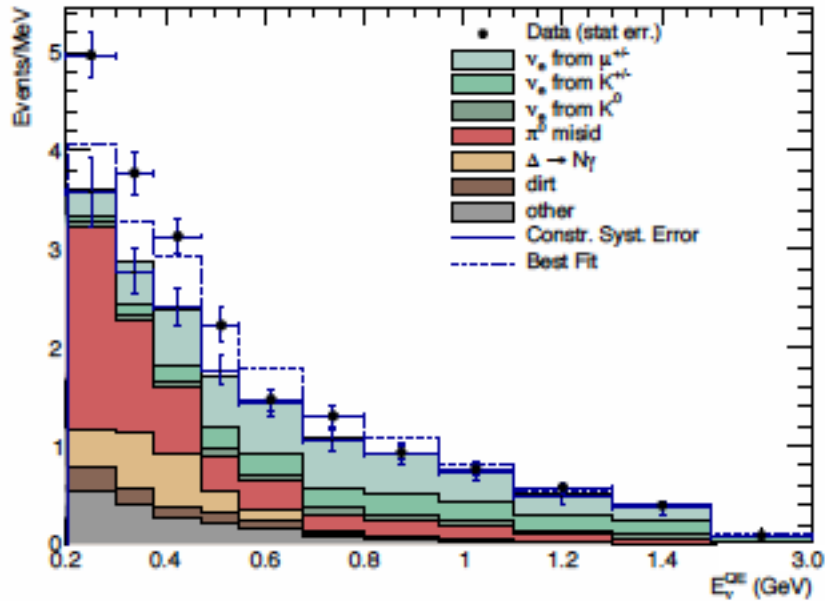
Updated Neutrino Mode Analysis  
MiniBooNE Collaboration 1805.12028

Complements earlier antineutrino  
results collected 2002-2010

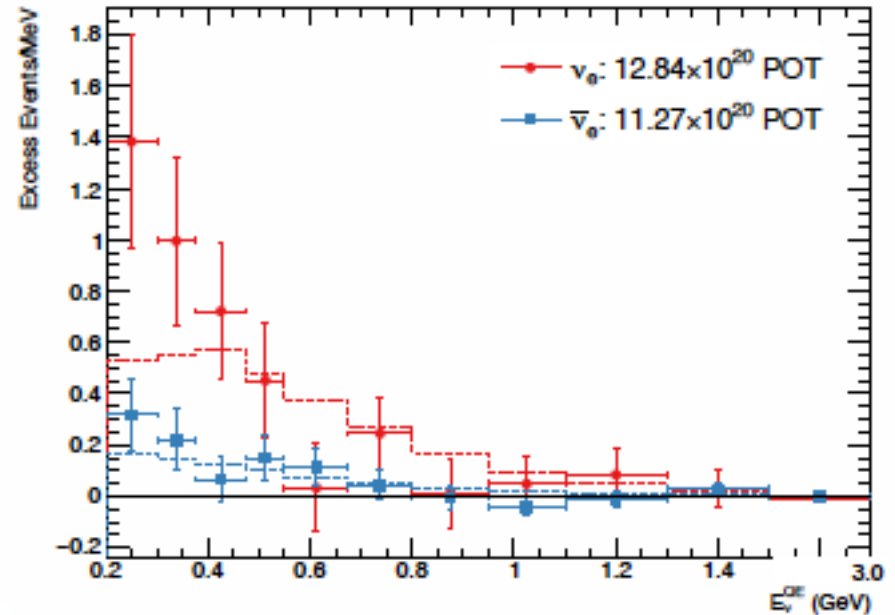
Courtesy Gordan Krnjaic

# MiniBooNE Anomaly

Neutrino mode only



Both excesses, BG subtracted



$$E_{\nu}^{(\text{reconst.})} = \frac{2m_n E_e + m_p^2 - m_n^2 - m_e^2}{2(m_n - E_e + \cos \theta_e \sqrt{E_e^2 - m_e^2})}$$

Measure charged lepton energy/angle

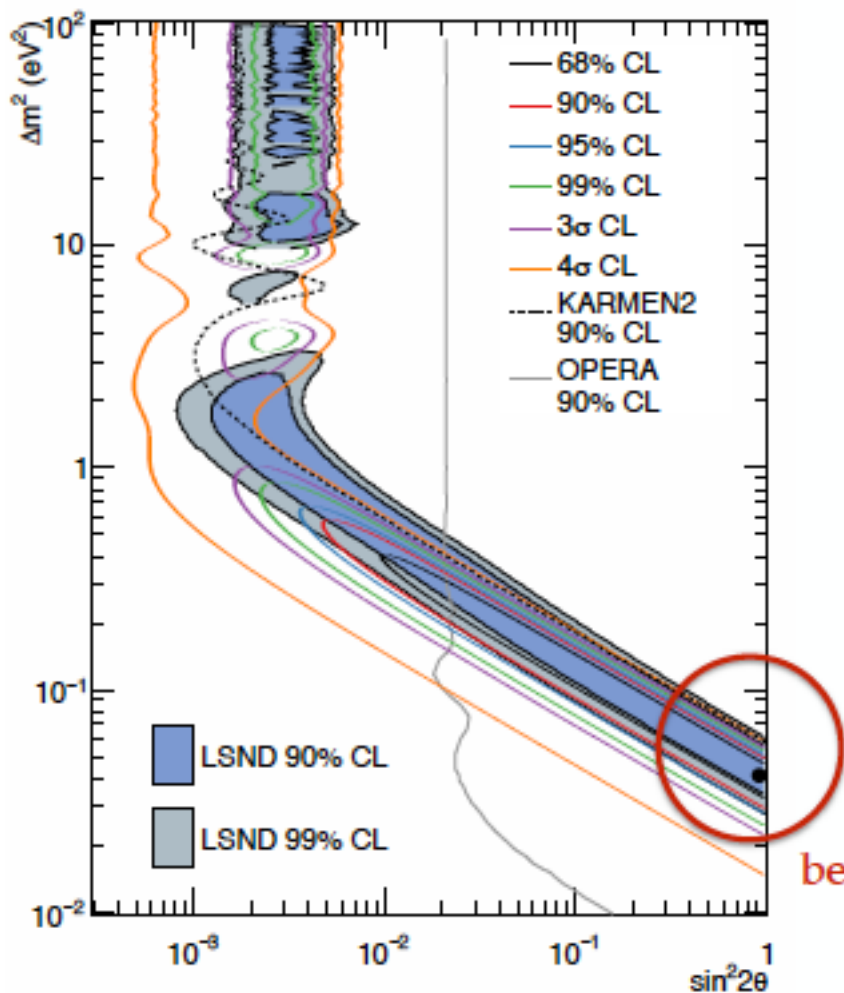
Observed  $\sim 400$  events, PMNS predicts 0

Combined  $\nu/\bar{\nu}$  modes :  $4.8\sigma$  excess

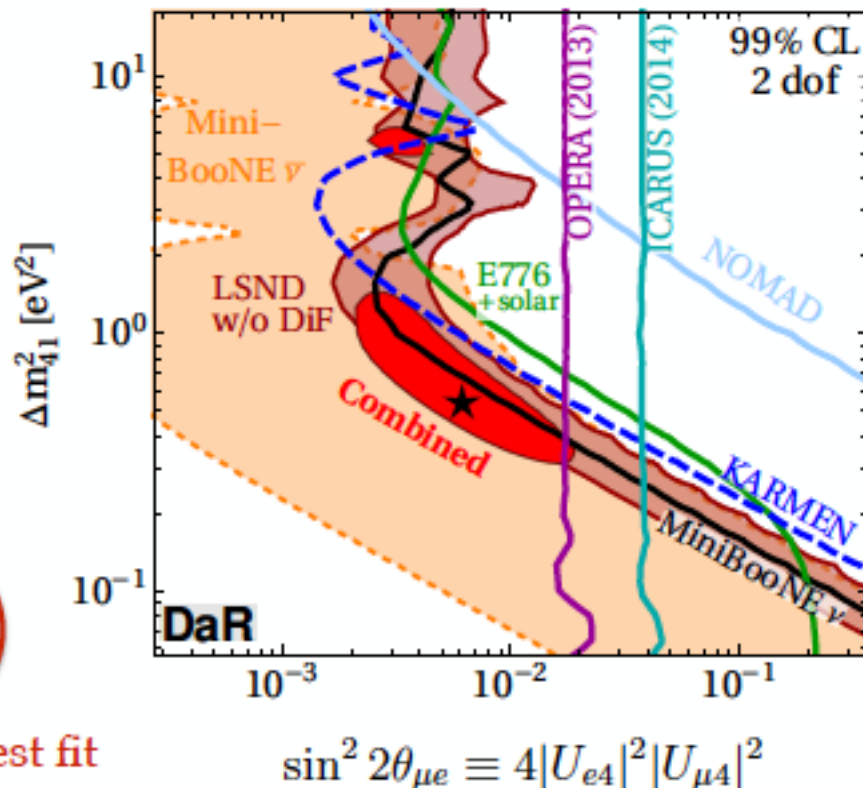
MiniBooNE Collaboration 1805.12028



# Can eV-scale Sterile Neutrino explain LSND+ MiniBooNE Anomaly?



Best fits [appearance only]

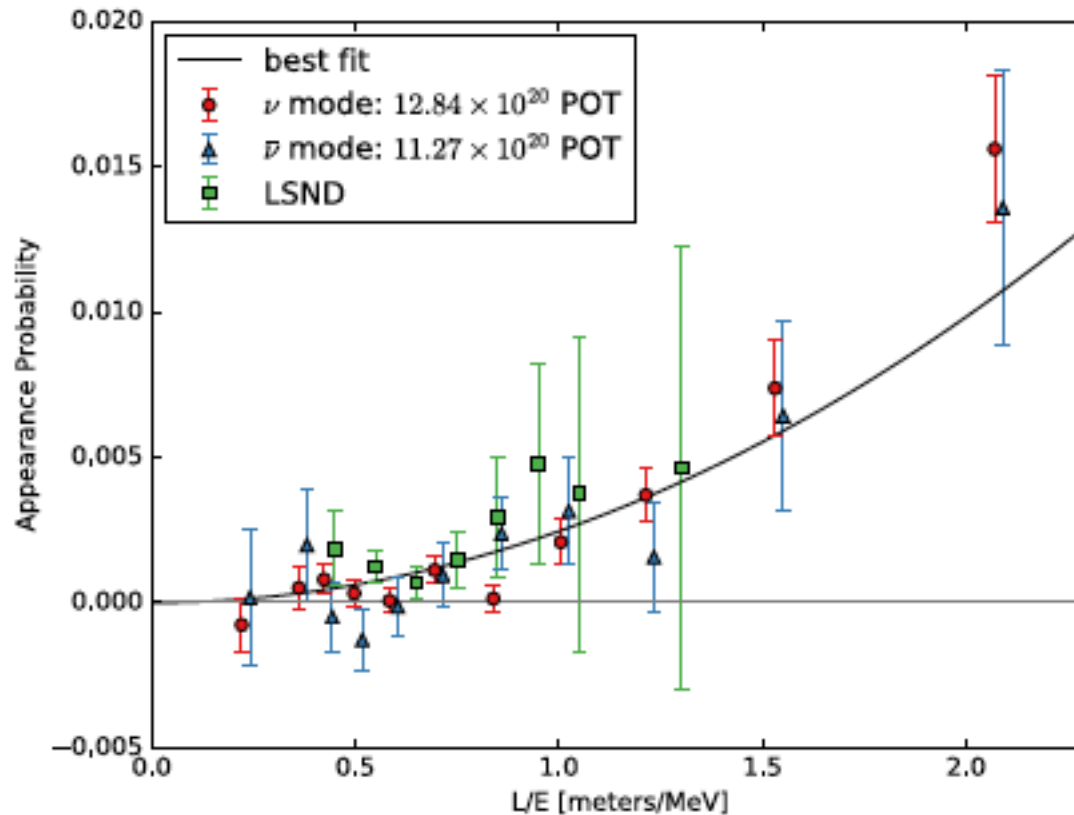


Combined 6.1 sigma excess!

MiniBooNE Collaboration 1805.12028

Dentler et. al. 1803.10661

# Can eV-scale Sterile Neutrino explain LSND+ MiniBooNE Anomaly?

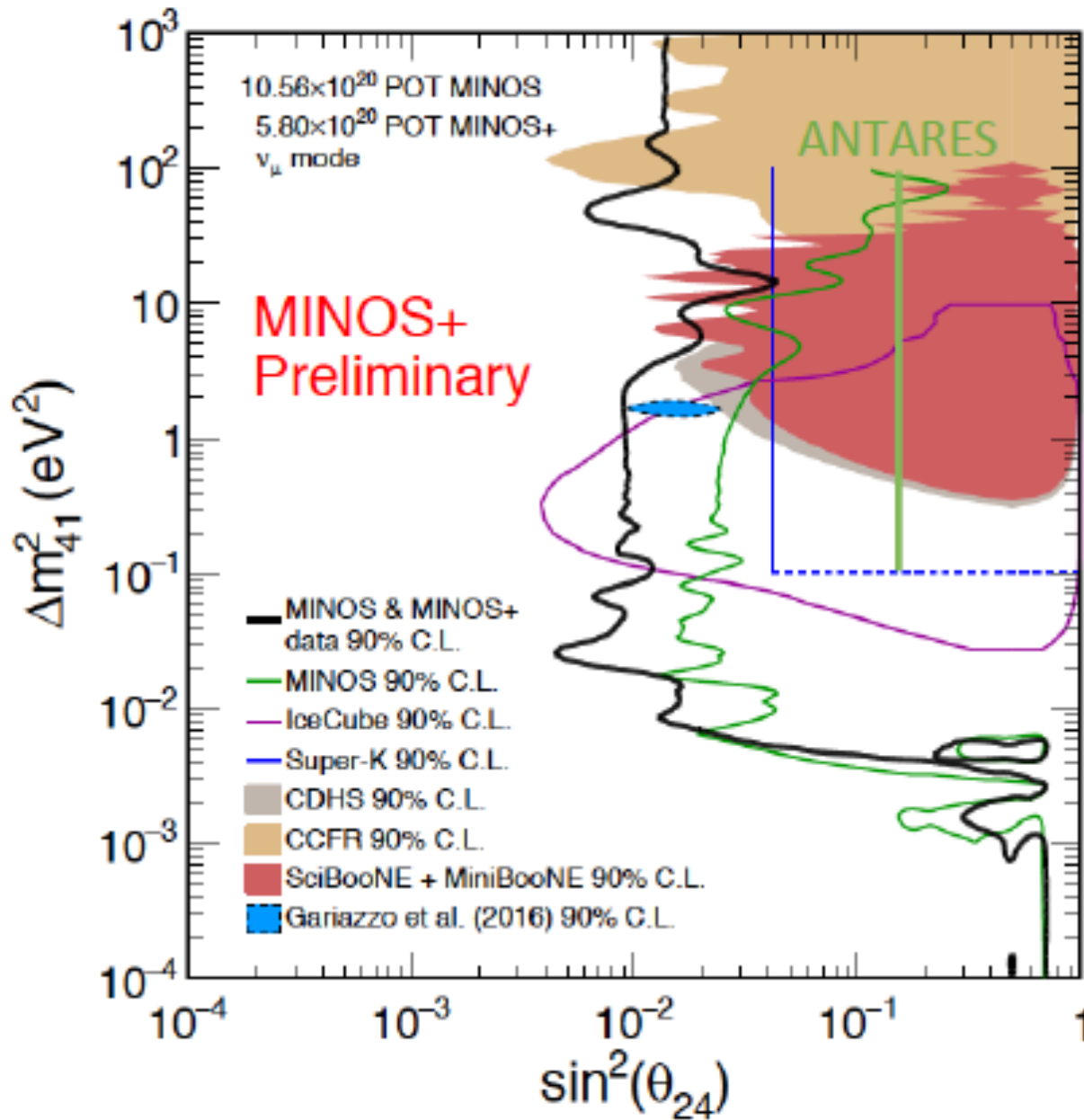


@ face value the scaling seems to support LSND/MB compatibility

MiniBooNE Collaboration 1805.12028

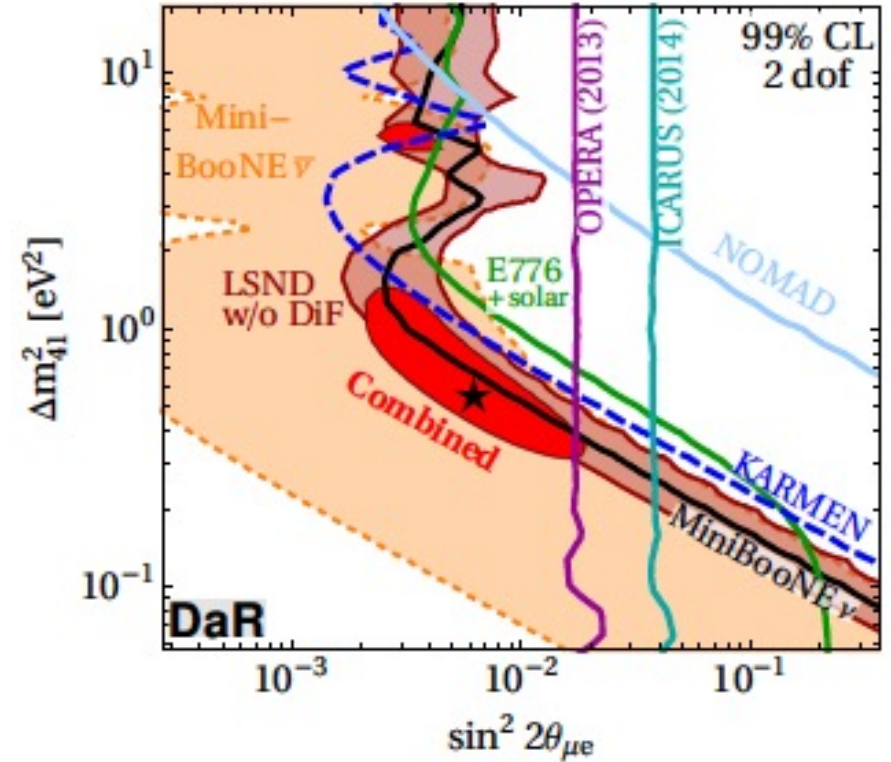
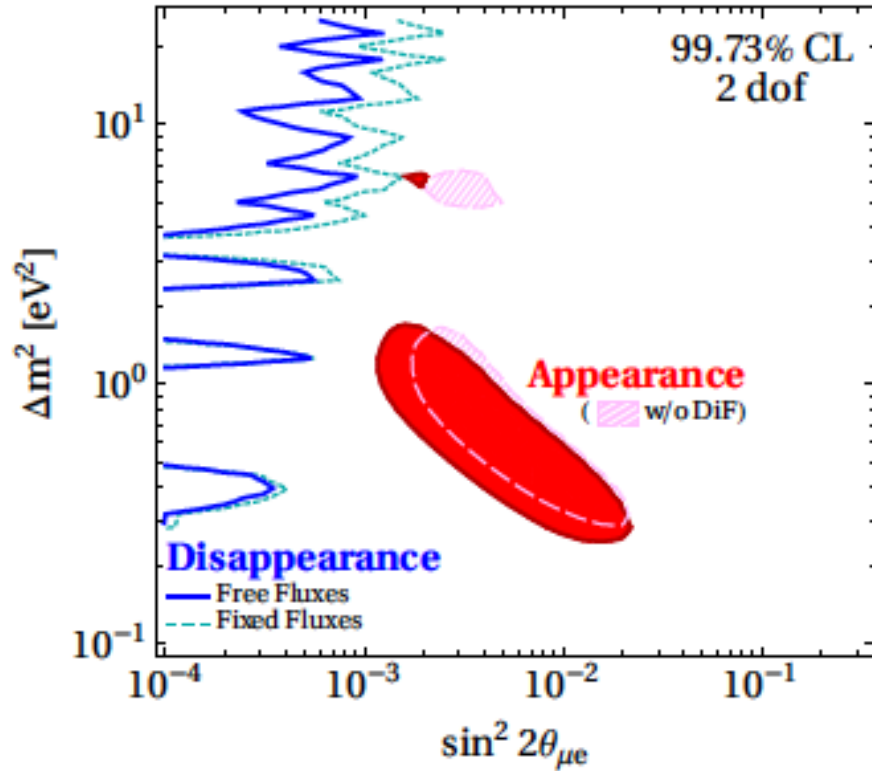
Courtesy Gordan Krnjaic

# Let us Zoom the Muon Disappearance Results



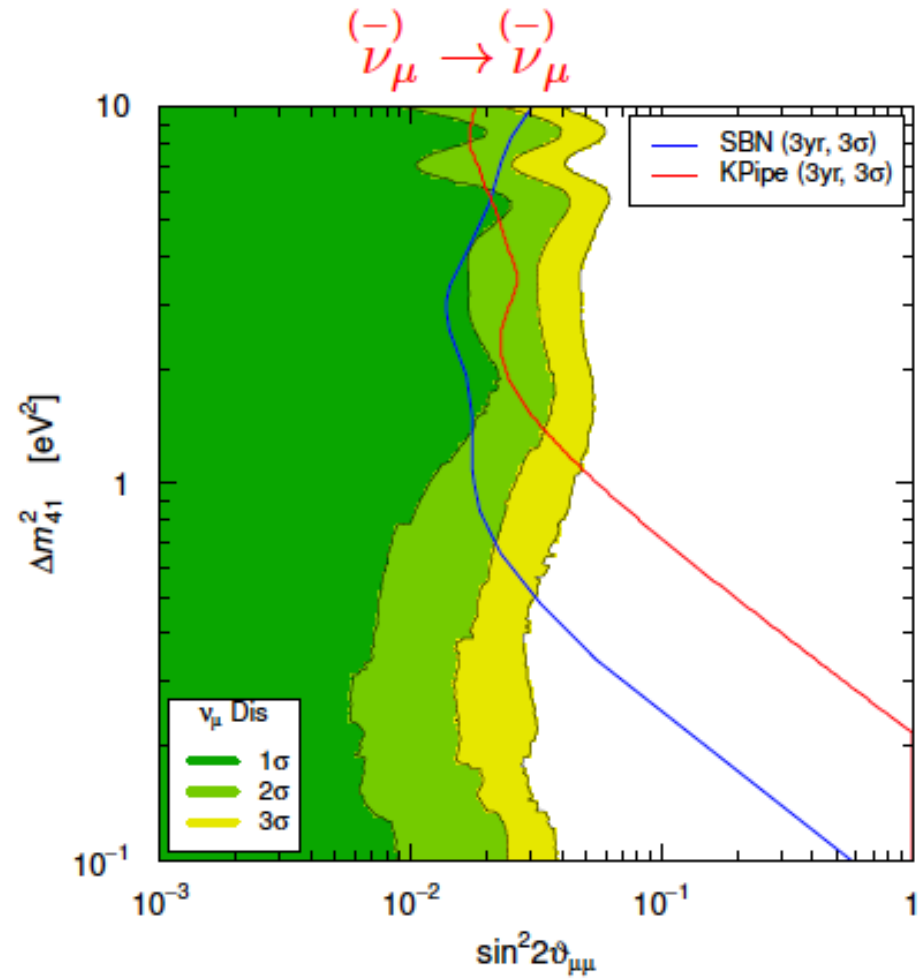
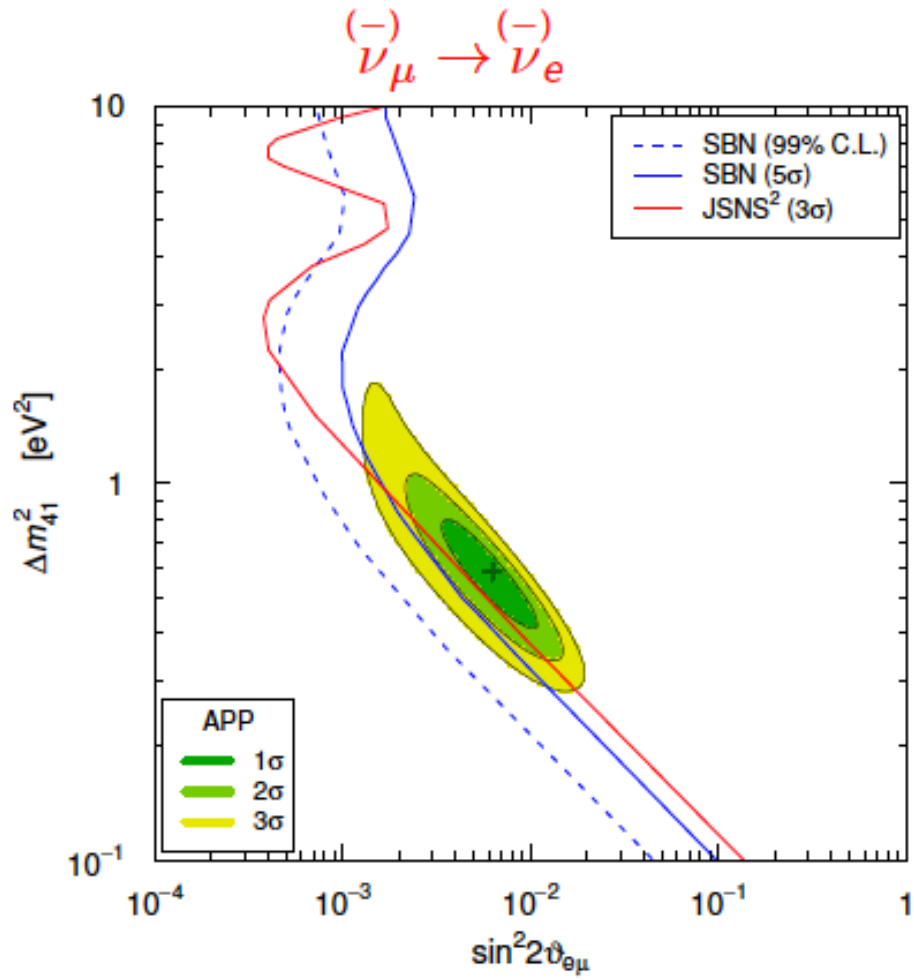
No anomaly!  
No evidence!  
Only constraints!

# Significant tension between appearance and disappearance results



Dentler, Hernandez-Cabezudo, Kopp, Machado, Maltoni, Martinez-Soler, Schwetz, JHEP 08 (2018) 010

# New Dedicated Experiments



Fermilab SBN Program: See Talk by Miquel Nebot-Guinot

Status of JSNS<sup>2</sup> and JSNS<sup>2</sup> - II: See Talk by ChangDong Shin

Attend the Topical Discussion on this topic

## Octant of $\theta_{23}$ in Danger with a Light Sterile Neutrino

Sanjib Kumar Agarwalla,<sup>1,2,\*</sup> Sabya Sachi Chatterjee,<sup>1,2,†</sup> and Antonio Palazzo<sup>3,4,‡</sup>

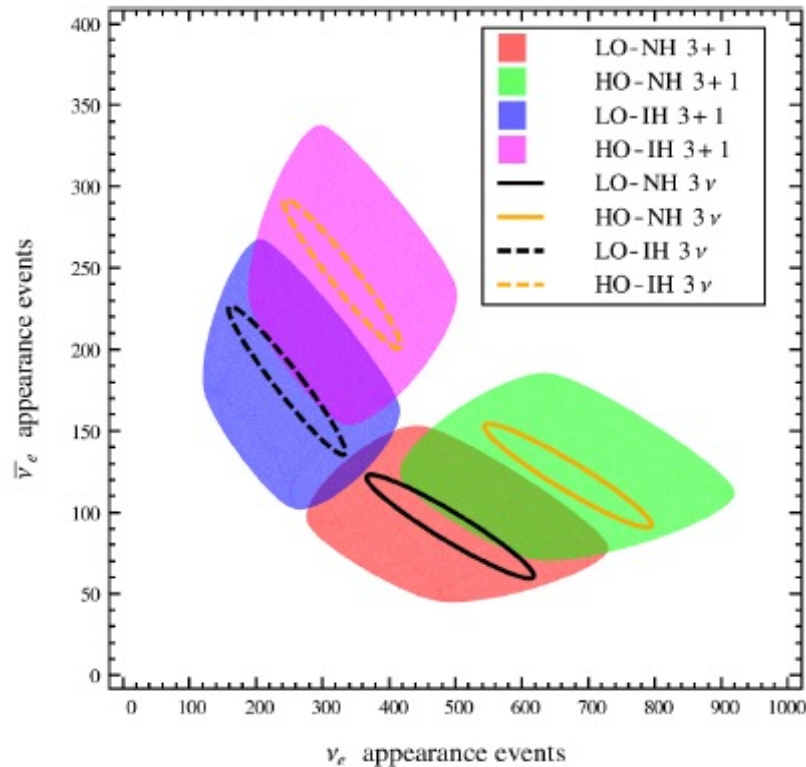
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<sup>2</sup>*Homi Bhabha National Institute, Training School Complex, Anushakti Nagar, Mumbai 400085, India*

<sup>3</sup>*Dipartimento Interateneo di Fisica "Michelangelo Merlin", Via Amendola 173, 70126 Bari, Italy*

<sup>4</sup>*Istituto Nazionale di Fisica Nucleare, Sezione di Bari, Via Orabona 4, 70126 Bari, Italy*

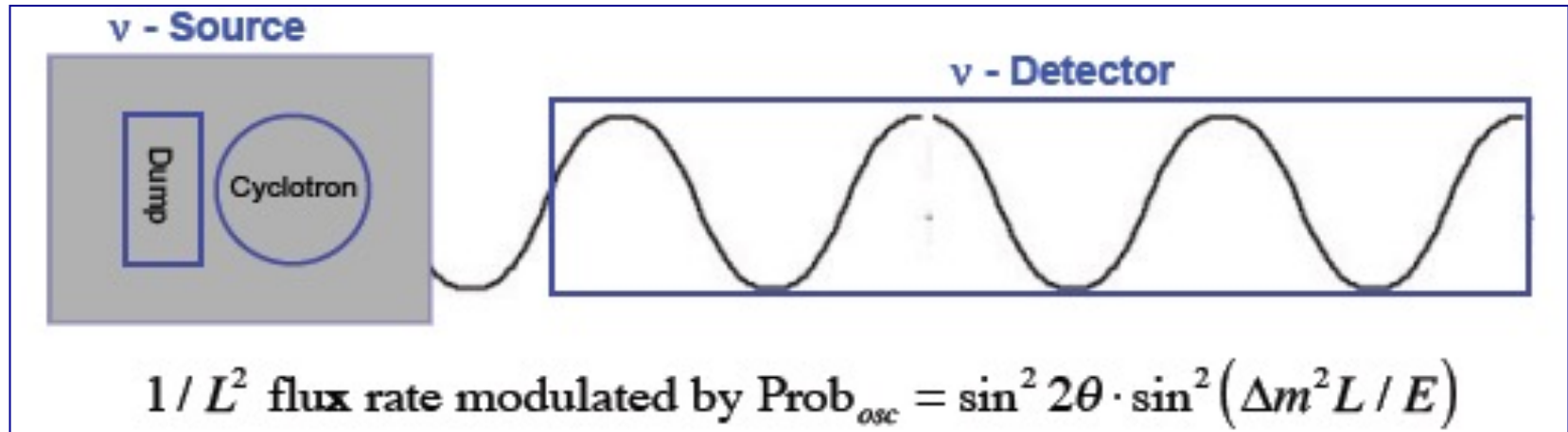
(Received 23 May 2016; revised manuscript received 5 December 2016; published 20 January 2017)



$$\sin^2\theta_{23} = 0.42 \text{ (LO)} \text{ and } 0.58 \text{ (HO)}$$

- Three-flavor ellipses due to variation in  $\delta_{13}$  in  $[-\pi \text{ to } \pi]$
- Four-flavor blobs due to variation in  $\delta_{13}$  and  $\delta_{14}$  in  $[-\pi \text{ to } \pi]$
- Due to new CP phases, sensitivity towards octant lost in DUNE

# Short-Baseline Neutrino Oscillation Wave



Agarwalla, Conrad, Shaevitz, arXiv:1105.4984 [hep-ph]; Agarwalla, Huber, arXiv:1007.3228 [hep-ph]

## Neutrino Sources

- Decay-at-rest beam from proton beam dump
- Small core reactor source
- Very high activity radioactive source

- Observe the  $L/E$  dependence of the event rates within a long  $\nu$  detector
- Background distribution is either independent of  $L$  or goes like  $1/L^2$
- Powerful verification of the short-baseline oscillation/new physics

# What Cosmology can infer about Light Sterile Neutrino?

- One eV-scale sterile neutrino is in severe tension with cosmology
- In standard scenario, sterile neutrino is produced via oscillations at  $T \gtrsim \text{MeV}$
- There are two key measurements from cosmology on massive neutrinos

□  $N_{\text{eff}} \lesssim 3.16$

- Measure the energy density due to relativistic particles
- Presence of an extra neutrino species would imply  $N_{\text{eff}} \sim 4$

□  $\sum m_\nu \lesssim 0.12 \text{ eV}$

- Sum of neutrino masses
- It affects structure formation and one light eV-scale sterile neutrino compatible with anomalies would suggest  $\sum m_\nu \sim 1 \text{ eV}$

Planck TT, TE, EE + low E + lensing + BAO, arXiv:1807.06209 [11053 Citations]



## *Concluding Remarks*

Light sterile neutrinos were interesting, are still interesting, and will remain interesting in near future as well.....

Discovery of a light sterile neutrino would prove that there is new physics beyond the SM at low energies, which is completely orthogonal to new physics searches at high energies at the LHC.....

Let us continue our effort to look for them at any mass scale and at any energies.....

*Thank you*

# What Cosmology can infer about Light Sterile Neutrino?

Sterile neutrinos require  $\sin^2 2\theta_{\mu e} > 10^{-3}$ ,  $m_4 < \text{few eV}$

Generic early universe thermalization

$$\Gamma > H \implies \sin^2 2\theta_{\mu e} G_F^2 T^5 > \sqrt{g_*} \frac{T^2}{m_{\text{Pl}}} \implies n_4 \sim n_\nu$$

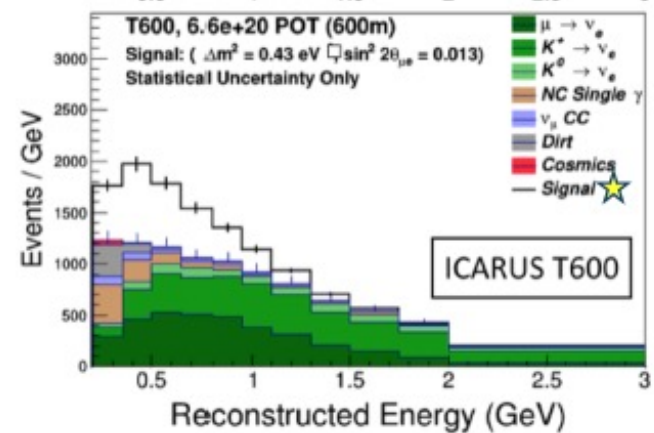
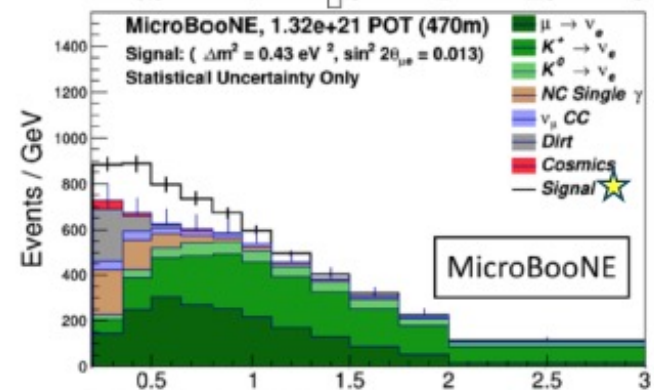
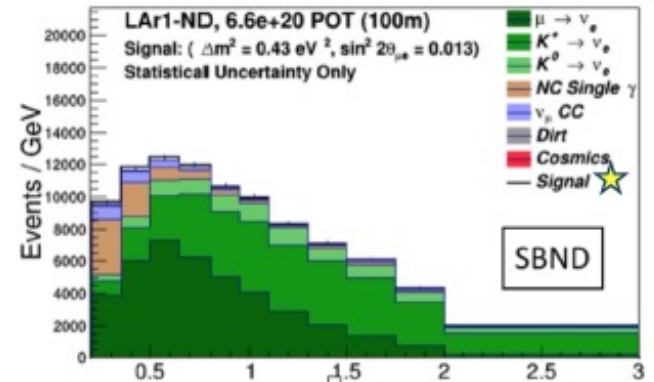
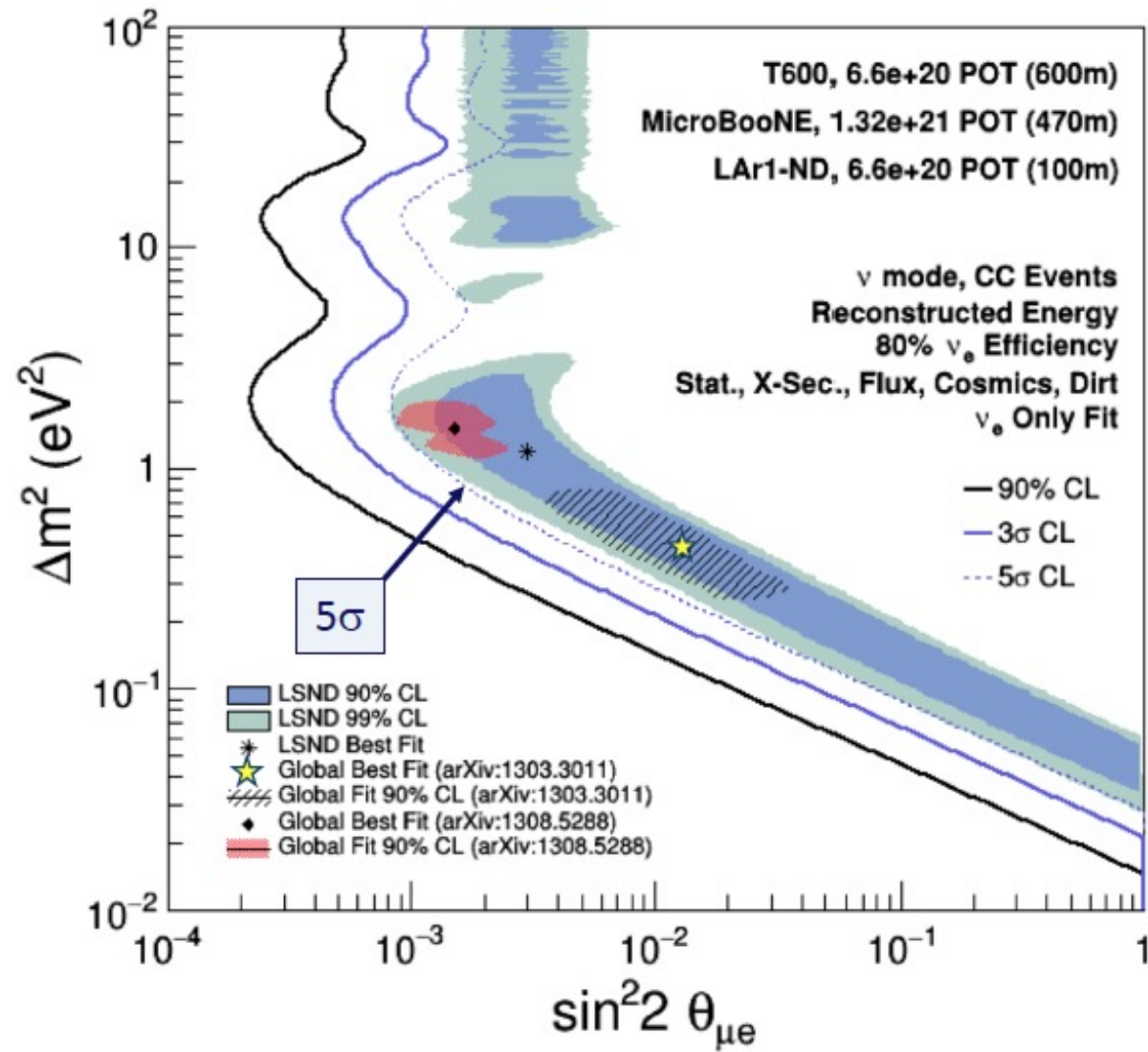
Excluded by BBN/CMB  $N_{\text{eff}} = 2.99 \pm 0.17$  Planck 1807.06209

Unless max temperature satisfies  $T_{\text{max}} \lesssim 15 \text{ MeV} \left( \frac{10^{-3}}{\sin^2 2\theta_{\mu e}} \right)^{1/3}$

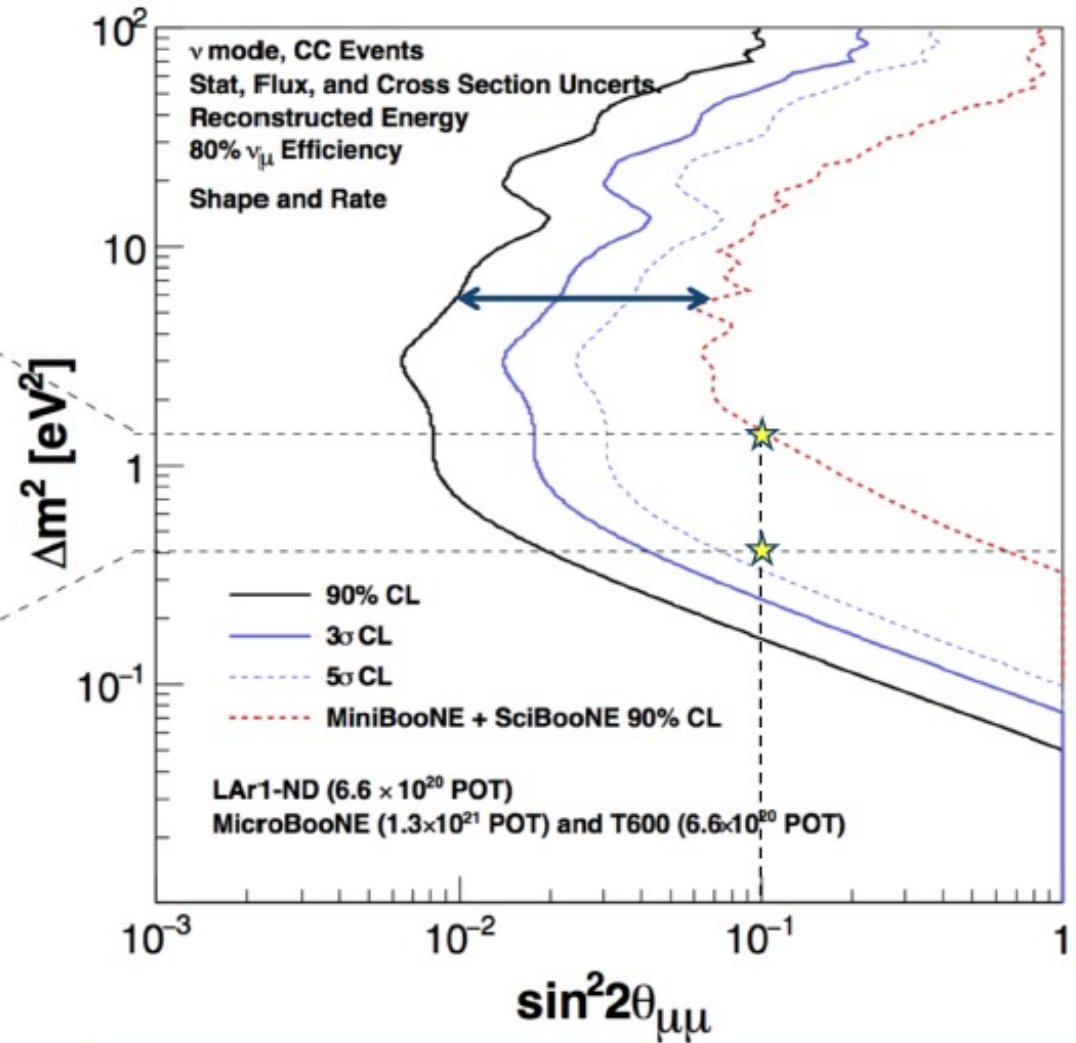
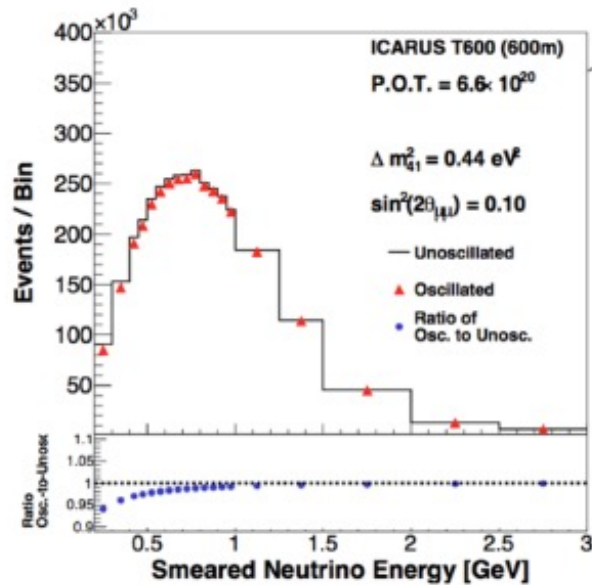
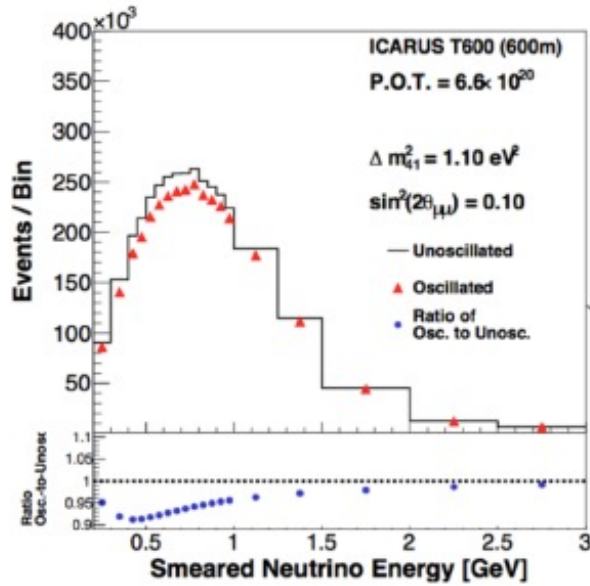
Courtesy Gordan Krnjaic

Interesting discussion in arXiv:1806.10629v1 by Chu, Dasgupta, Dentler, Kopp, Saviano

# Fermilab: Short-Baseline Neutrino Appearance Oscillation Sensitivity



# Fermilab: Short-Baseline $\nu$ Disappearance Oscillation Sensitivity



Very little background. Near detector key to controlling flux and cross-section uncertainties.